

Ab initio many-body calculations of single-nucleon transfer reactions with deuteron projectile

[arXiv:1602.04404]



FUSTIPEN
French-U.S. Theory Institute for Physics with Exotic Nuclei
GANIL, Caen, France (March 14-18 2016)

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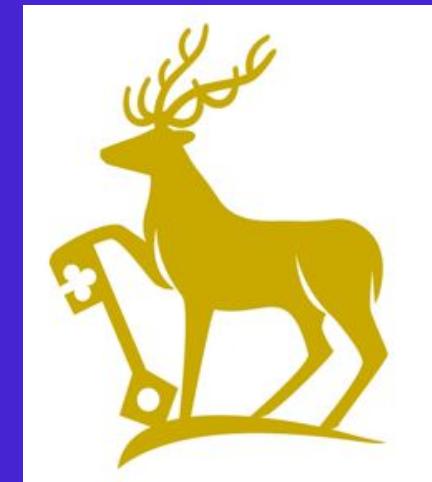
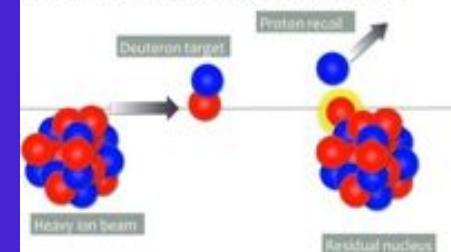


Figure 1 from Kate L Jones 2013 Phys. Scr. 2013 014020

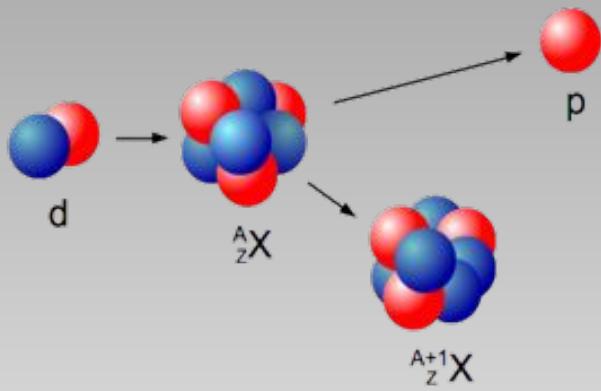


Outline

- Motivations for the study of transfer reactions and interest in ${}^7\text{Li}(d,p){}^8\text{Li}$ reaction
- The No-Core Shell Model with Resonating Group Method (NCSM/RGM) and with continuum (NCSMC)
- Results on ${}^7\text{Li}(d,p){}^8\text{Li}$ and ${}^7\text{Li}(d,d){}^7\text{Li}$ reactions and resonances of ${}^9\text{Be}$ above d - ${}^7\text{Li}$ threshold:
 - (Eigen)phase shifts
 - Cross sections
- Conclusions & perspectives

Deuteron-nucleus reaction: experimental motivations

Intense experimental activity (direct and inverse kinematics):



(d,p) reaction
in direct
kinematics

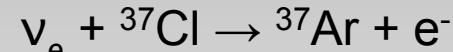
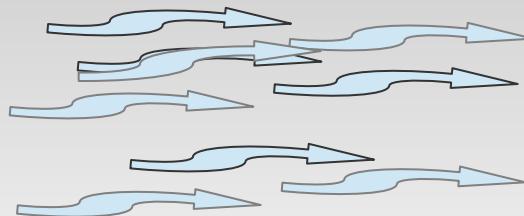
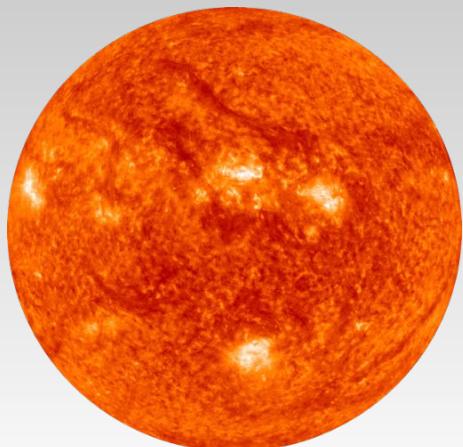
- Structure and spectroscopy of nuclei
- Nucleosynthesis and nuclear fusion applications ($^3\text{H}(d,n)^4\text{He}$ reaction)
- Surrogate for (p/n) capture reactions
- Calibration reaction for measurement of processes of interest

$^7\text{Li}(d,p)^8\text{Li}$ transfer reaction



Calibration reaction for astrophysical process: $^7\text{Li}(d,p)^8\text{Li}$ as target calibration for $^7\text{Be}(p,\gamma)^8\text{B}$

Solar neutrino problem:



R. Davis Jr takes a dip
At Homestake Mine (1971)

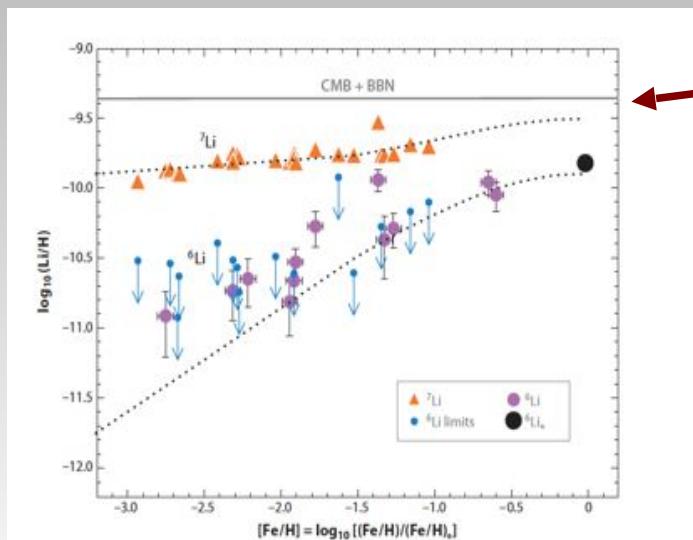
$^7\text{Li}(d,p)^8\text{Li}$ transfer reaction



Calibration reaction for astrophysical process: $^7\text{Li}(d,p)^8\text{Li}$ as target calibration for $^7\text{Be}(p,\gamma)^8\text{B}$



Possible mechanism of destruction of ^7Li in the context of baryon-inhomogeneous models of the primordial nucleosynthesis



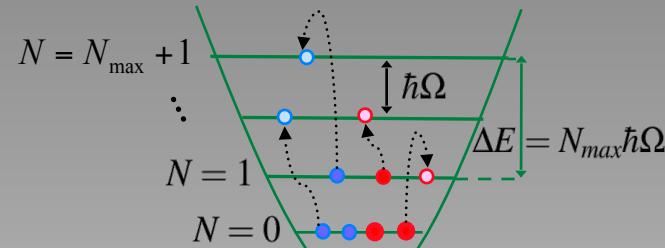
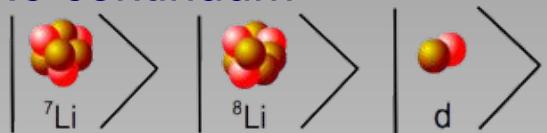
Primordial Lithium problem:

- 4-5 σ discrepancy between observed and calculated (CMB+BBN) abundance of ^7Li
- Nuclear solution to the problem: d - ^7Li destruction mechanism is ruled out (but only in a standard BBN scenario PRC 47, 2369 1993)

No-core shell model combined with the resonating group method (NCSM-RGM) and NCSM with continuum (NCSMC)

No-core shell model (NCSM):

- > A-nucleon wave function expansion in the harmonic-oscillator (HO) basis
- > Short- and medium-range correlations
- > No continuum

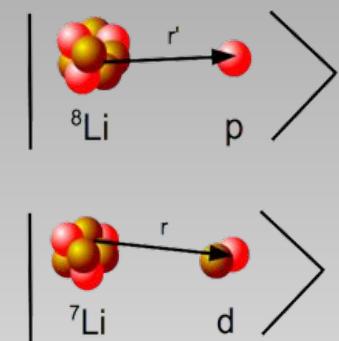


P. Navrátil et al. PRL 84, 5728 (2000)

NCSM+Resonating group method (NCSM-RGM):

- > Microscopic approach to describe the scattering of clusters
- > Long range correlations (relative motion of clusters)

K. Wildermuth, Y.C. Tang A unified theory of the nucleus 1977



NCSM with continuum (NCSMC):

S. Baroni, P. Navrátil, and S. Quaglioni, PRL 110, 022505 (2013); PRC 87, 034326 (2013)

$$|\Psi_A^{J^\pi T} \rangle = \sum_\lambda c_\lambda |{}^9\text{Be} \rangle + \sum_{\tilde{\nu}} \int dr r^2 \frac{g_{\tilde{\nu}}^{J^\pi T}(r)}{r} \hat{A}_{\tilde{\nu}} |{}^7\text{Li} \rangle_d + \sum_{\tilde{\nu}'} \int dr' r'^2 \frac{g_{\tilde{\nu}'}^{J^\pi T}(r')}{r'} \hat{A}_{\tilde{\nu}'} |{}^8\text{Li} \rangle_p$$

Variational amplitudes
(unknowns of the many-body problem)

NCSM-RGM and NCSMC equations

$$\mathcal{H} = T_{\text{rel}}(r) + \mathcal{V}_{\text{rel}} + \bar{V}_C(r) + H_{(A-a)} + H_{(a)}$$

Internal A-nucleon
microscopic Hamiltonian

Coupled-channel equations solved for the amplitude c_λ and g_v

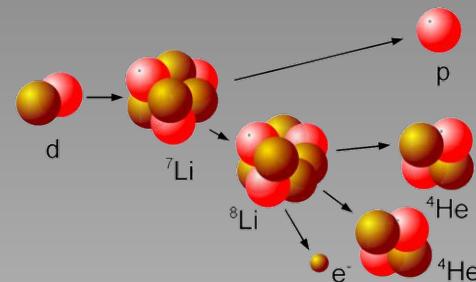
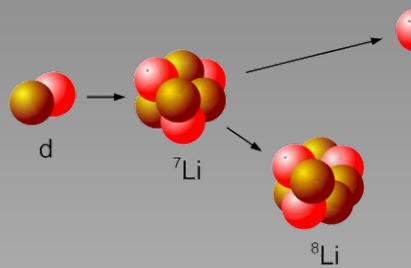
$$\left(\begin{array}{c} E_\lambda \delta_{\lambda\lambda'} \\ \langle \text{Be}^9 | \mathcal{H} \mathcal{A}_{\bar{\nu}} | \text{Li}^8 \text{ Li}^7 \text{ d} \rangle \\ \langle \text{Li}^7 \text{ d} \text{ Li}^8 \text{ p} | \mathcal{A}_{\bar{\nu}'} \mathcal{H} | \text{Be}^9 \rangle \end{array} \right) \left(\begin{array}{c} c_\lambda \\ g_{\bar{\nu}, \bar{\nu}'} \end{array} \right) = \left(\begin{array}{c} \langle \text{Be}^9 | \mathcal{H} \mathcal{A}_{\bar{\nu}} | \text{Li}^8 \text{ Li}^7 \text{ d} \rangle \\ \langle \text{Li}^7 \text{ d} \text{ Li}^8 \text{ p} | \mathcal{A}_{\bar{\nu}'} \mathcal{H} \mathcal{A}_{\bar{\nu}} | \text{Be}^9 \rangle \\ \langle \text{Li}^7 \text{ d} \text{ Li}^8 \text{ p} | \mathcal{A}_{\bar{\nu}'} \mathcal{H} \mathcal{A}_{\bar{\nu}} | \text{Li}^8 \text{ Li}^7 \text{ d} \rangle \end{array} \right) \left(\begin{array}{c} c_\lambda \\ g_{\bar{\nu}, \bar{\nu}'} \end{array} \right)$$

$$E \left(\begin{array}{c} \delta_{\lambda\lambda'} \\ \langle \text{Li}^7 \text{ d} \text{ Li}^8 \text{ p} | \mathcal{A}_{\bar{\nu}'} | \text{Be}^9 \rangle \\ \langle \text{Li}^7 \text{ d} \text{ Li}^8 \text{ p} | \mathcal{A}_{\bar{\nu}'} \mathcal{A}_{\bar{\nu}} | \text{Be}^9 \rangle \end{array} \right) \left(\begin{array}{c} c_\lambda \\ g_{\bar{\nu}, \bar{\nu}'} \end{array} \right)$$

Coupled-channel microscopic *R*-matrix method on Lagrange mesh provides Scattering matrix and Asymptotic Normalization Coefficients by matching internal solution to known asymptotic

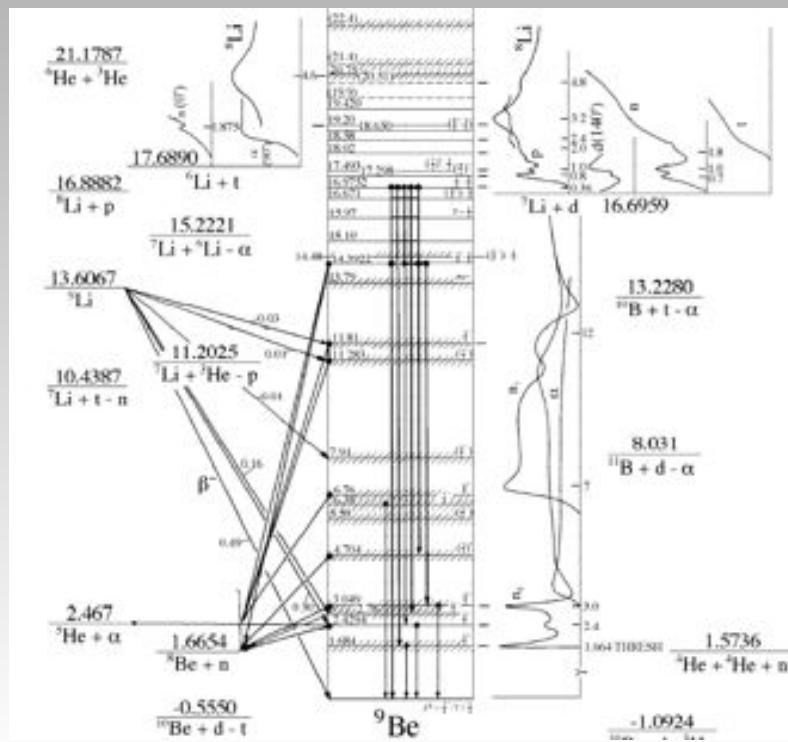
M. Hesse, J.M. Sparenberg, F. Van Raemdonck, and D. Baye, Nucl Phys. A 640, 37 (1988)

$^7\text{Li}(d,p)^8\text{Li}$ reaction and structure of ^9Be



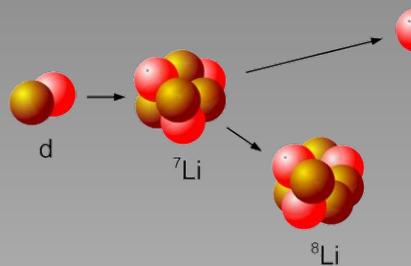
“Model space” reaction

“Real world” reaction

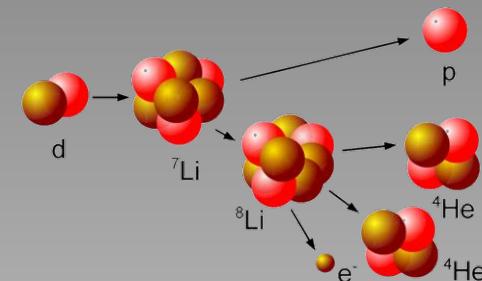


${}^9\text{Be}$ ground state is stable
All excited states are unbound

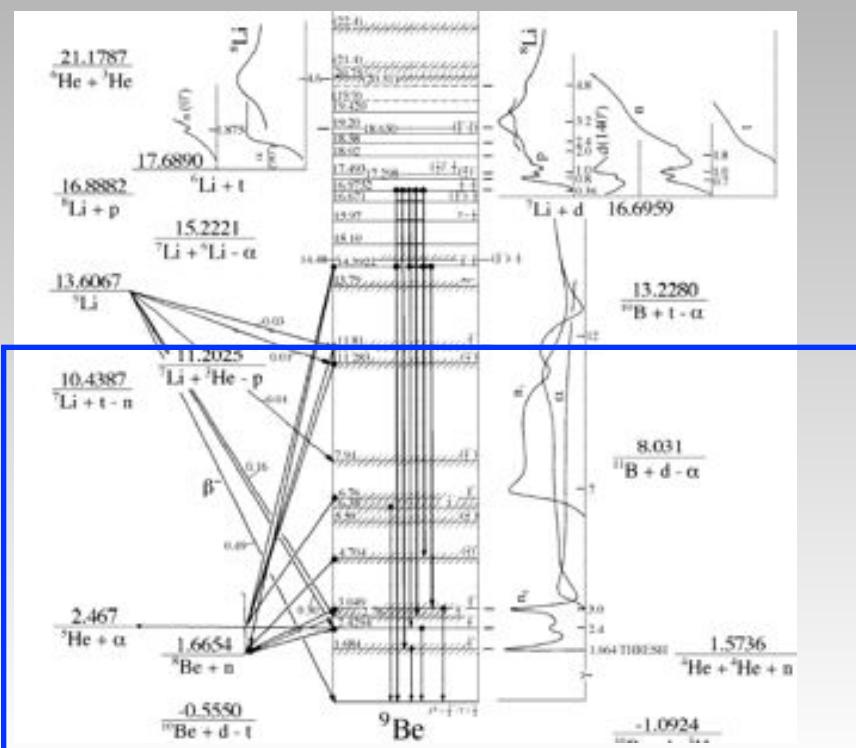
$^7\text{Li}(d,p)^8\text{Li}$ reaction and structure of ^9Be



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“Real world” reaction



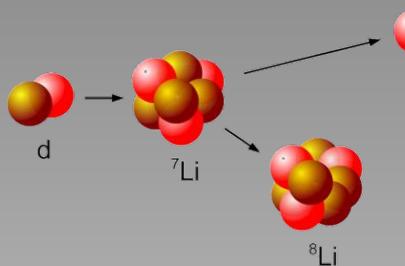
^9Be ground state is stable
All excited states are unbound

Inclusion of the continuum:

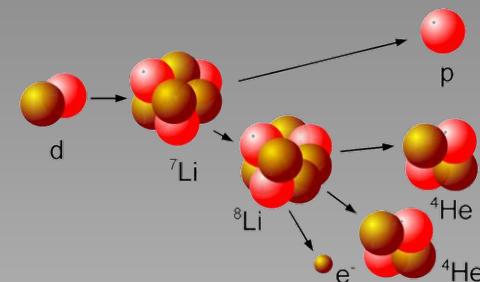
- Low-energy spectrum: $n-^8\text{Be}$ ($n-\alpha-\alpha$)

[J. Langhammer, P. Navrátil, S. Quaglioni, G. Hupin, A. Calci, and R. Roth, PRC(R) 91, 021301 (2015)]

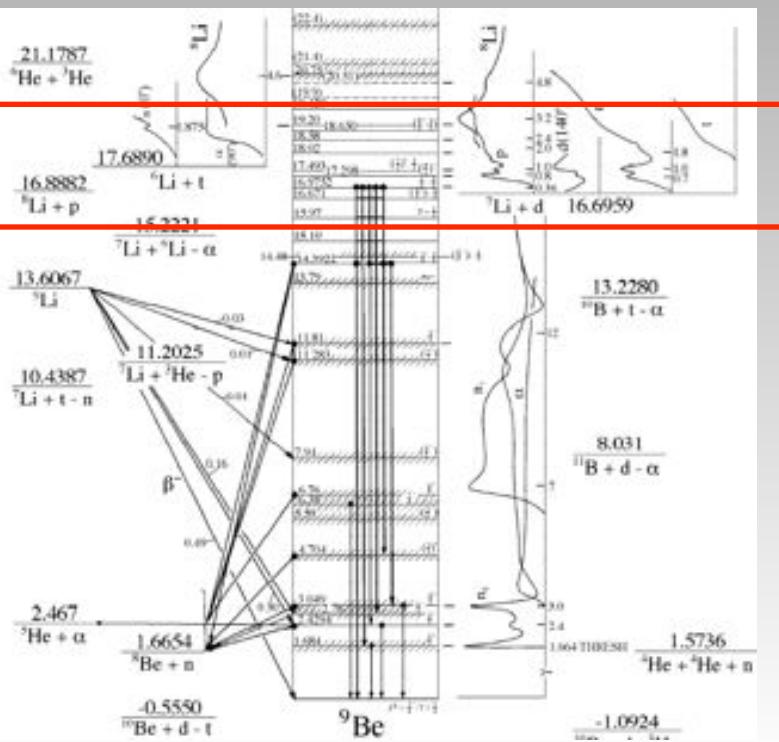
$^7\text{Li}(d,p)^8\text{Li}$ reaction and structure of ^9Be



“Model space” reaction



“Real world” reaction

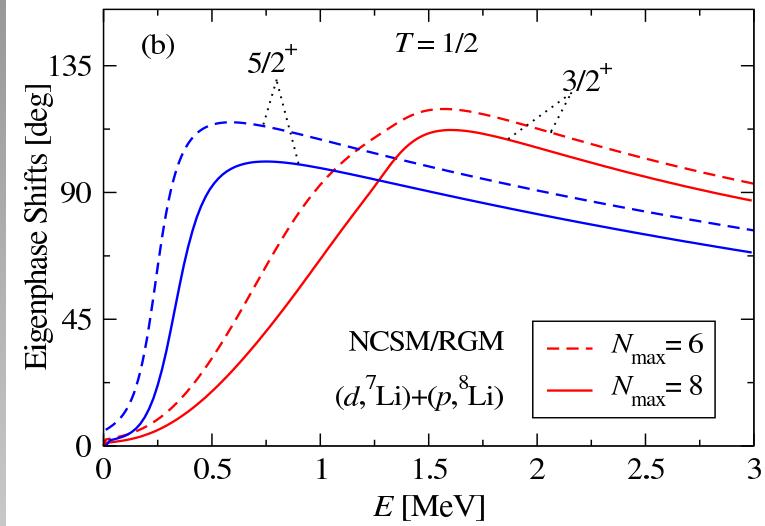
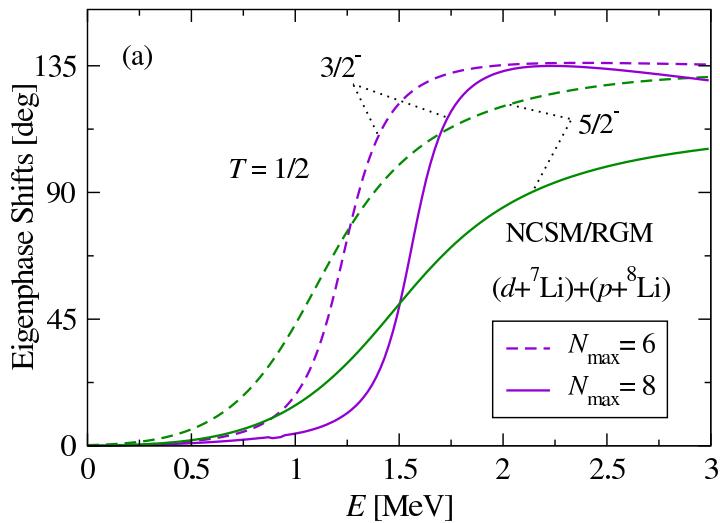


${}^9\text{Be}$ ground state is stable
All excited states are unbound

Inclusion of the continuum:

- Low-energy spectrum: $n-{}^8\text{Be}$ ($n-\alpha-\alpha$)
- High-energy spectrum: $d-{}^7\text{Li}$, $p-{}^8\text{Li}$

$(d, {}^7\text{Li}) + (\text{p}, {}^8\text{Li})$ coupled NCSM-RGM calculation Eigenphase shifts



Model space ($N_{\max}=6,8$ $\hbar\Omega=20$ MeV):

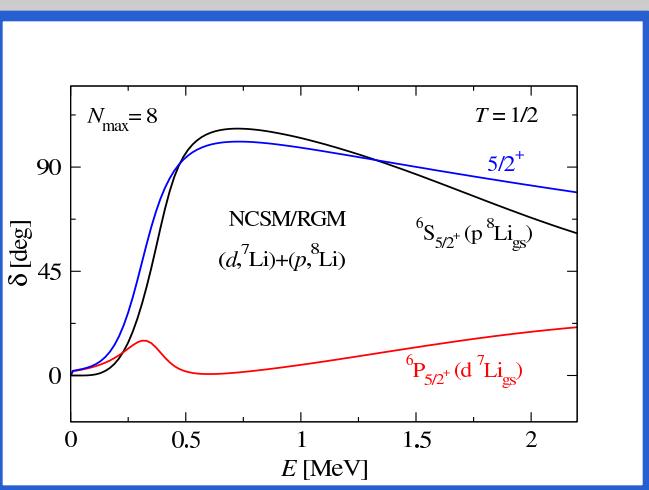
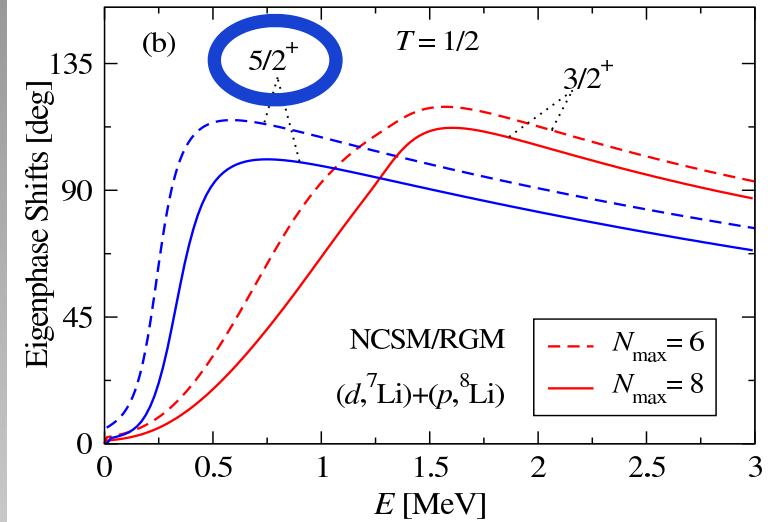
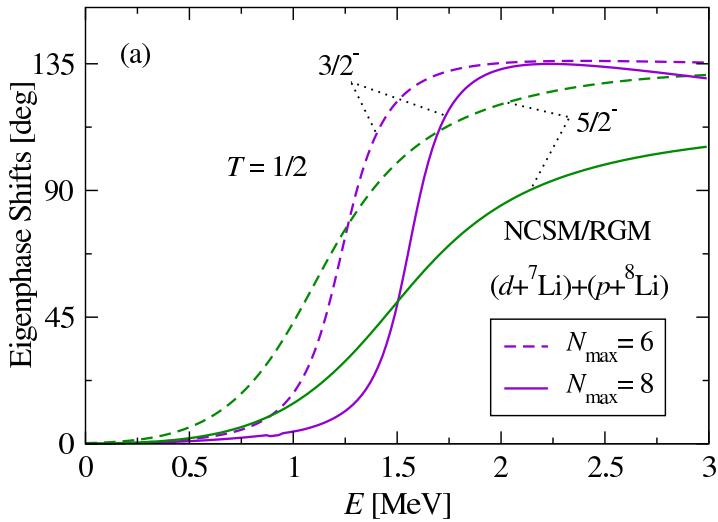
$$|\text{d}(\text{d}^*)+{}^7\text{Li}_{\text{gs}}\rangle + |\text{d}(\text{d}^*)+{}^7\text{Li}_{1\text{ex}}\rangle + |\text{p}+{}^8\text{Li}_{\text{gs}}\rangle + |\text{p}+{}^8\text{Li}_{1\text{ex}}\rangle + |\text{p}+{}^8\text{Li}_{2\text{ex}}\rangle + |\text{p}+{}^8\text{Li}_{3\text{ex}}\rangle$$

Virtual breakup of the deuteron: 4 pseudostates

Chiral nuclear interaction:

Entem-Machleidt SRG-evolved ($\Lambda=2.02$ fm $^{-1}$) NN force at N³LO (cutoff 500 MeV)

$(d, {}^7\text{Li}) + (\text{p}, {}^8\text{Li})$ coupled NCSM-RGM calculation Eigenphase shifts

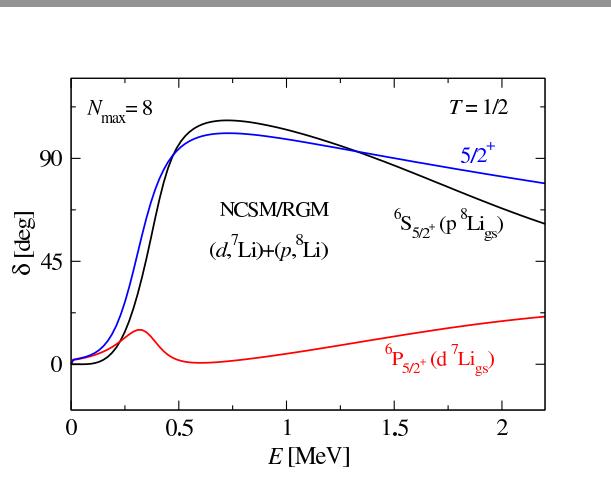


Dominant partial waves above p + ${}^8\text{Li}$ threshold: $3/2^{-,+}$, $5/2^{-,+}$

Main phase shifts for $5/2^+$:

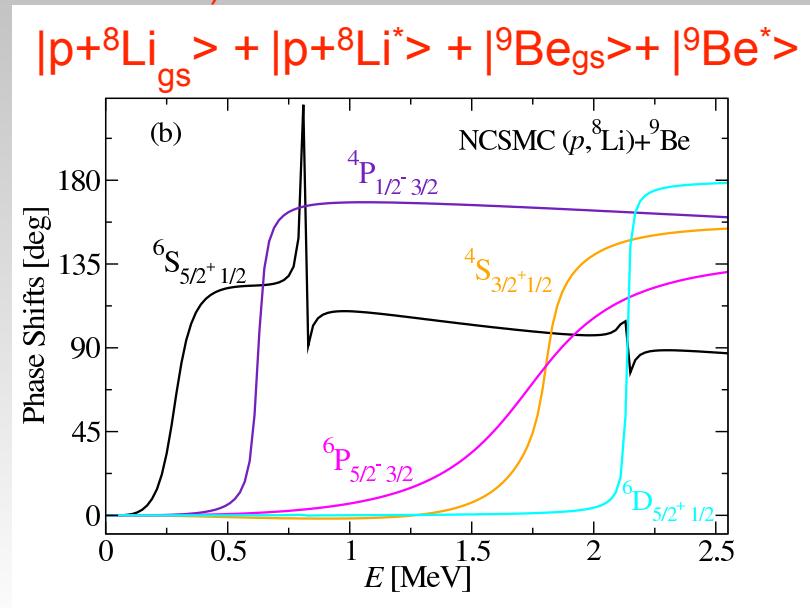
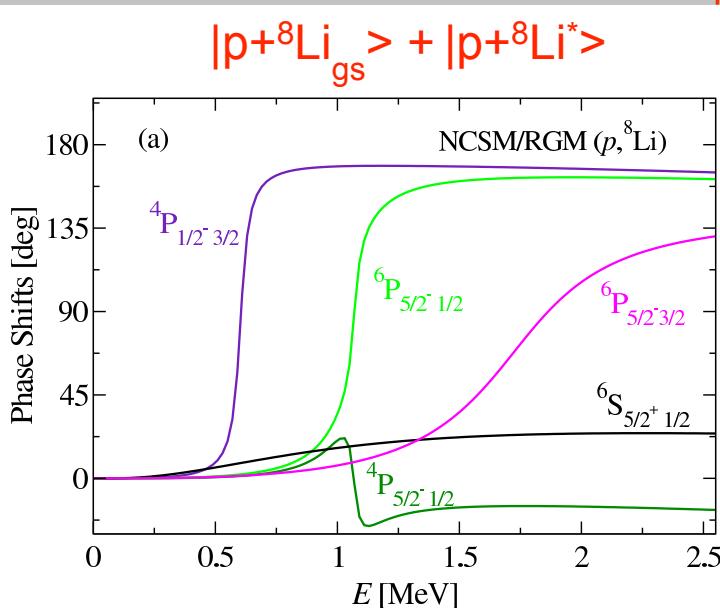
- P-wave in $(d, {}^7\text{Li})$
- Resonant S-wave in $(\text{p}, {}^8\text{Li})$

(p, ${}^8\text{Li}$) UNcoupled NCSM-RGM calculation Phase shifts

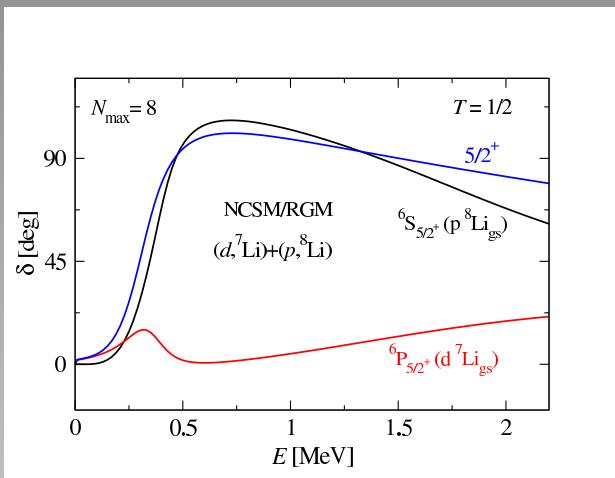


- $T=3/2$ resonances ($1/2^-$, $5/2^-$) reproduced.
- S-wave phase shift in $5/2^+$ strongly suppressed in NCSM-RGM calculation.
- Effects of the short-range correlations in NCSMC calculation: 1) Resonant S-wave in $5/2^+$ enhanced; 2) P-wave in $J=5/2^-$ e $T=1/2$ becomes bound.

Model space ($N_{\max}=8$ $\hbar\Omega=20$ MeV):

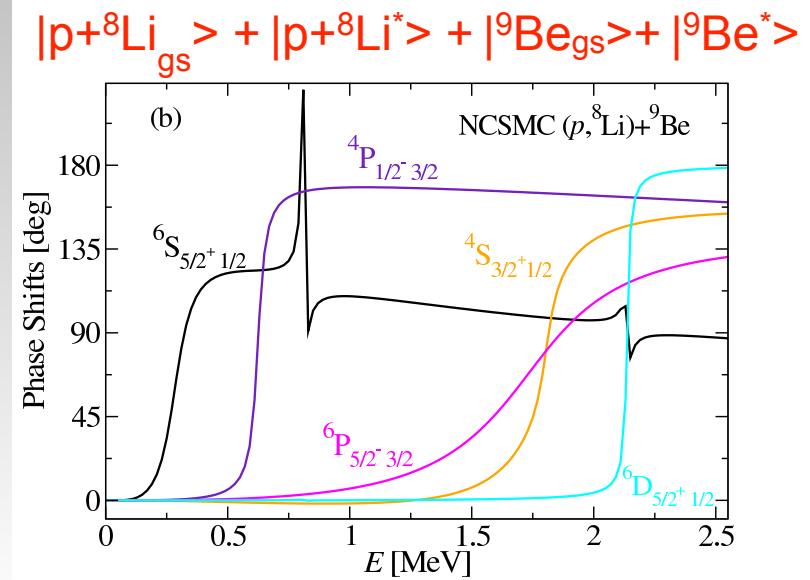
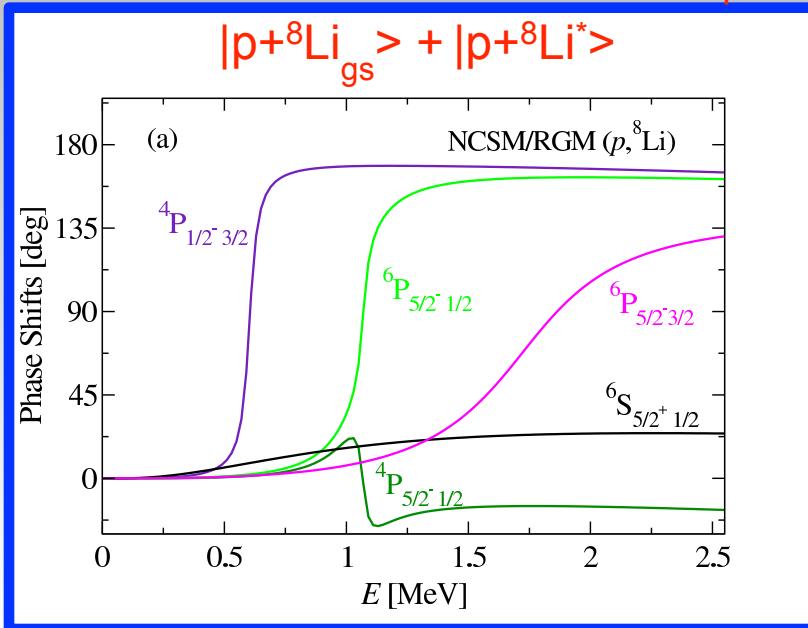


(p, ${}^8\text{Li}$) UNcoupled NCSM-RGM calculation Phase shifts

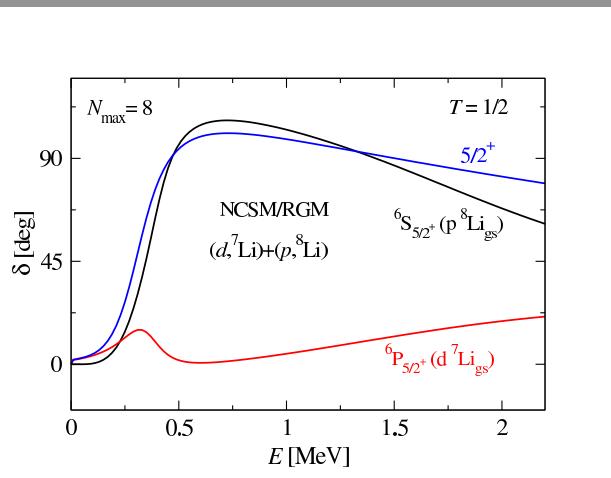


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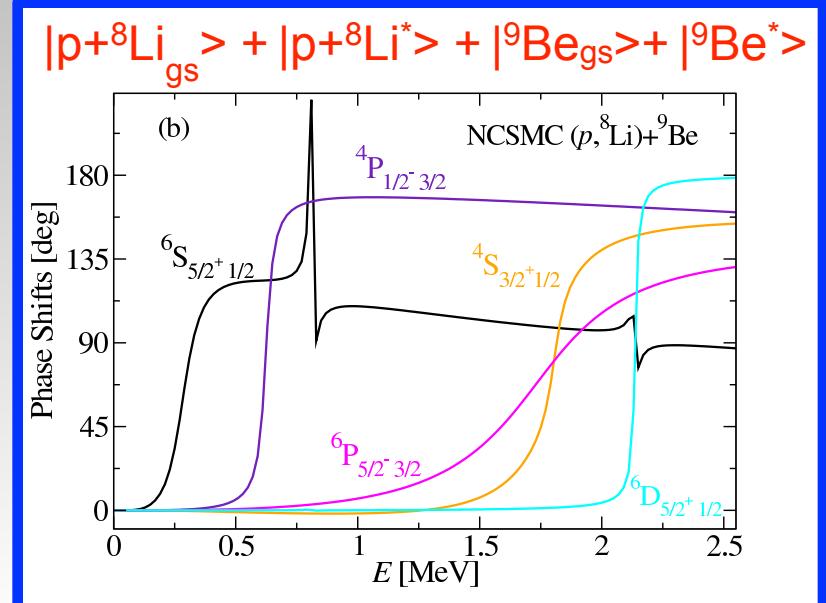
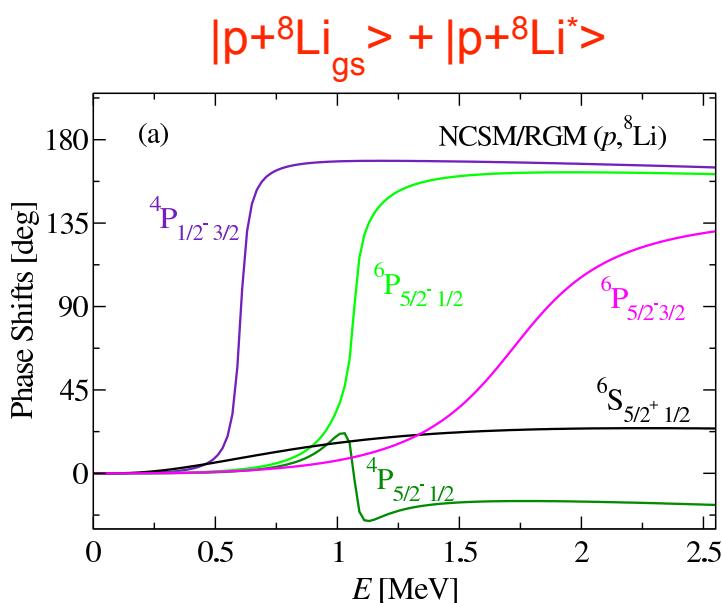


(p, ${}^8\text{Li}$) UNcoupled NCSM-RGM calculation Phase shifts

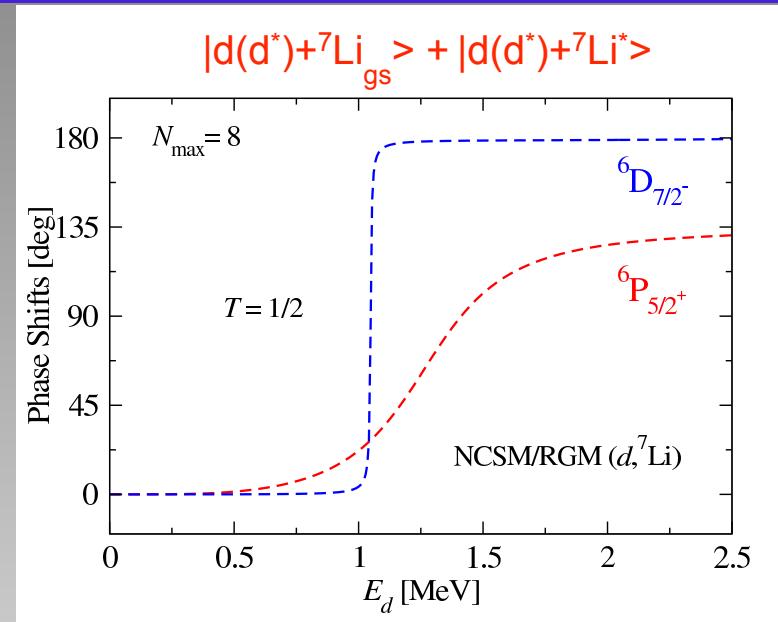
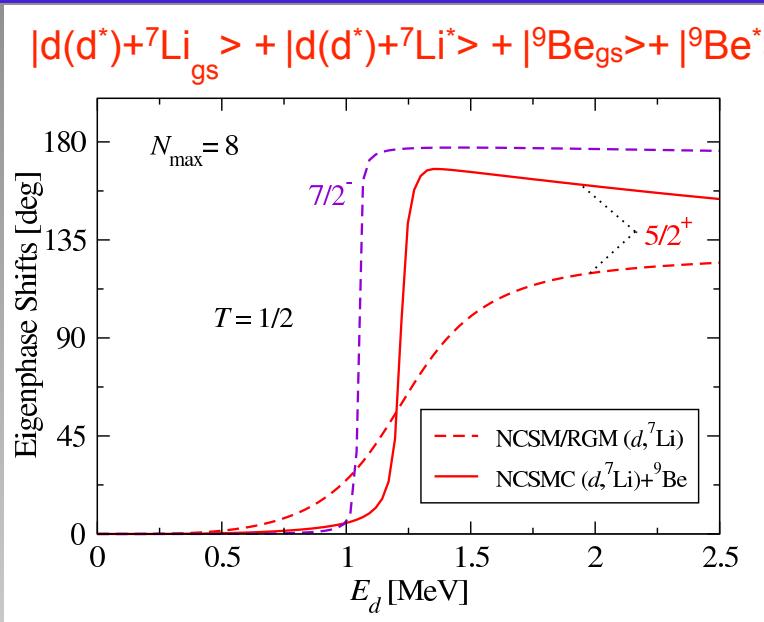


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Model space ($N_{\max}=8$ $\hbar\Omega=20$ MeV):

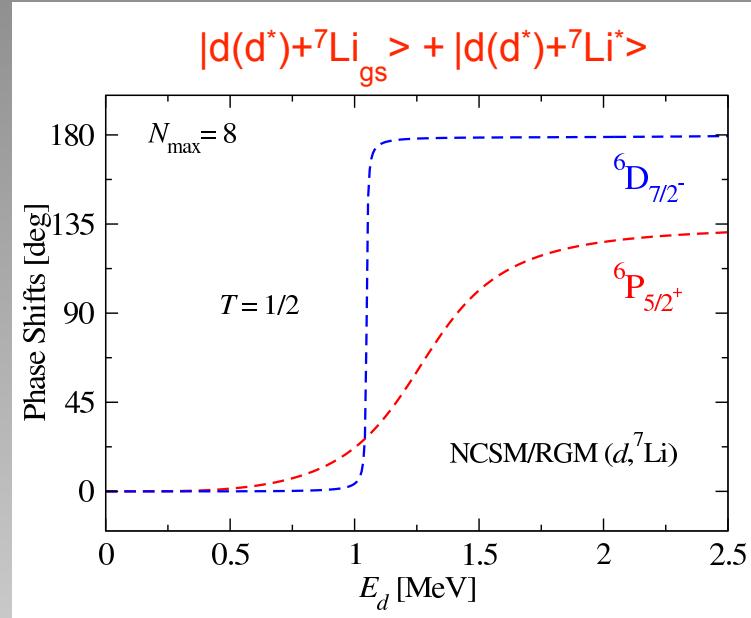
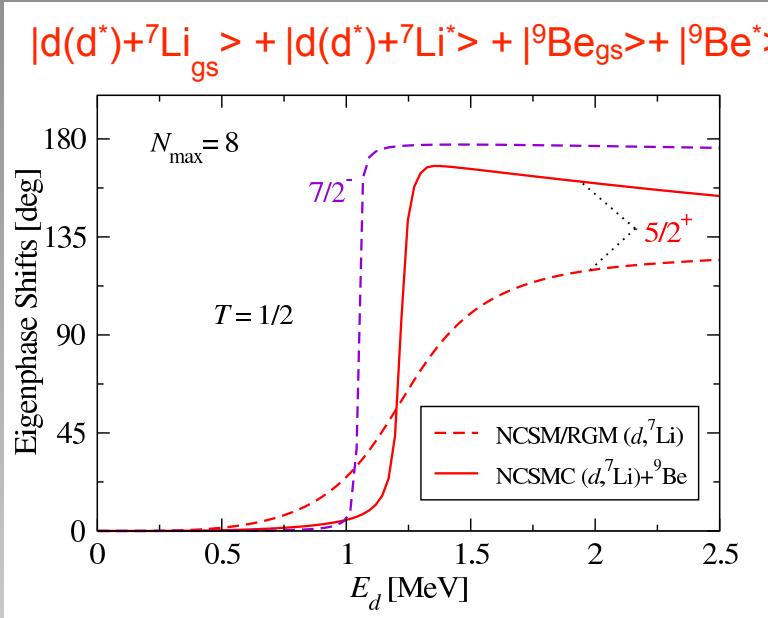


$(d, {}^7\text{Li})$ UNcoupled NCSM-RGM and NCSMC calculation

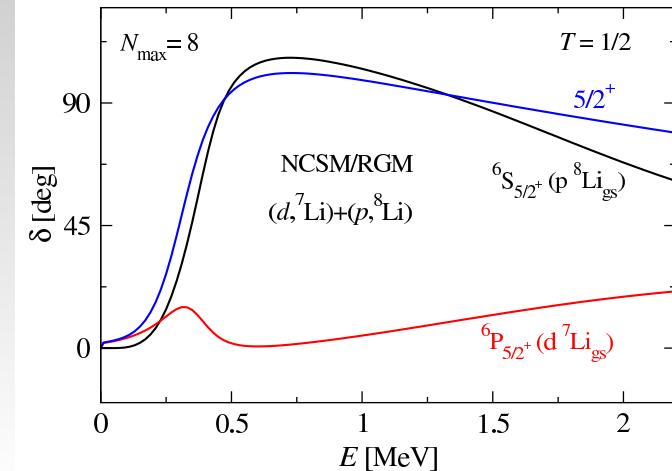


- Dominant resonances:
 $J=7/2^-$ in D-wave and $5/2^+$ in P-wave.
- Effect of the short-range correlations on $J=5/2^+$ in NCSMC calculation:
Decreased width of the resonance.
- Coupling effect:
Quenching of ${}^6\text{P}_{5/2^+}$ resonance

$(d, {}^7\text{Li})$ UNcoupled NCSM-RGM and NCSMC calculation

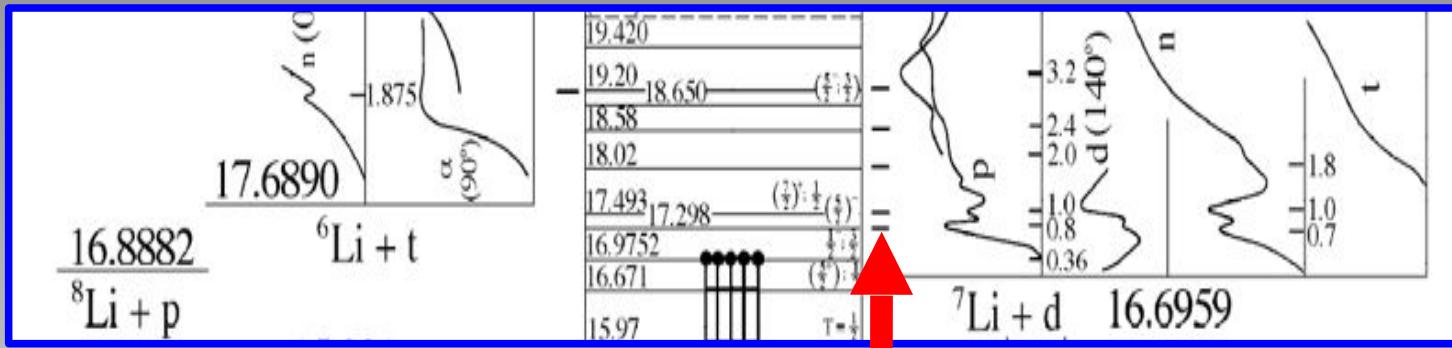


- Dominant resonances:
 $J=7/2^-$ in D-wave and $5/2^+$ in P-wave.
- Effect of the short-range correlations on $J=5/2^+$ in NCSMC calculation:
Decreased width of the resonance.
- Coupling effect:
Quenching of $6P_{5/2}^+$ resonance

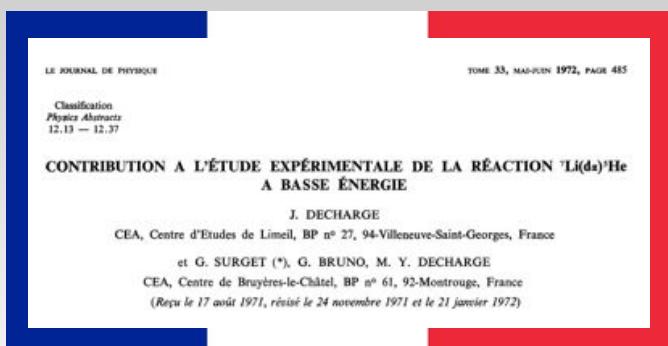


Spin-parity assignment of 0.78 MeV resonance of ${}^9\text{Be}$

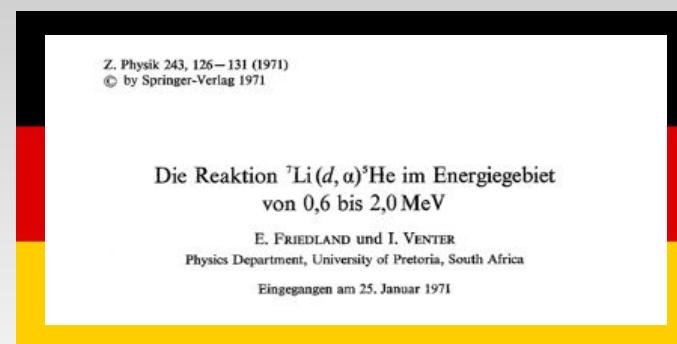
${}^9\text{Be}$ spectrum above d - ${}^7\text{Li}$ threshold



Low peak in the experimental total cross section:
 $E(5/2^-) \sim 0.78 \text{ MeV}$ above the threshold
(Uncertain spin-parity assignment)

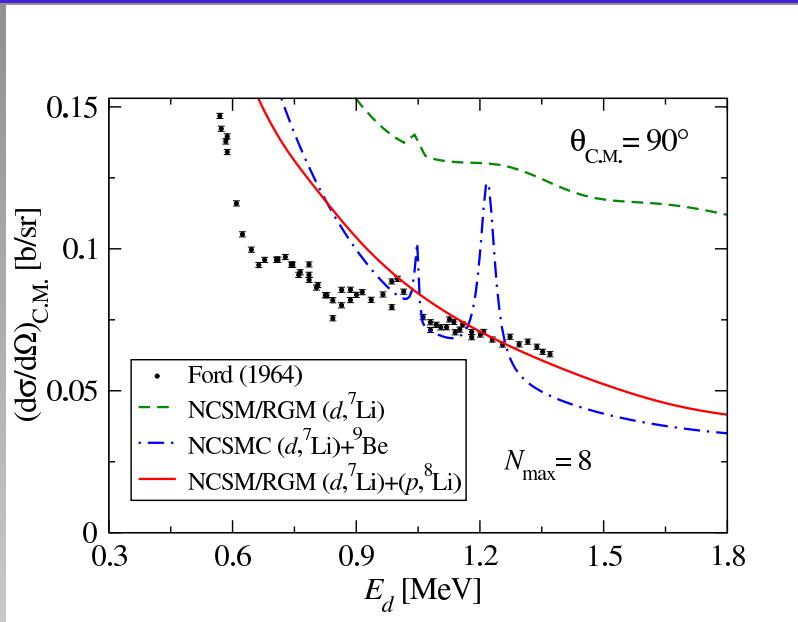


$E(\sim 17.3 \text{ MeV}) \rightarrow 3/2^-, 5/2^-$



$E(\sim 17.3 \text{ MeV}) \rightarrow 3/2^+, 5/2^+, 7/2^+$

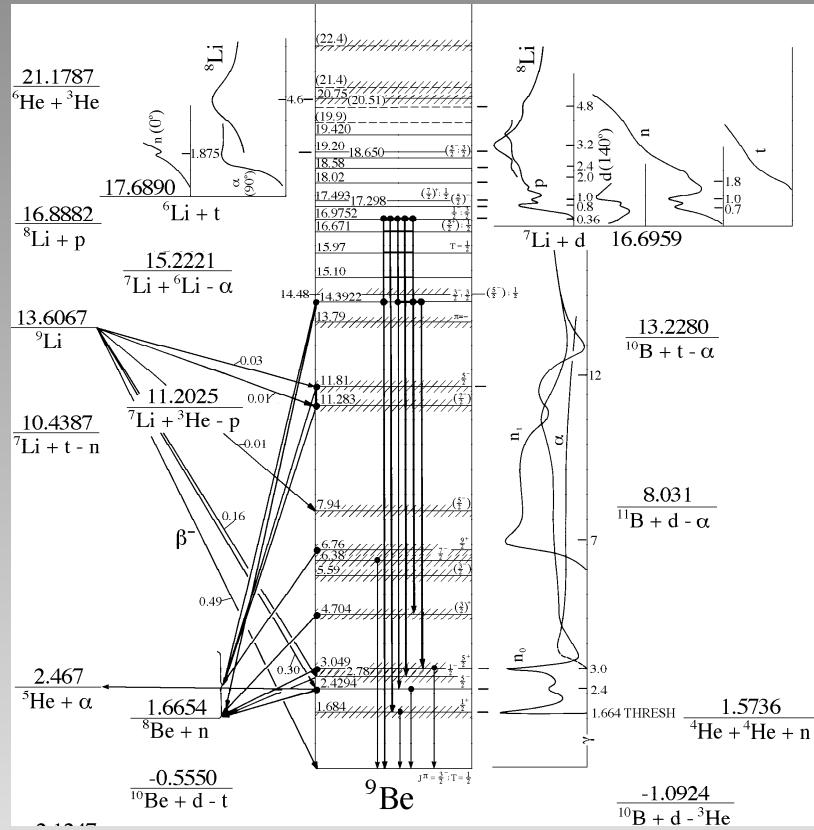
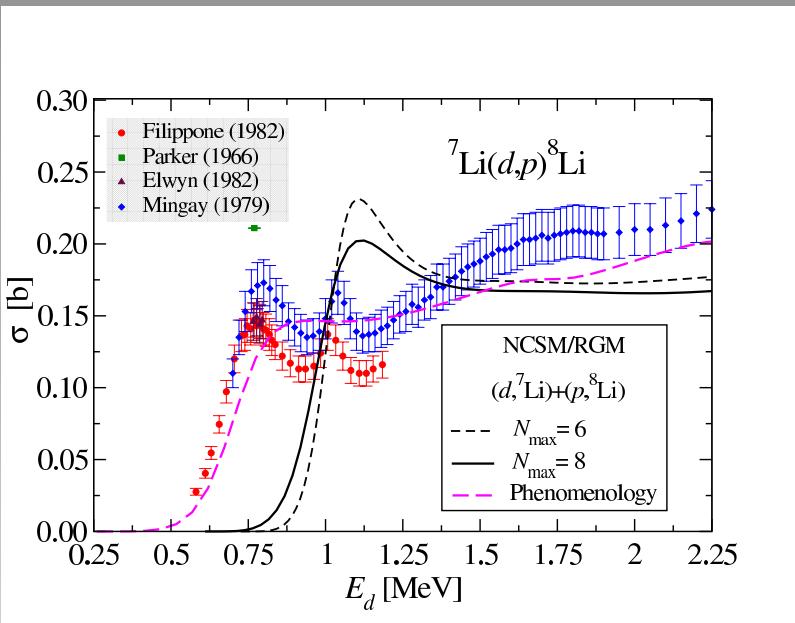
${}^7\text{Li}(d,d){}^7\text{Li}$ cross section



Experimental resonant peaks
at 0.8 MeV (S-wave)
and 1.0 MeV (P-wave)
(‘elastic’ process not ideal probe
for the ${}^9\text{Be}$ resonant states)

- Peak structure (1 MeV and 1.2 MeV) in uncoupled calculations ($J=7/2^-$ in D-wave and $5/2^+$ in P-wave). ...shifted at higher energy (missing bare 3N? SRG parameter dependence?)
- Effect of the short-range correlations in NCSMC calculation:
Increased lifetime of the resonance too narrow peaks (lack of p- ${}^8\text{Li}$ decay channel. other mass partition?)
- Qualitative trend of the data reproduced by NCSMC and coupled NCSM-RGM calculations ...still not-converged calculation at $N_{\max}=8$

$^7\text{Li}(d,p)^8\text{Li}$ total cross section



Included channels:

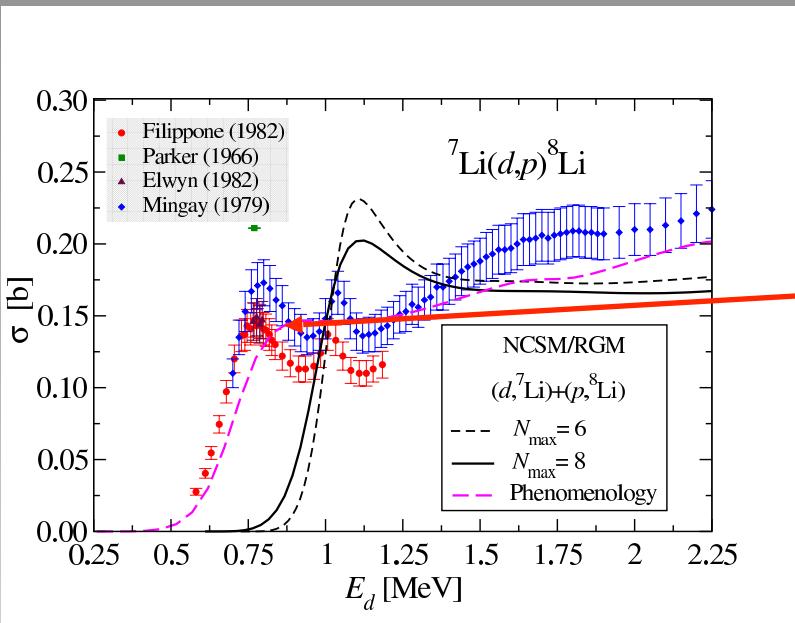
- (1) $p, ^8\text{Li}$
- (2) $d, ^7\text{Li}$
- (3) coupling (d,p)
- (4) virtual breakup of d

Not-included channels:

- (1) $^8\text{Be}, n$
- (2) $^6\text{Li}, t$

Channel	Exp. thresh. [MeV]	Theo.thresh. [MeV]
$d, ^7\text{Li}$	-41.470	-40.124
$p, ^8\text{Li}$	-41.278	-39.659
Q-value	-0.193	-0.465

${}^7\text{Li}(d,p){}^8\text{Li}$ total cross section

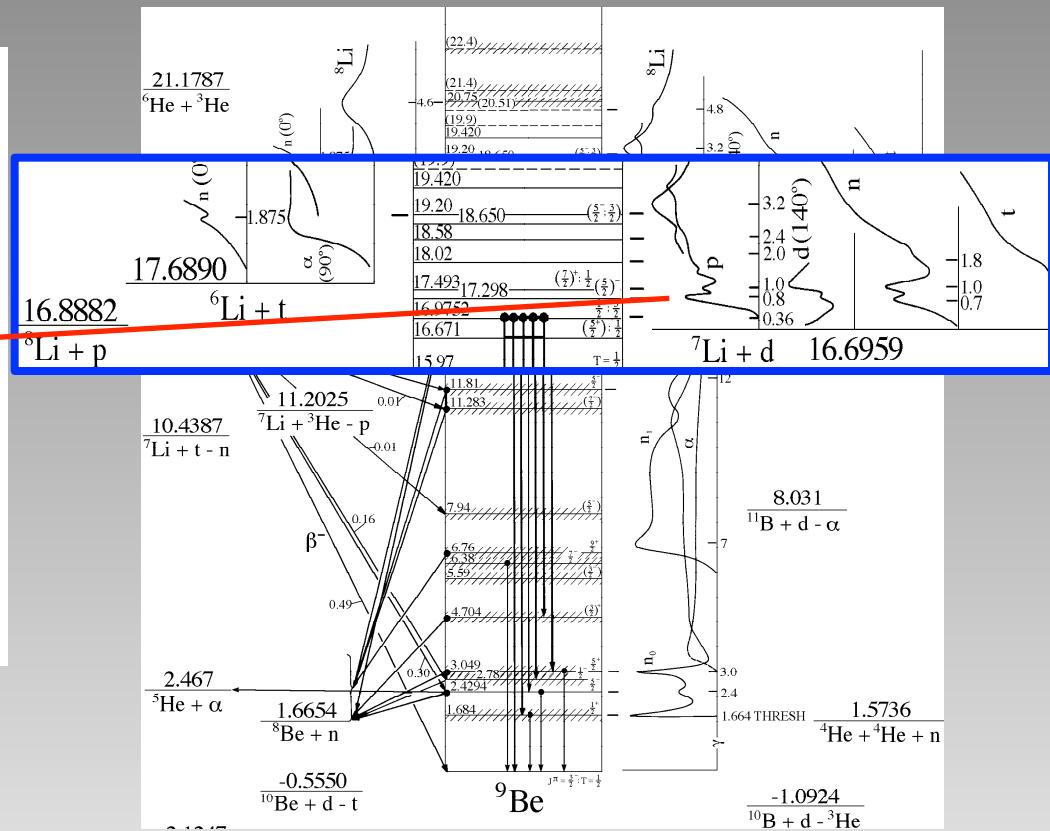


Included channels:

- (1) $p, {}^8\text{Li}$ (2) $d, {}^7\text{Li}$ (3) coupling (d,p)
- (4) virtual breakup of d

Not-included channels:

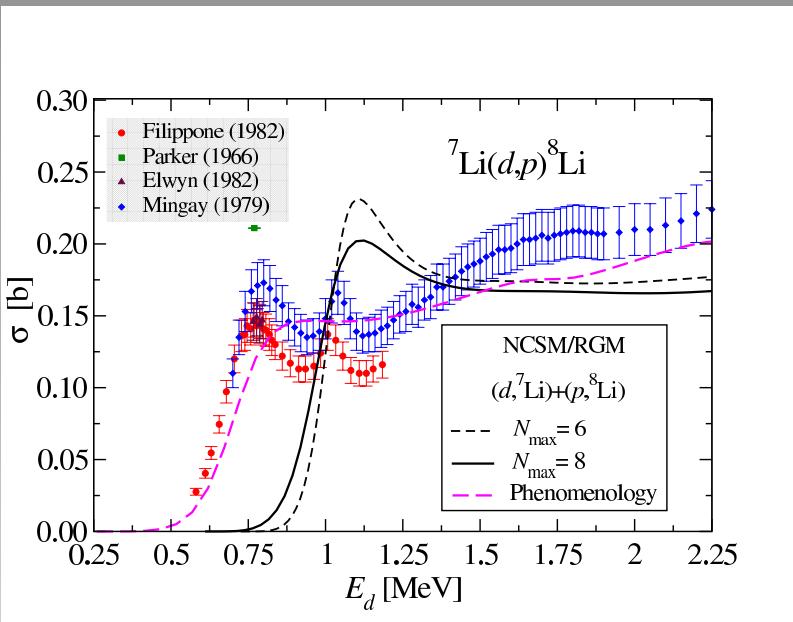
- (1) ${}^8\text{Be}, n$ (2) ${}^6\text{Li}, t$



Experimental recommended value
 $0.147 \pm 0.011 \text{ b } (\Gamma \approx 0.2 \text{ MeV})$ at 0.78 MeV
 of deuteron kinetic energy

Calibration peak for
 ${}^7\text{Be}(p,\gamma){}^8\text{B}$ radiative capture reaction

${}^7\text{Li}(d,p){}^8\text{Li}$ total cross section



- Position of the first resonant peak overestimated by ~ 0.33 MeV (see Q-value)
- Peak at 17.493 MeV (${}^9\text{Be}$ spectrum) not reproduced (missing ${}^8\text{Be}(\alpha,\alpha)\text{n}$? 3N forces?)

$d, {}^7\text{Li}, {}^8\text{Li}$ NCSM energies adjusted to reproduce the experimental Q-value of the reaction

Included channels:

- (1) $p, {}^8\text{Li}$
- (2) $d, {}^7\text{Li}$
- (3) coupling (d,p)
- (4) virtual breakup of d

position of first peak
slightly overestimated

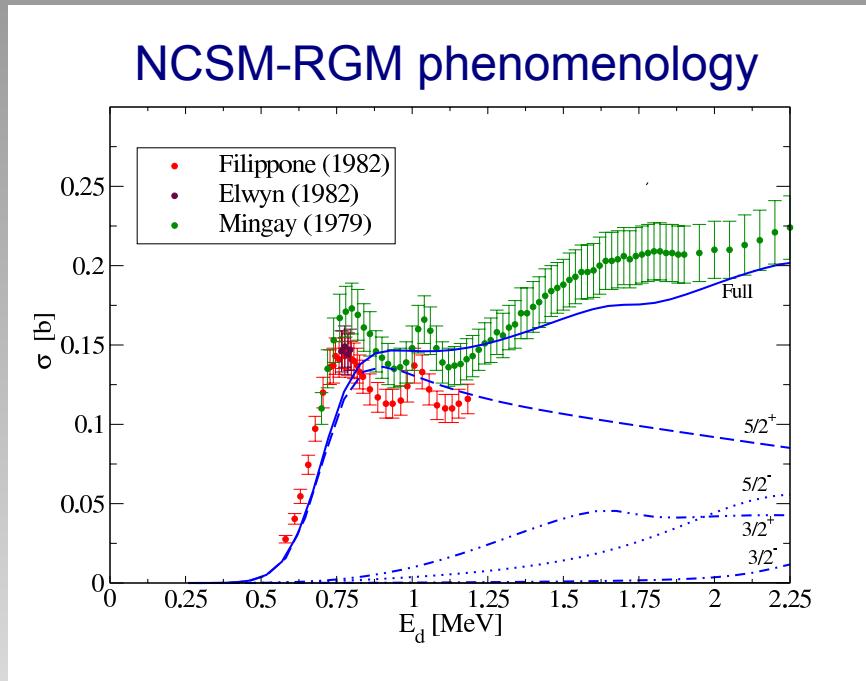
Not-included channels:

- (1) ${}^8\text{Be}, \text{n}$
- (2) ${}^6\text{Li}, \text{t}$

Channel	Exp. thresh. [MeV]	Theo.thresh. [MeV]
$d, {}^7\text{Li}$	-41.470	-40.124
$p, {}^8\text{Li}$	-41.278	-39.659
Q-value	-0.193	-0.465

Impact of different partial waves on NCSM-RGM total cross section

${}^7\text{Li}(d,p){}^8\text{Li}$ cross section



- Confirmed dominant role played by $5/2^+$ partial wave
- Below ~ 2 MeV the cross section is dominated by positive-parity partial waves
- Increasing trend up to deuteron break-up fairly well reproduced (contribution from $5/2^-$ and $3/2^+$ partial wave)

Conclusions & Perspectives

First application of the NCSM-RGM for deuteron-projectile and p-shell nucleus as target:

- Inclusion of the “elastic” and coupling channel in the description of transfer reactions

Study of the ${}^7\text{Li}(d,p){}^8\text{Li}$ transfer reaction and of the ${}^9\text{Be}$ resonances above d - ${}^7\text{Li}$ threshold:

- Discussion of the spin-parity assignment of $E_d=0.78$ MeV resonance

To be done:

- Complete the calculation of the ${}^7\text{Li}(d,p){}^8\text{Li}$ transfer reaction in the NCSMC framework
- Include 3N force also for p-shell nuclei