



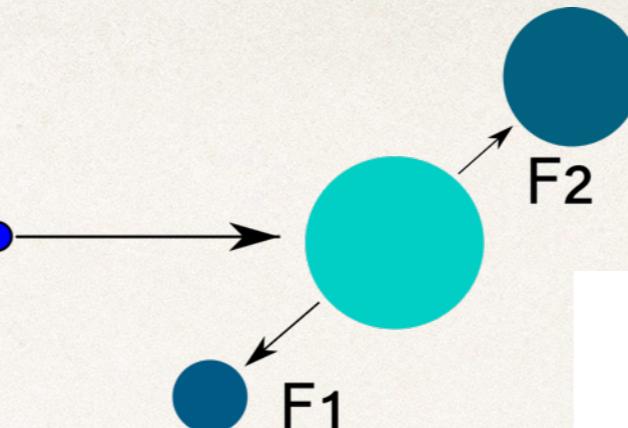
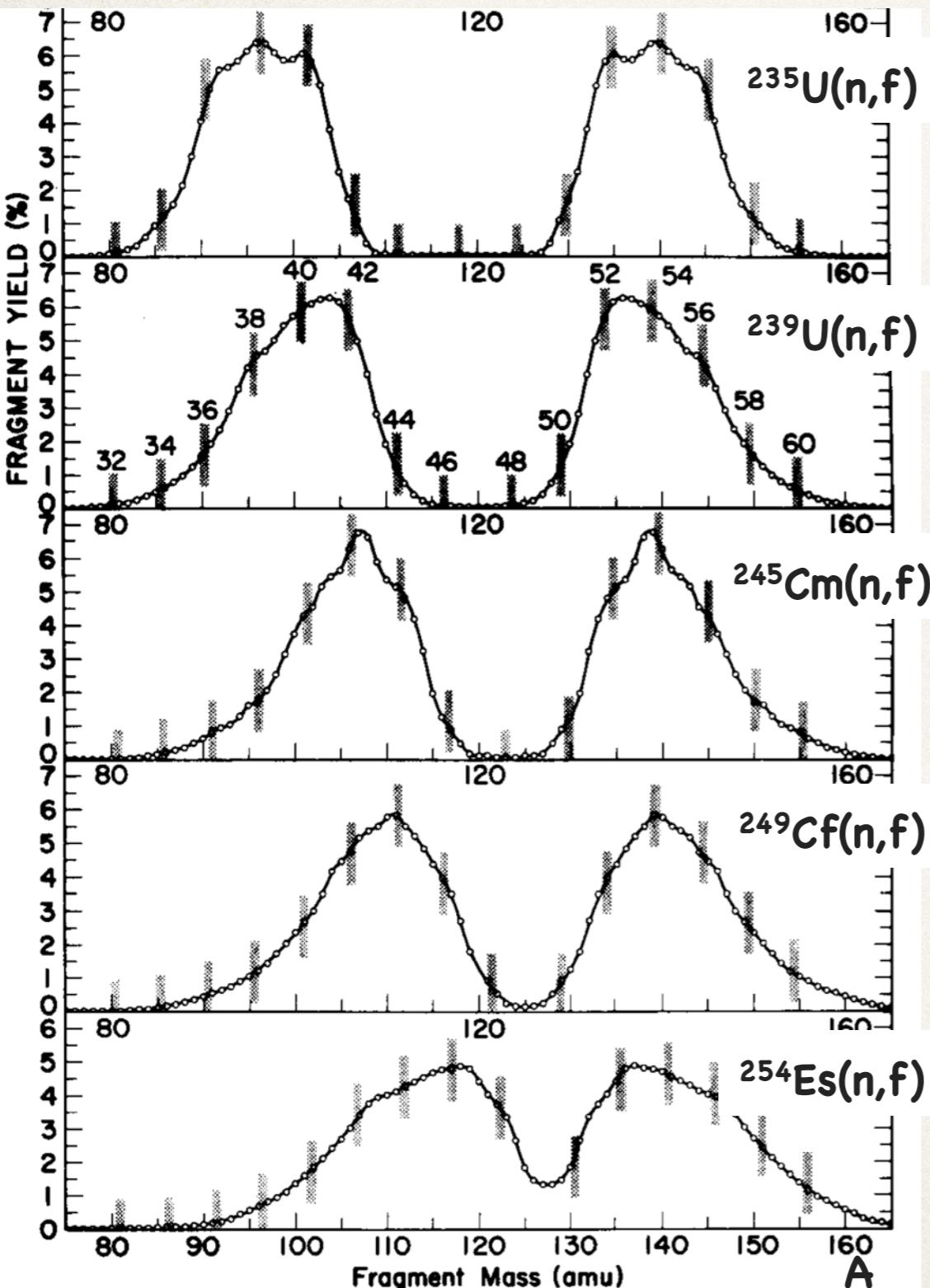
Dependence of Fission-Fragment Properties on Excitation Energy for N-rich Actinides

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Limitations of Direct Kinematics

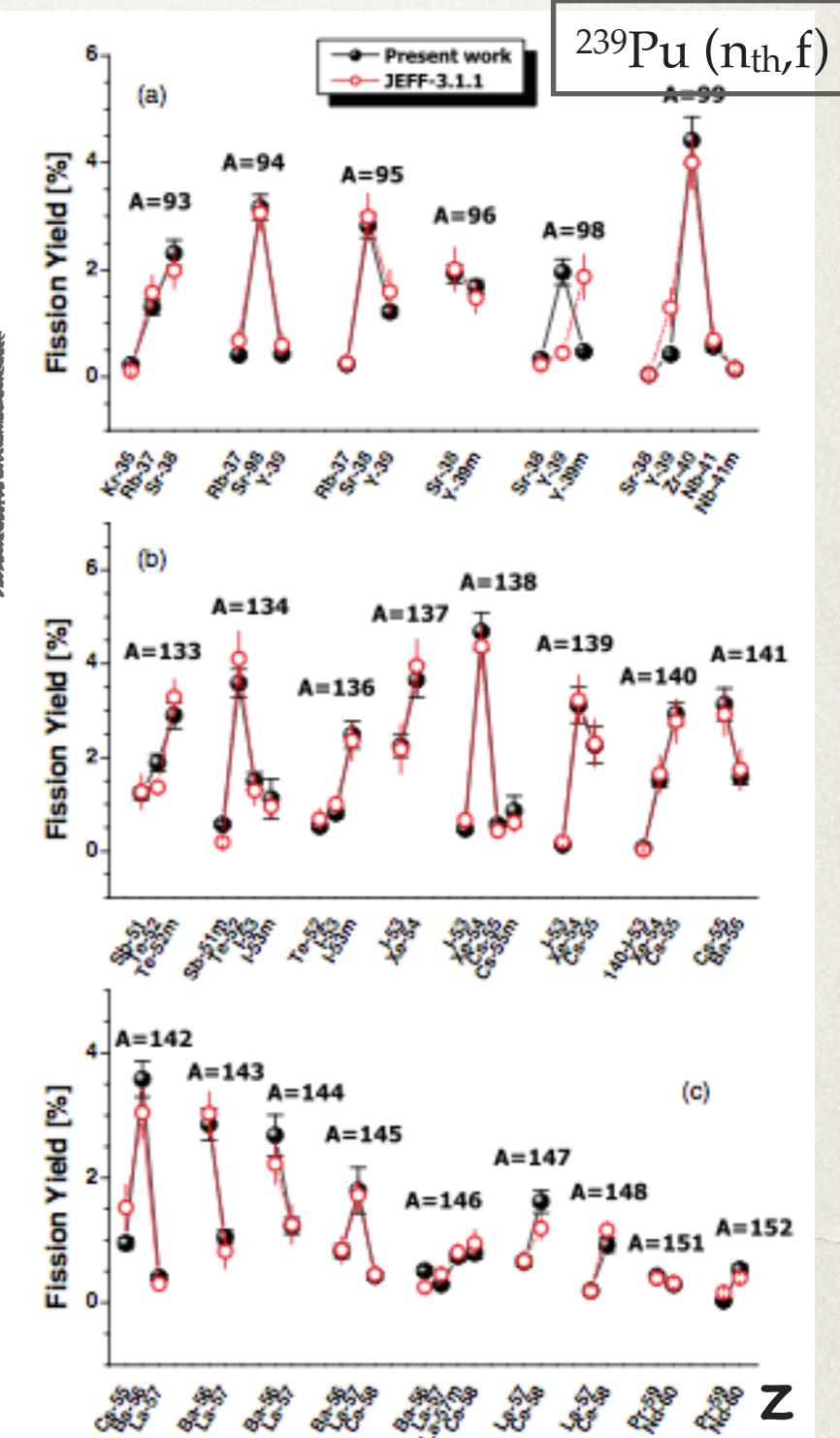
Good measurement of mass distribution



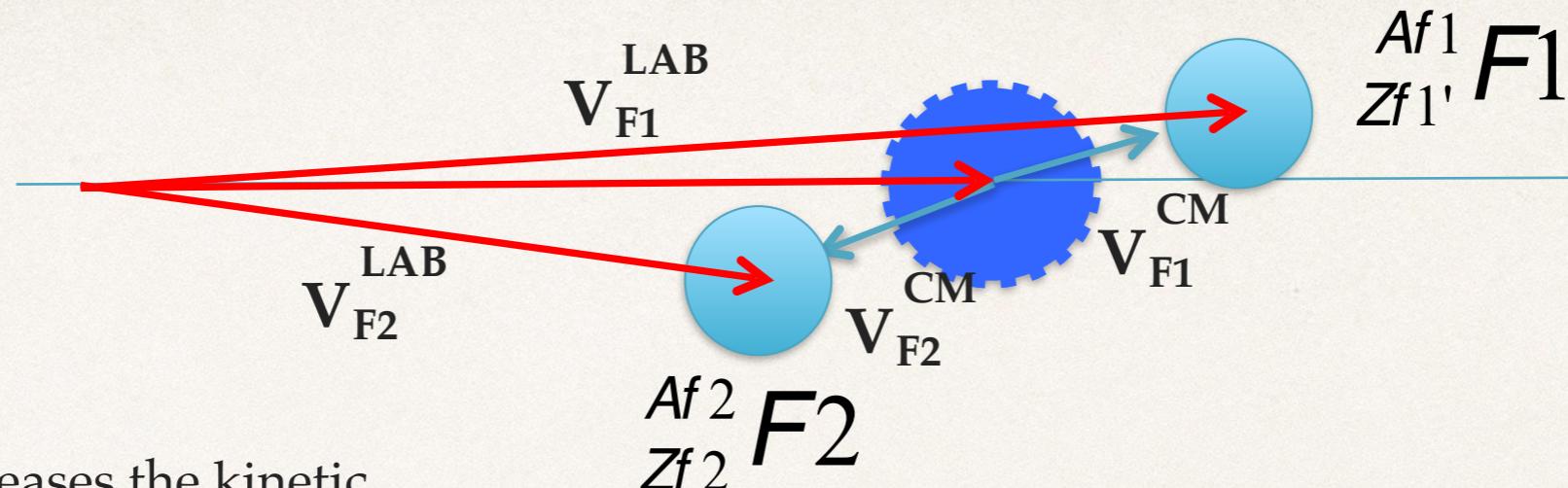
Incomplete Measurements of Isotopic distributions

limitations for heavy ff
- lifetimes
-unknown level schemes

Isotopic distribution by γ -spectroscopy

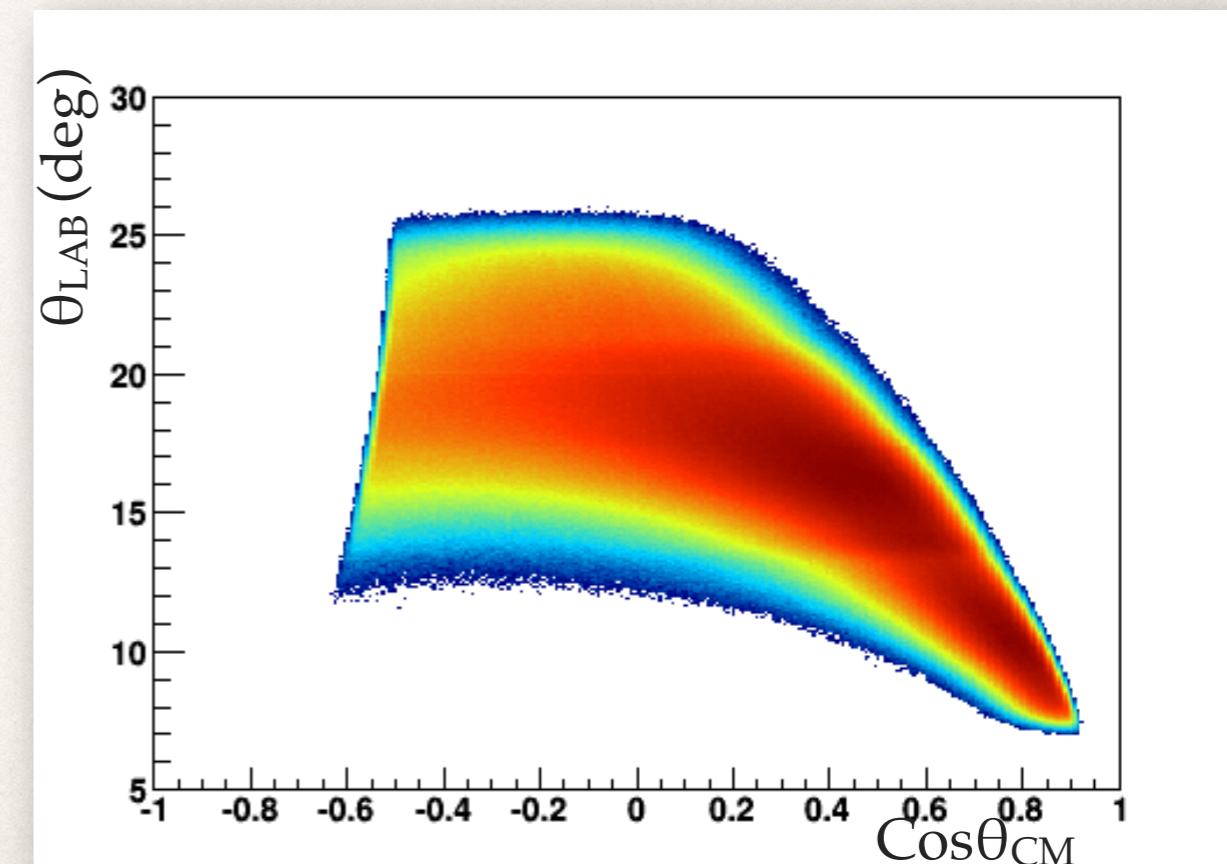
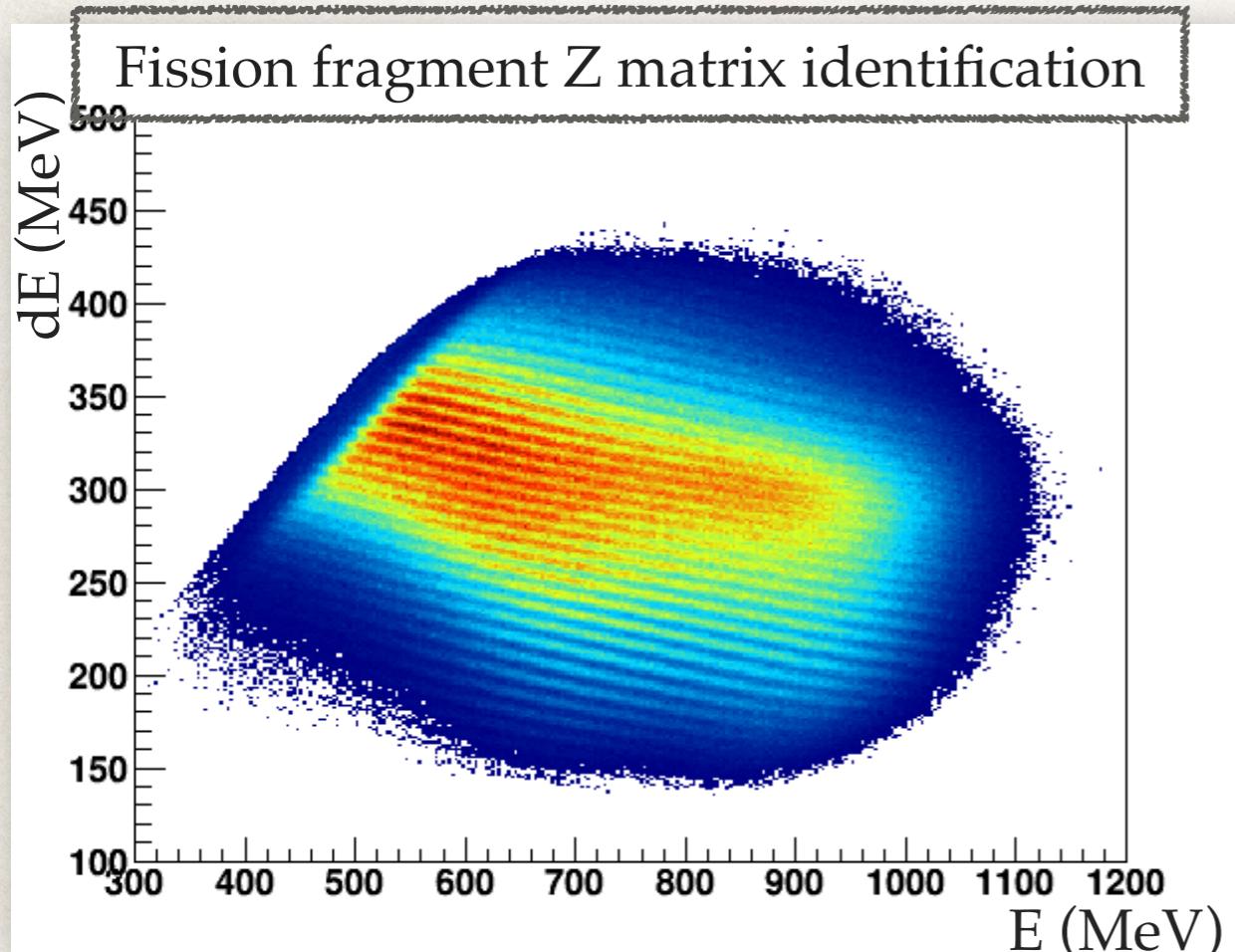


Goals of Inverse kinematics



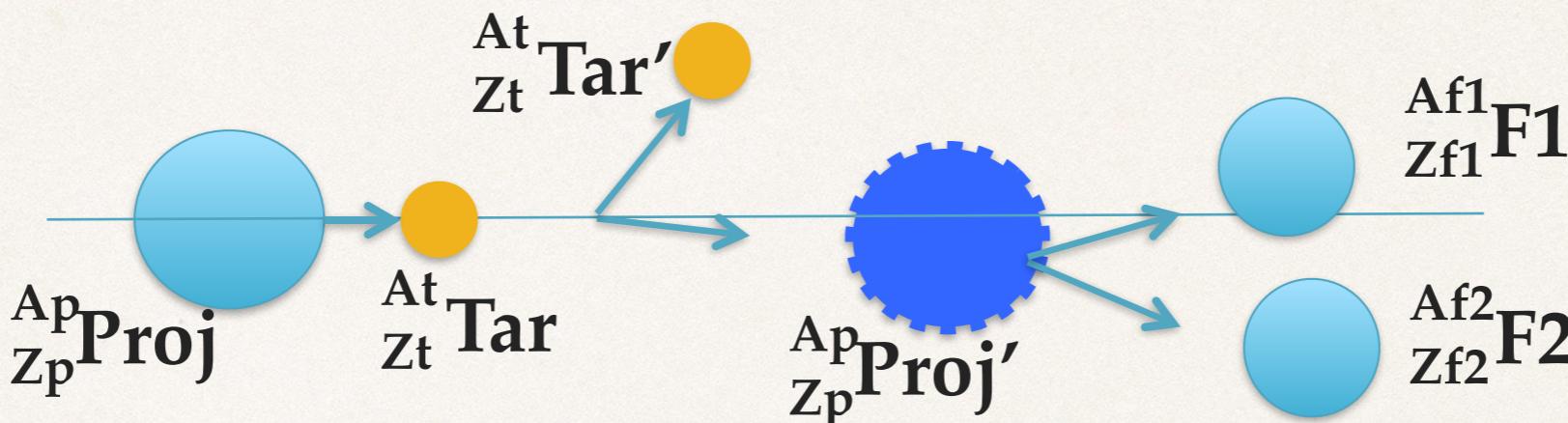
Kinematical boost increases the kinetic energy of the fission fragments providing **the capability of a direct identification**

Kinematical boost allows to keep a wide angular coverage in the CM frame when the size of the detectors is limited



Reaction Mechanism

$^{238}\text{U} + ^{12}\text{C}$ @ 6.14 AMeV



Fissioning Systems

^{242}Cf	^{243}Cf	^{244}Cf	^{245}Cf	^{246}Cf	^{247}Cf	^{248}Cf	^{249}Cf	^{250}Cf	^{251}Cf	^{252}Cf
^{241}Bk	^{242}Bk	^{243}Bk	^{244}Bk	^{245}Bk	^{246}Bk	^{247}Bk	^{248}Bk	^{249}Bk	^{250}Bk	^{251}Bk
^{240}Cm	^{241}Cm	^{242}Cm	^{243}Cm	^{244}Cm	^{245}Cm	^{246}Cm	^{247}Cm	^{248}Cm	^{249}Cm	^{250}Cm
^{239}Am	^{240}Am	^{241}Am	^{242}Am	^{243}Am	^{244}Am	^{245}Am	^{246}Am	^{247}Am	^{248}Am	^{249}Am
^{238}Pu	^{239}Pu	^{240}Pu	^{241}Pu	^{242}Pu	^{243}Pu	^{244}Pu	^{245}Pu	^{246}Pu	^{247}Pu	
^{237}Np	^{238}Np	^{239}Np	^{240}Np	^{241}Np	^{242}Np	^{243}Np	^{244}Np			
^{236}U	^{237}U	^{238}U	^{239}U	^{240}U	^{241}U	^{242}U				

Fissioning systems not accessible from any other mechanism

10% above Coulomb barrier

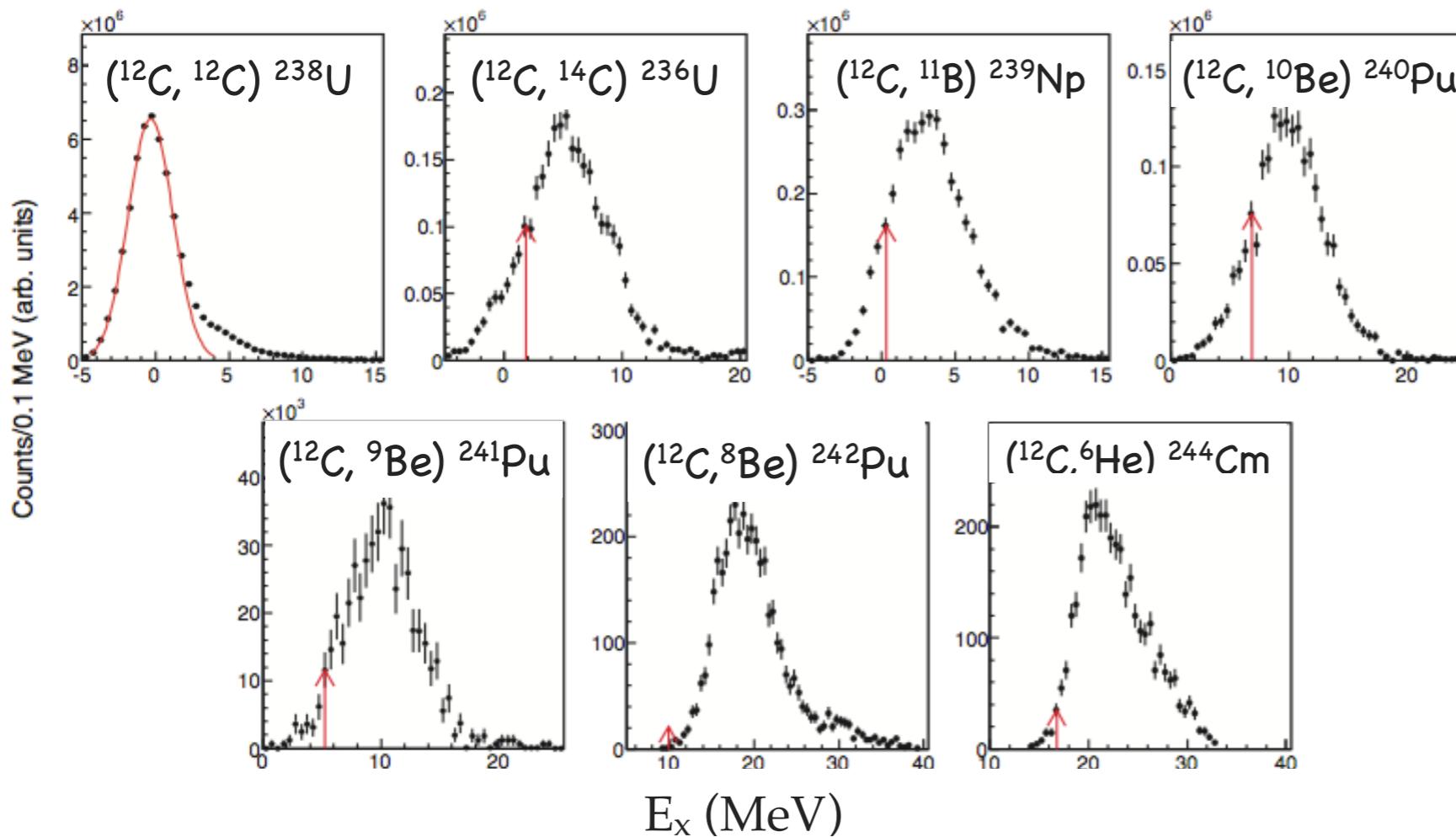
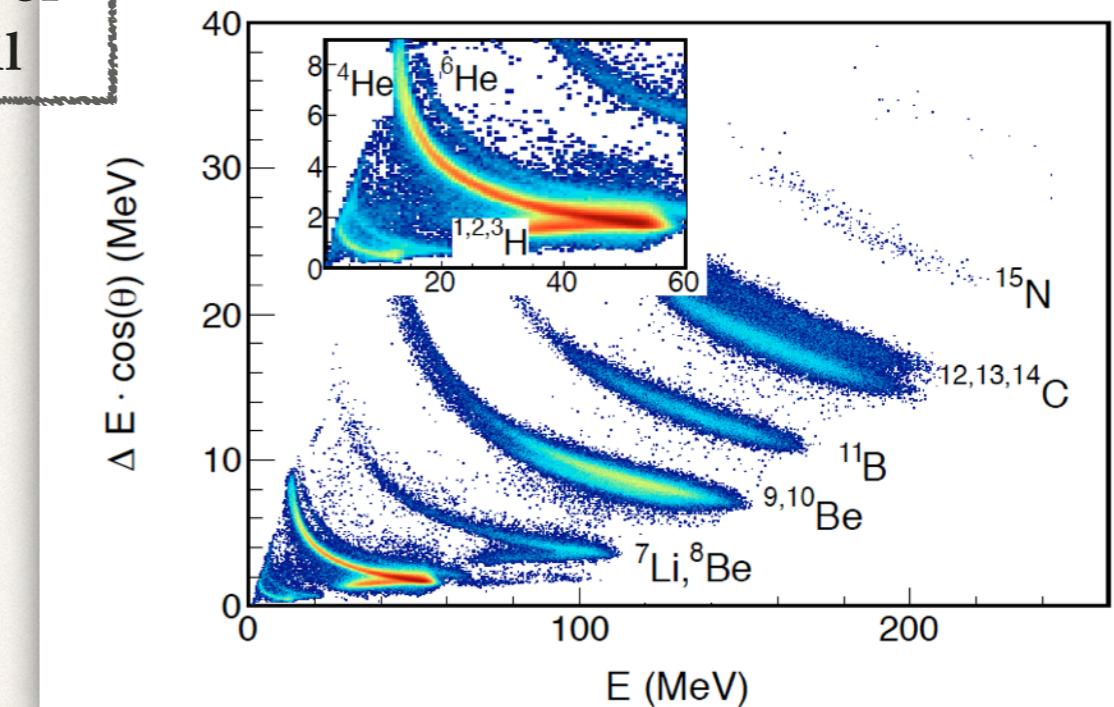
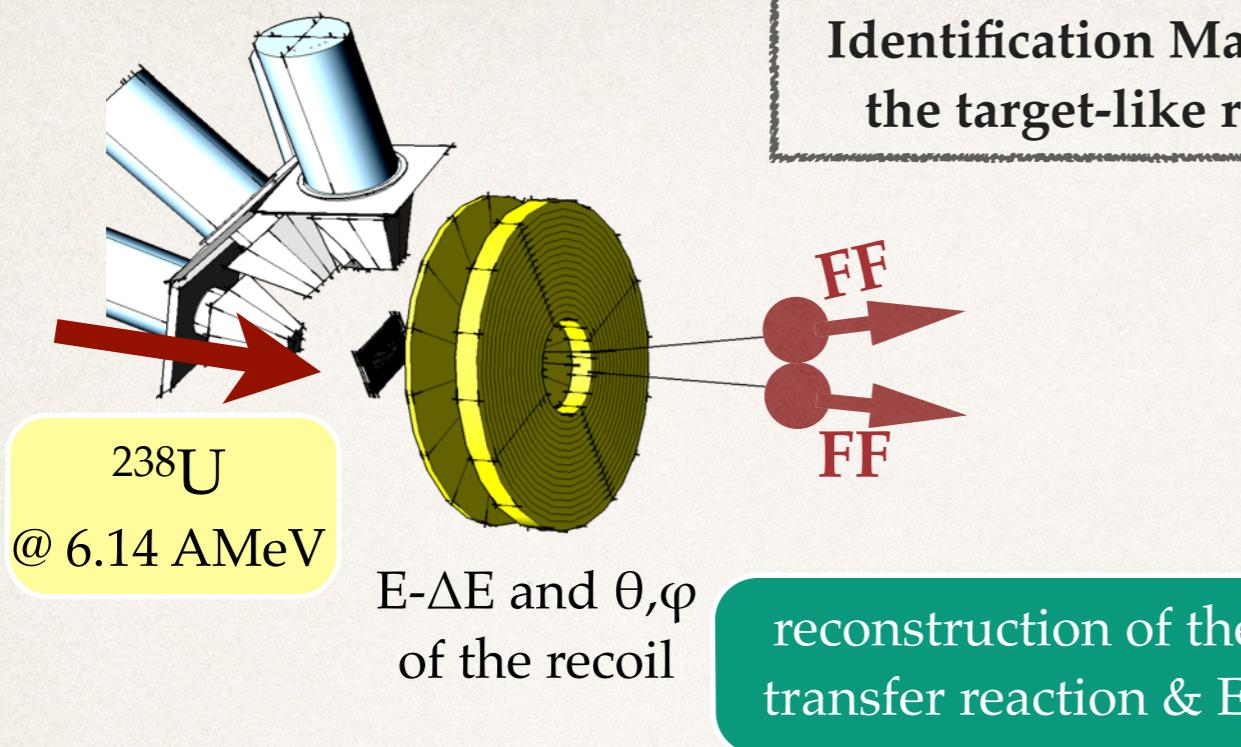
Transfer-Fission:

10 n-rich actinides produced with a distribution of E_x below 30 MeV

Fusion-Fission:

production of ^{250}Cf with $E_x = 46$ MeV
10 times more likely than any transfer channel

Transfer Reaction and Excitation Energy



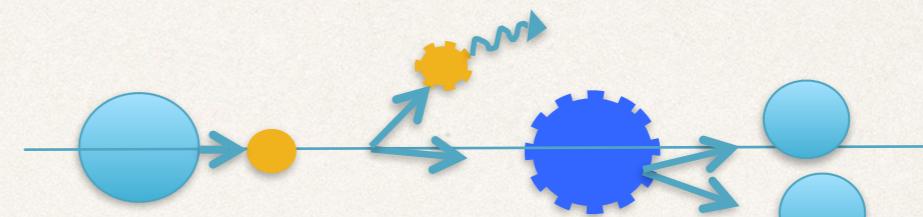
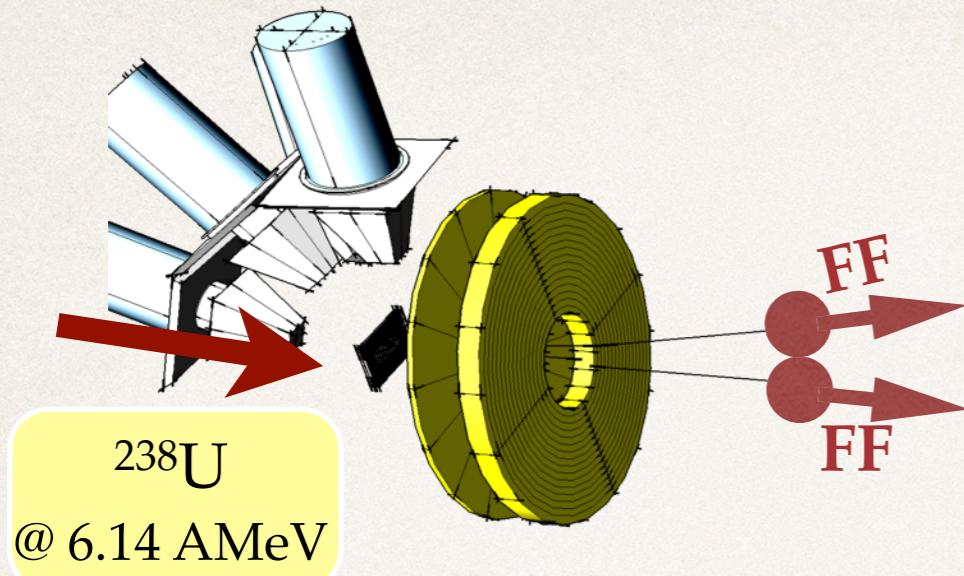
Distribution of E_x for the different fissioning systems from reconstruction of the binary reaction

Higher E_x for higher number of transferred nucleons

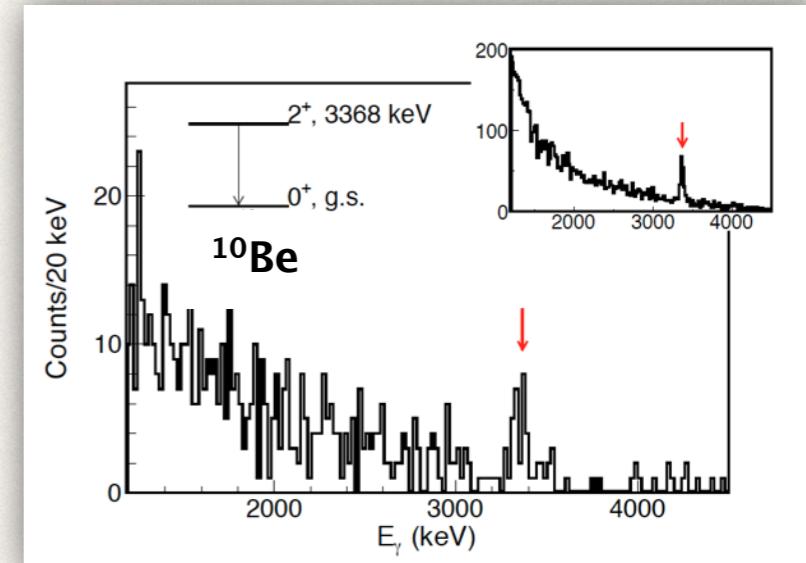
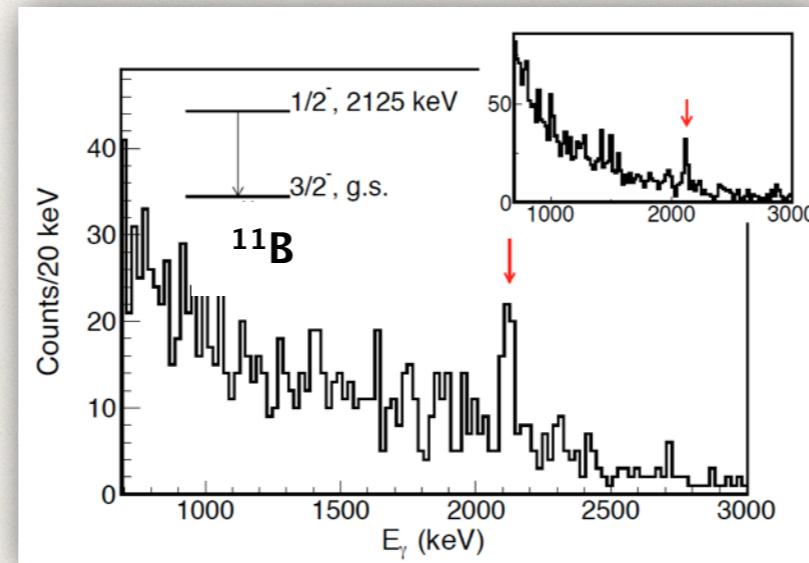
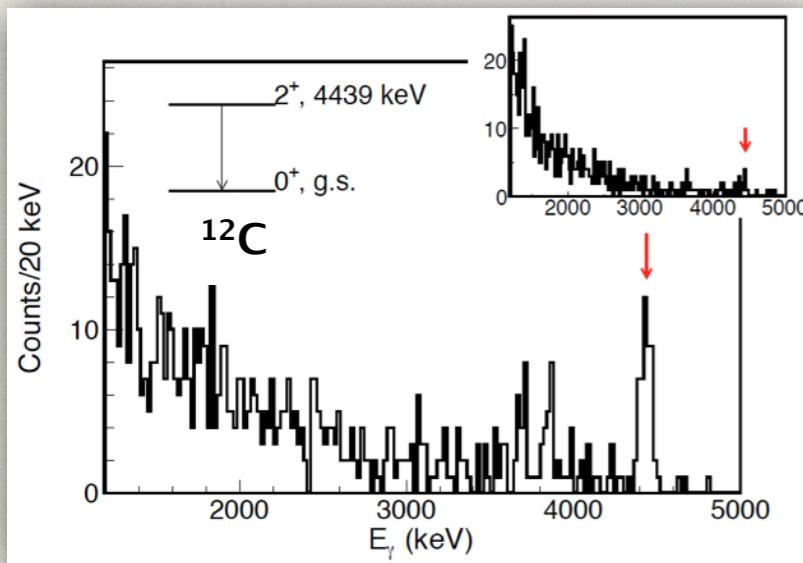
$E_x \sim 8$ MeV is comparable with fast-neutron fission

C. Rodriguez-Tajes et al.,
PRC (2014) 024614

Excitation of Target-like Recoil

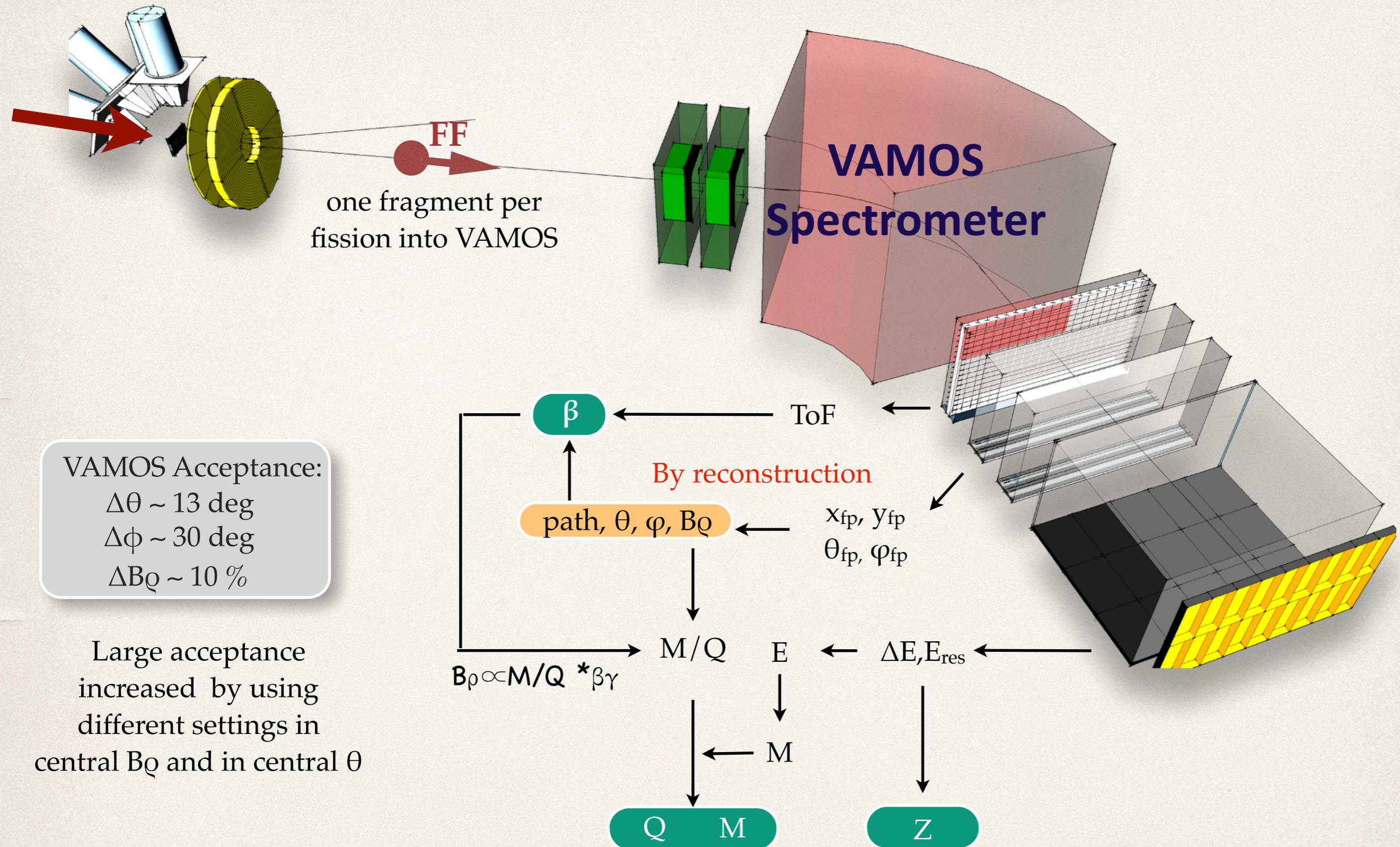


EXOGAM detector allow us to evaluate the excitation probability of the target-like nuclei



γ -rays measurements show excited states in ^{12}C , ^{11}B and ^{10}Be in coincidence with fission with $P_\gamma = 0.12-0.14$

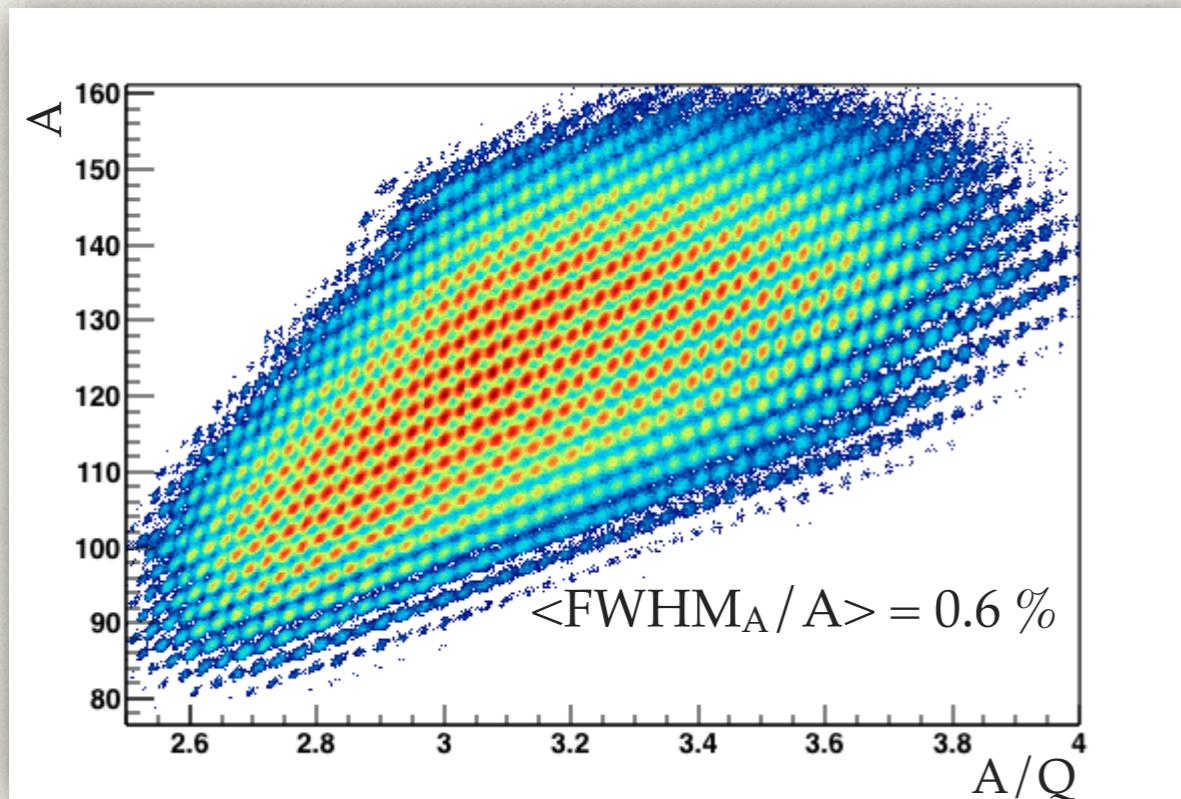
Fission Fragments Detection



M. Rejmund et al., NIM A 646 (2011) 184
 S. Pullanhotan et al., NIM A 593 (2008) 343

Fission Fragments Identification

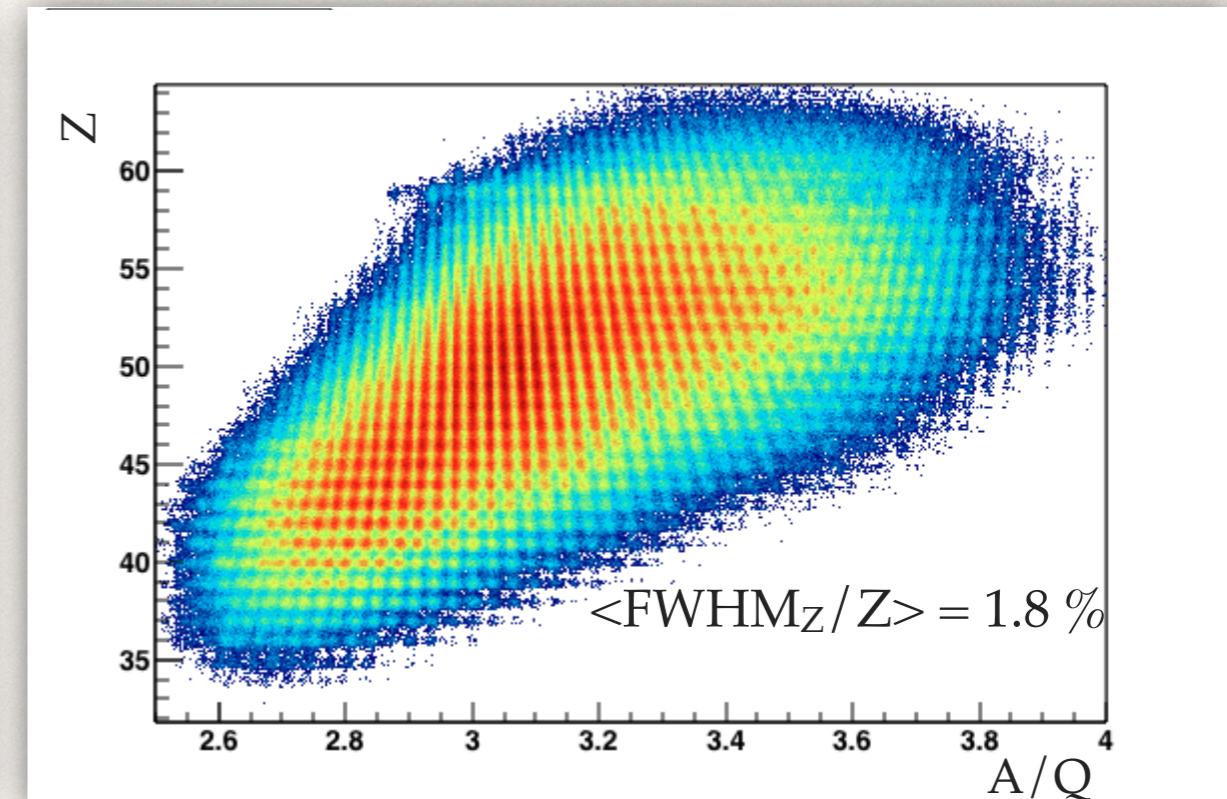
Mass Identification



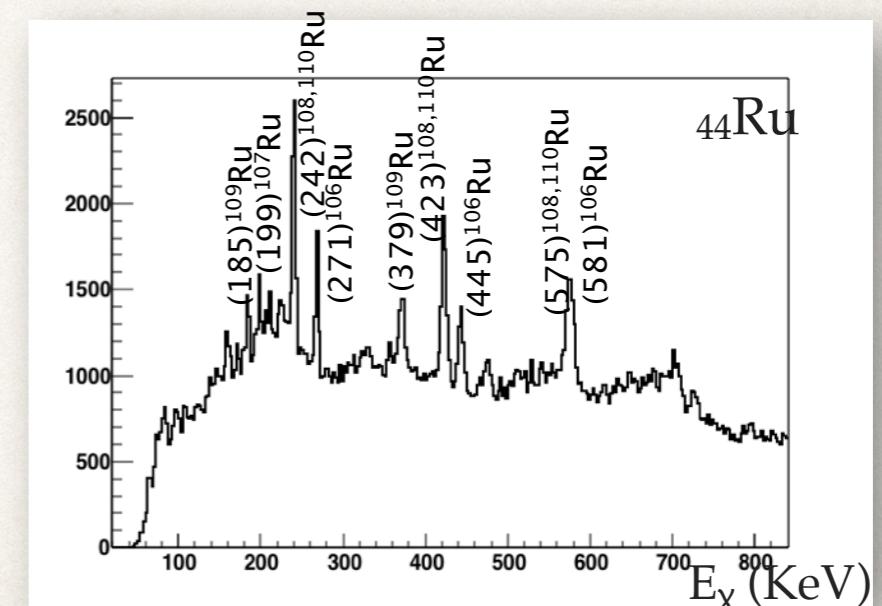
A/Q provides the Q separation and contributes to a better A resolution

More than 300 isotopes identified

Proton Number Identification



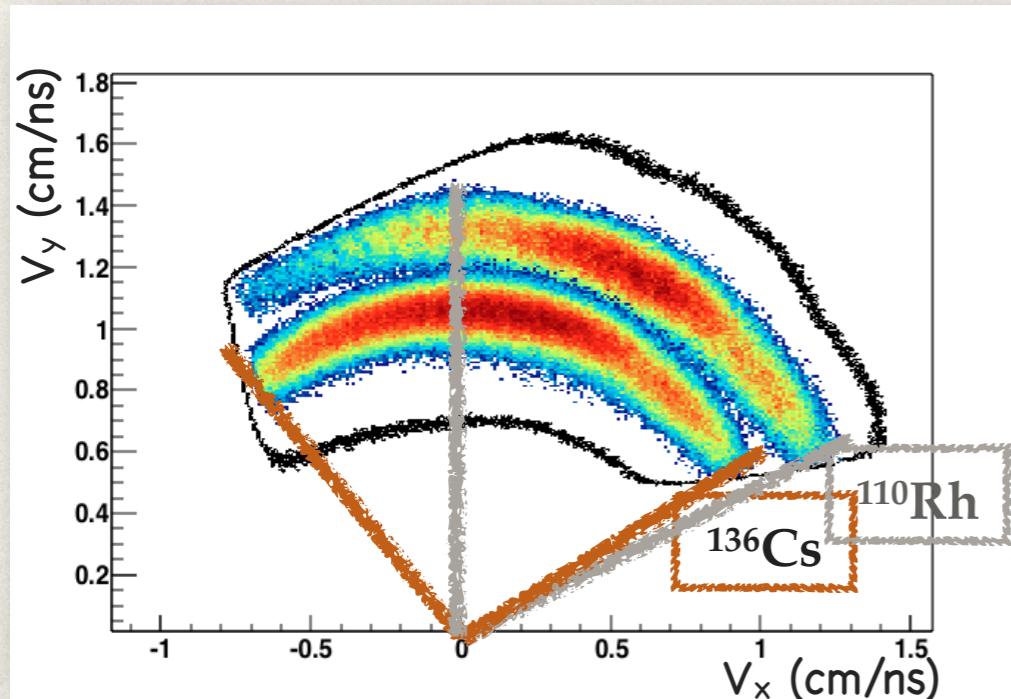
γ -rays in coincidence with fission fragments provide a cross check for the Z and A identification



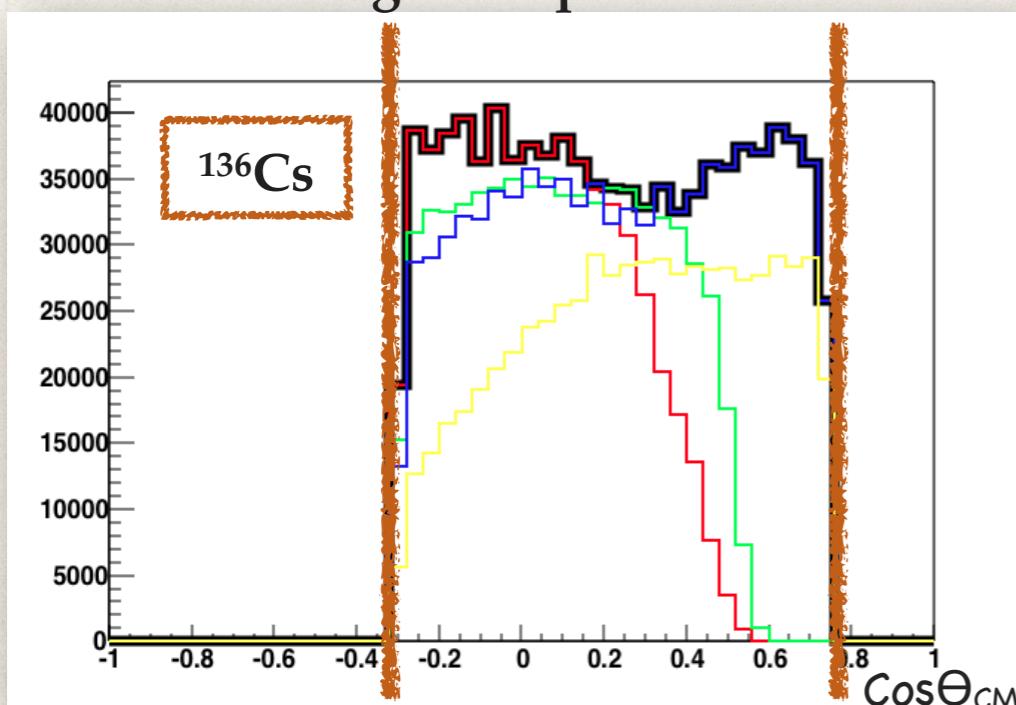
A. Shrivastava et al. PRC 80 (2009) 051305

Transmission through VAMOS

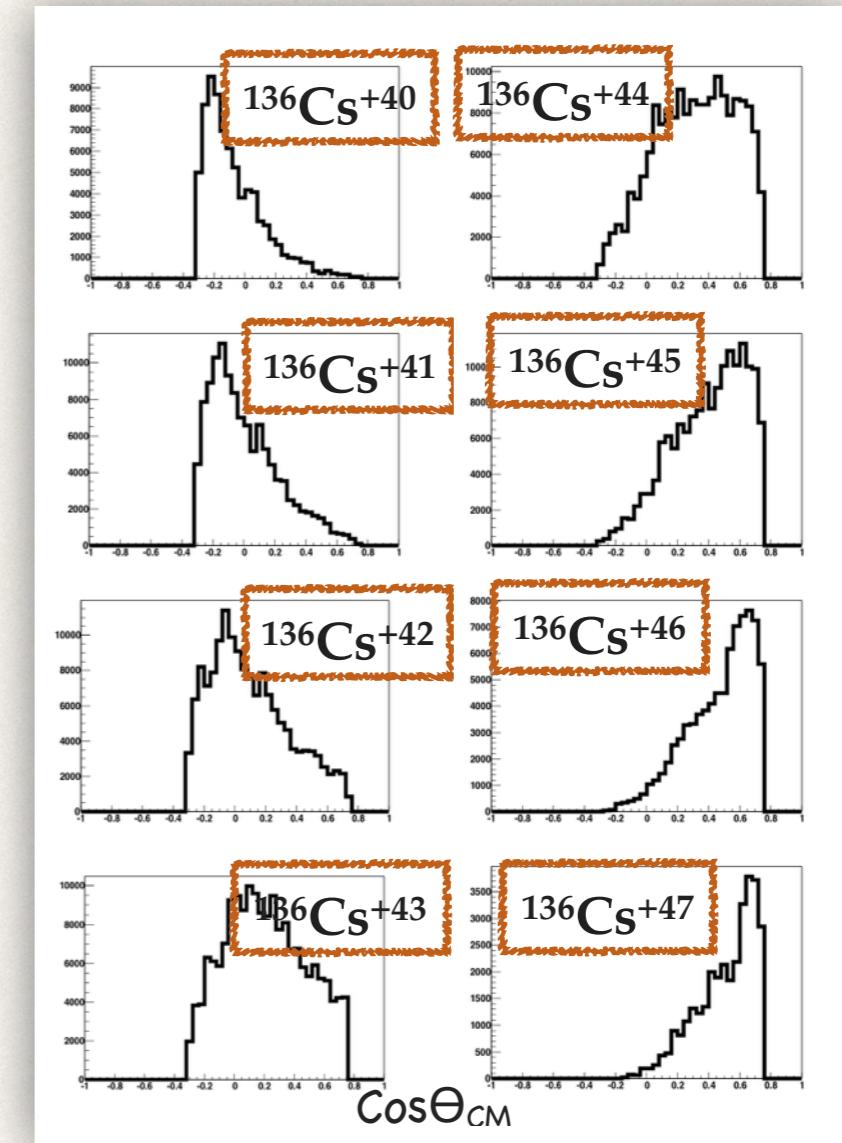
The detection is limited by the transmission



Beam normalization for different settings is required as well



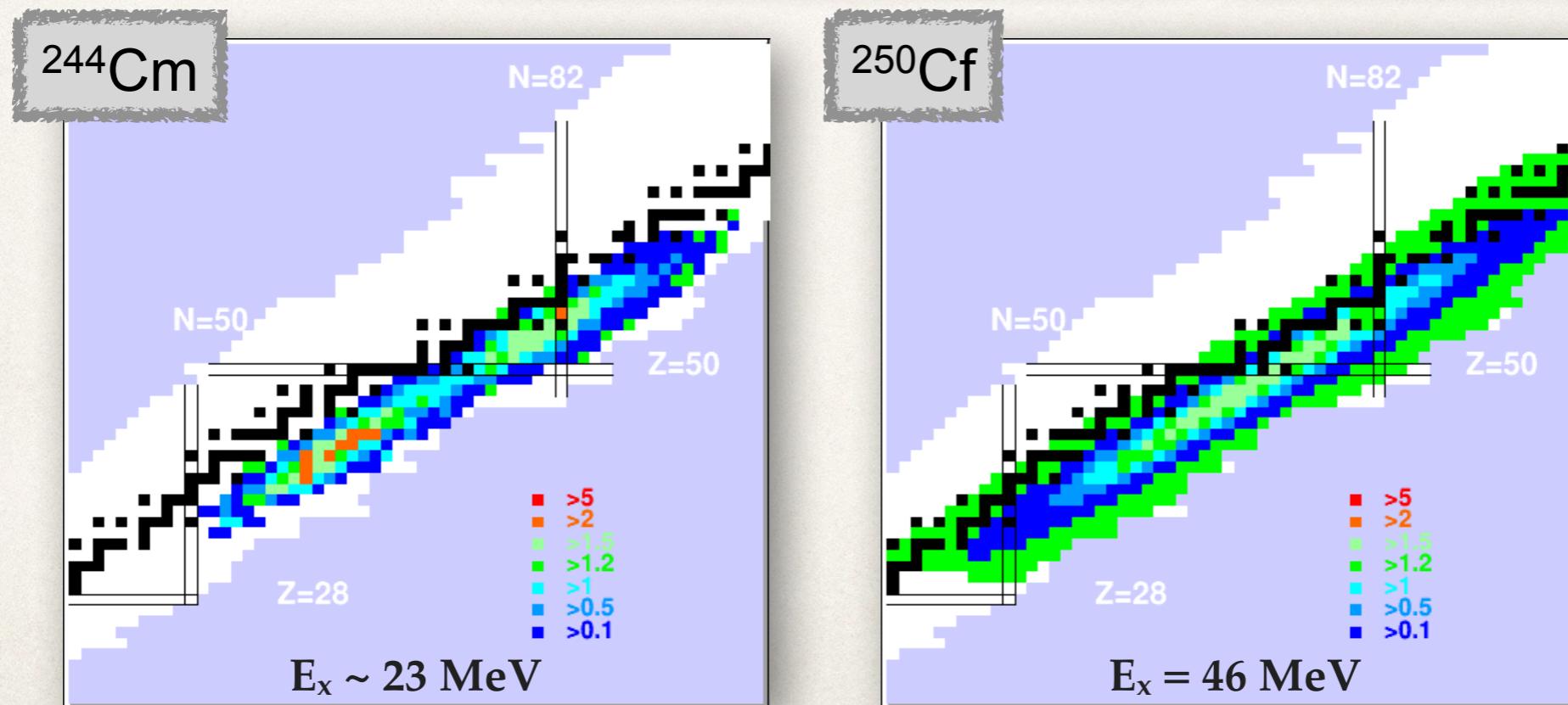
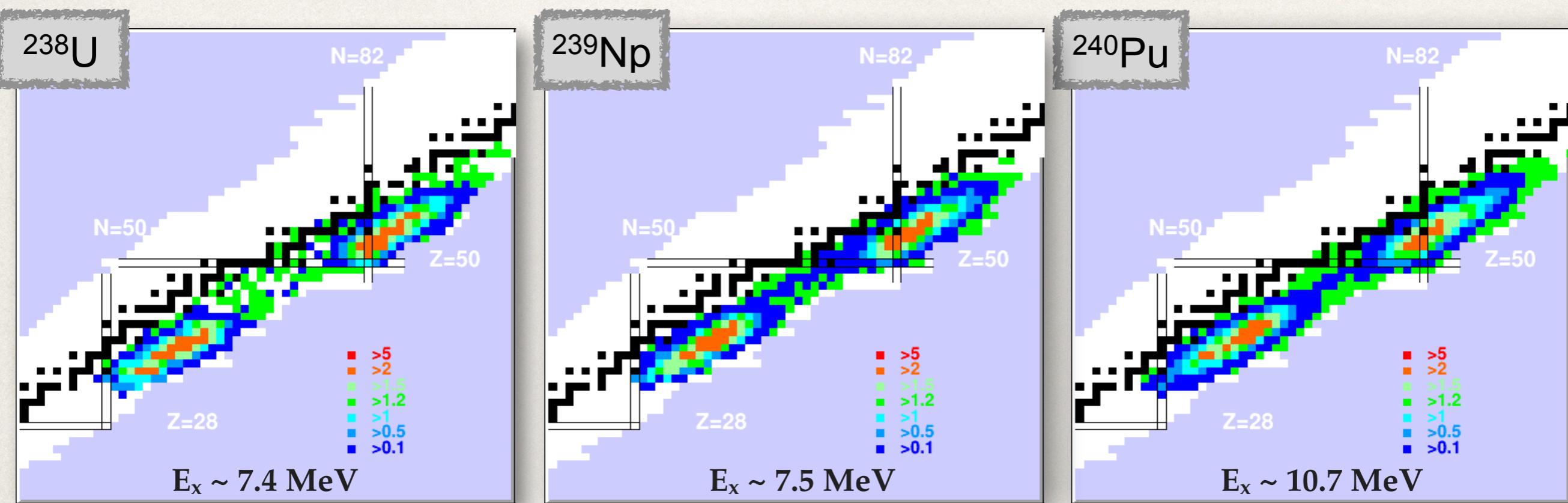
The transmission is different for different isotopes



We need to recover all the charge states per isotope and compensate the acceptance in the azimuthal and polar angles

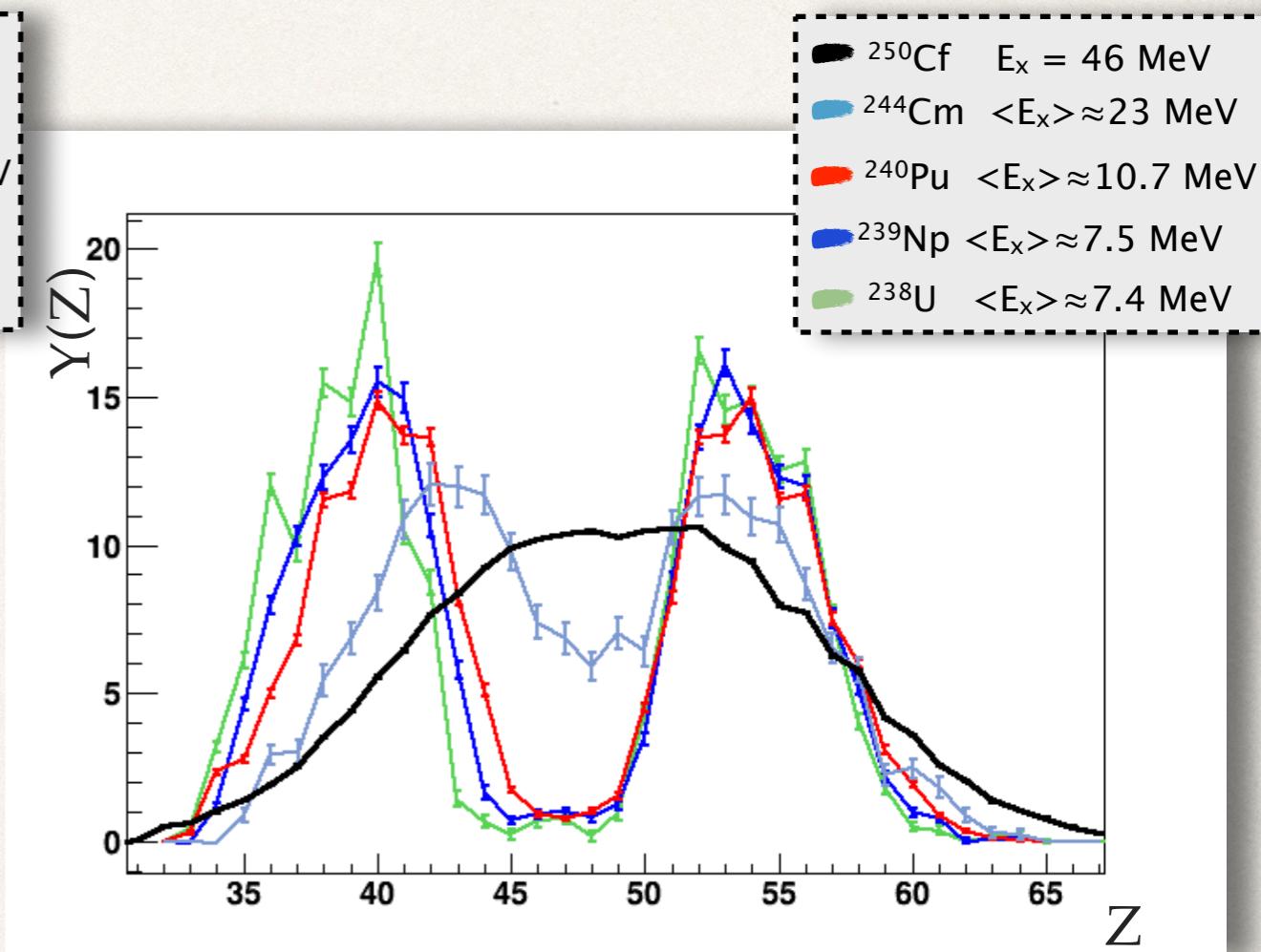
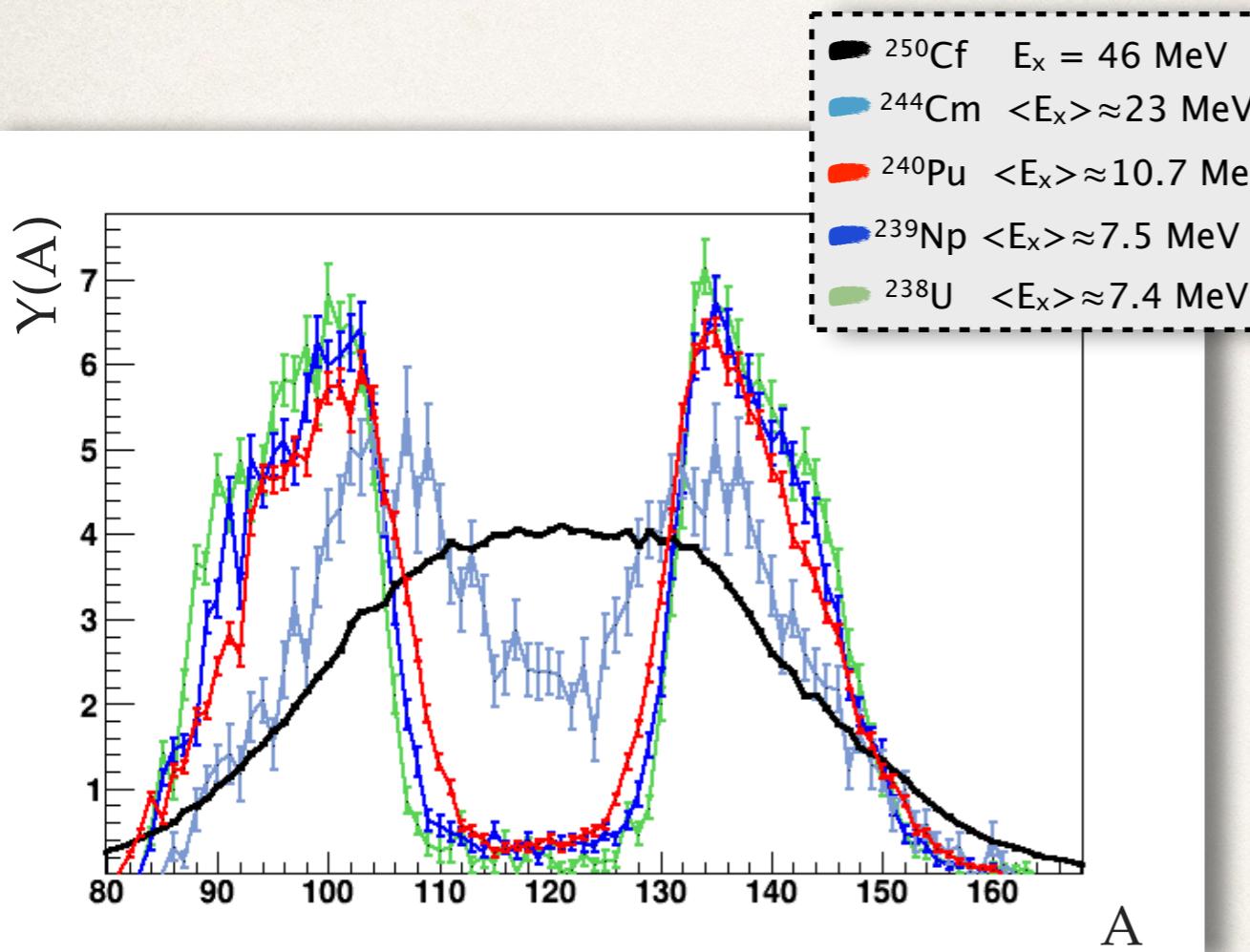
$$Y(Z, A) = I(Z, A) \frac{2}{\text{Range}(Z, A)}$$

Isotopic Fission-Fragment Distribution



Fission Yields

Mass-yields and Z-yields distribution of 5 different fissioning systems, most of them exotic nuclei



New complete measurements, difficult to produce by n-capture

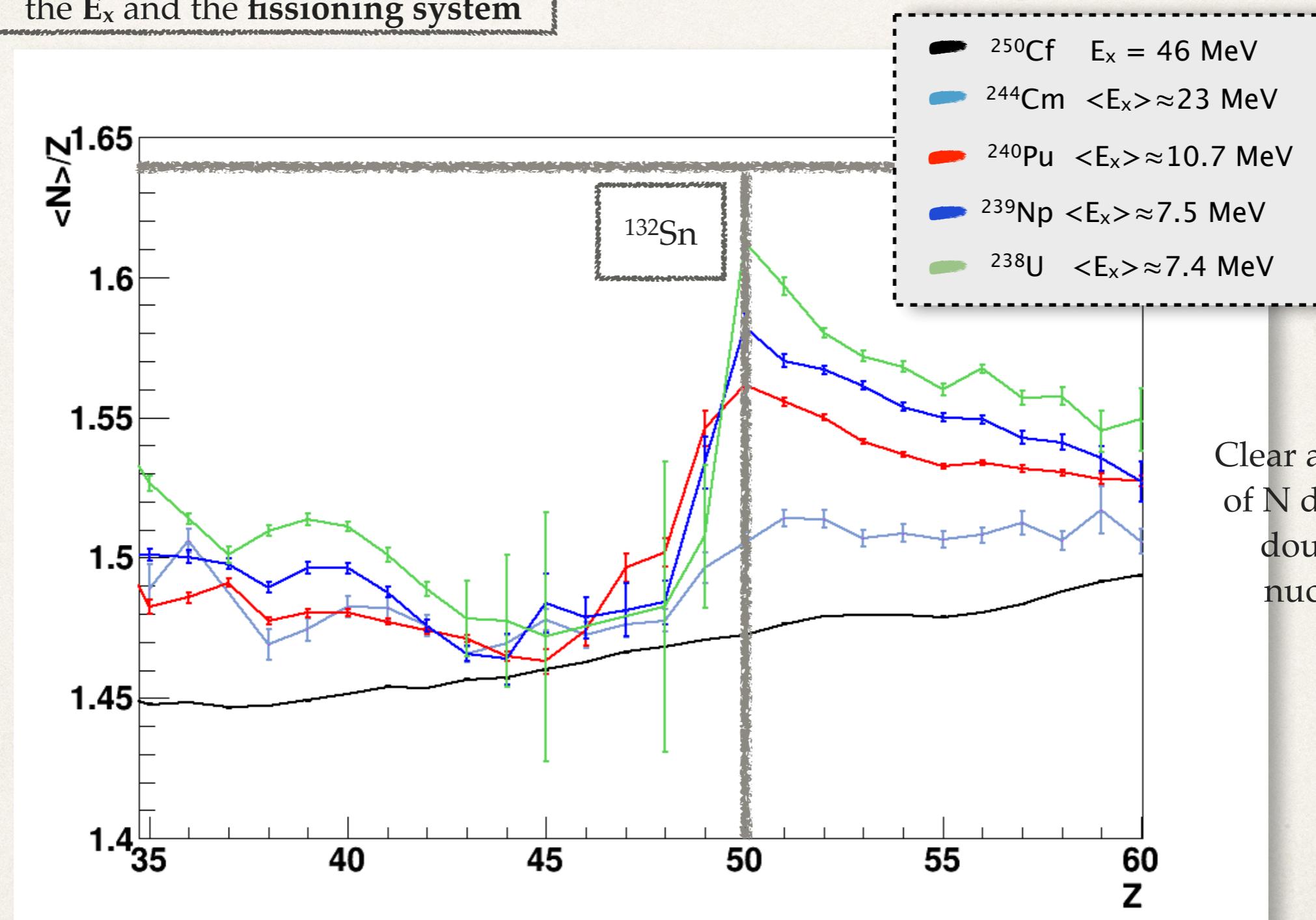
Measurements of fragment distributions of ^{239}Np is scarce ($T_{1/2}(^{238}\text{Np}) = 2.1 \text{ d}$)
There is no direct measurements of fragment distributions of ^{244}Cm

The contribution of the symmetric mode disappears for the systems at low excitation energy

The shift in Z of the light fragments reflects the atomic number of the fissioning system

Neutron Excess

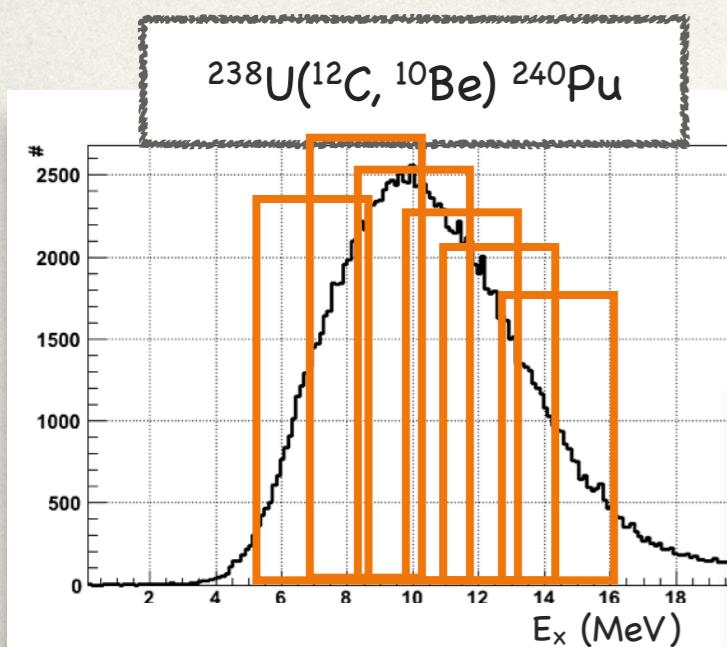
Evolution of the polarization with the E_x and the fissioning system



Charge Polarization present in all the systems

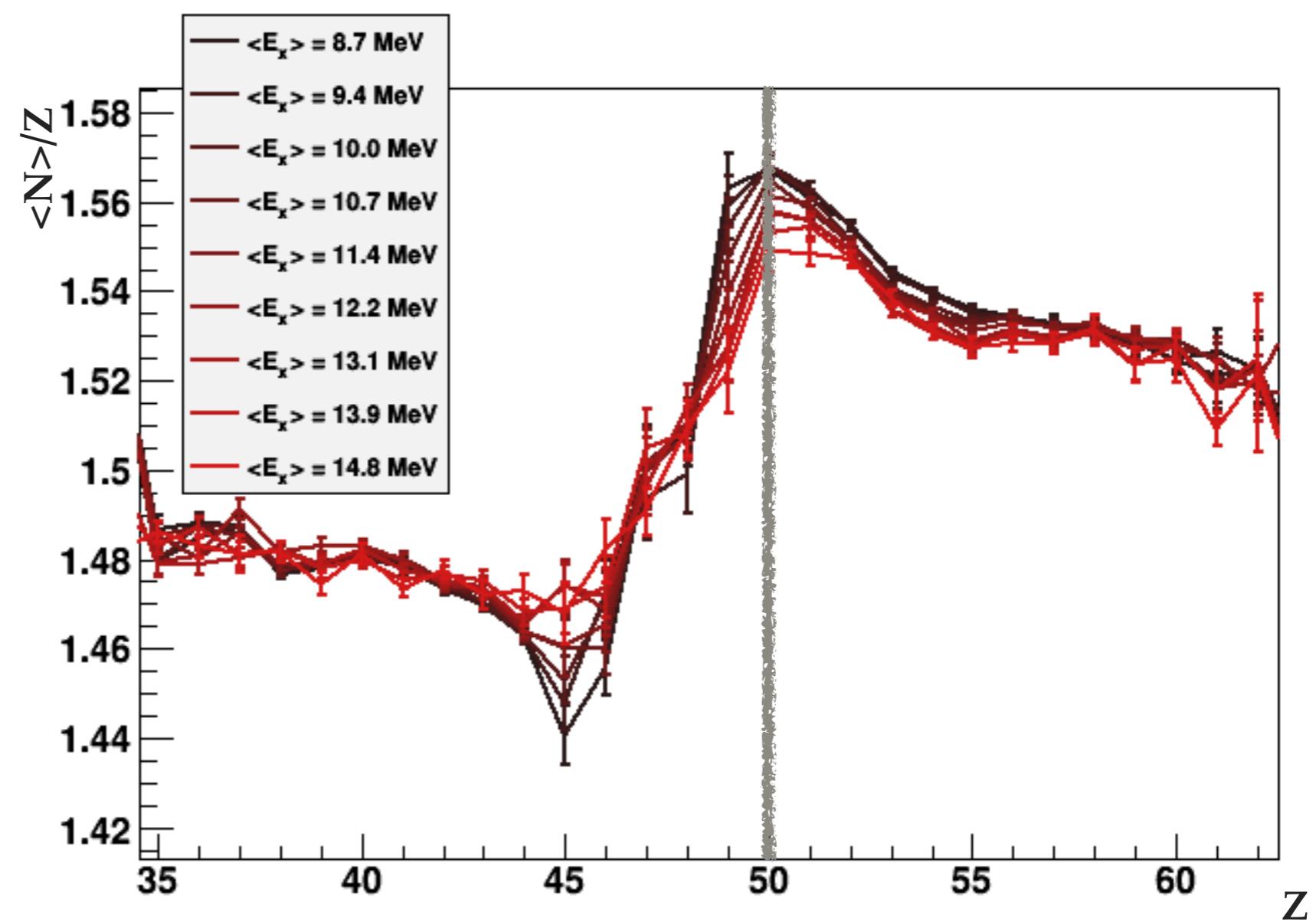
Neutron Excess

^{240}Pu



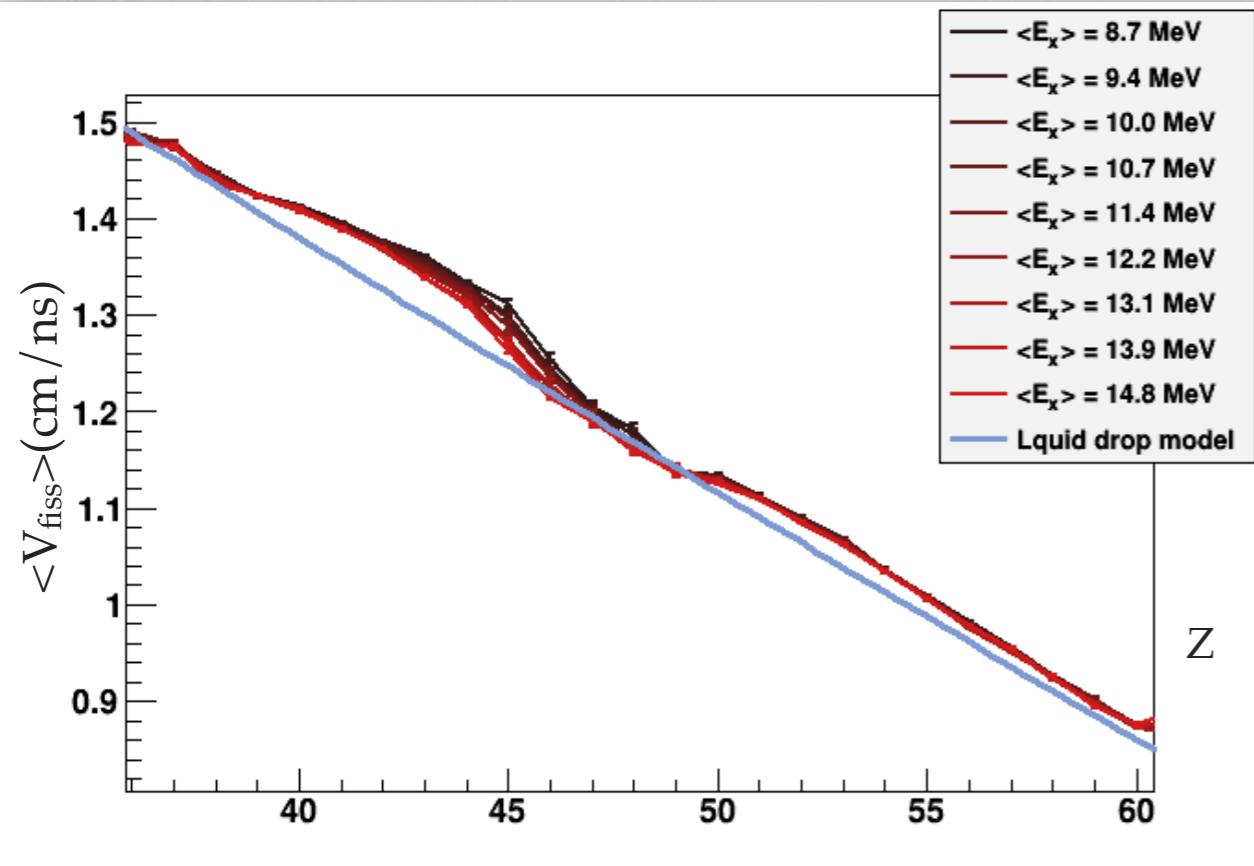
Different positions of gate along the E_x distribution

The $\langle N \rangle / Z$ ratio gets reduced around $Z \approx 50$ by increasing E_x , signature of a closed shell which effect is smaller for higher E_x .



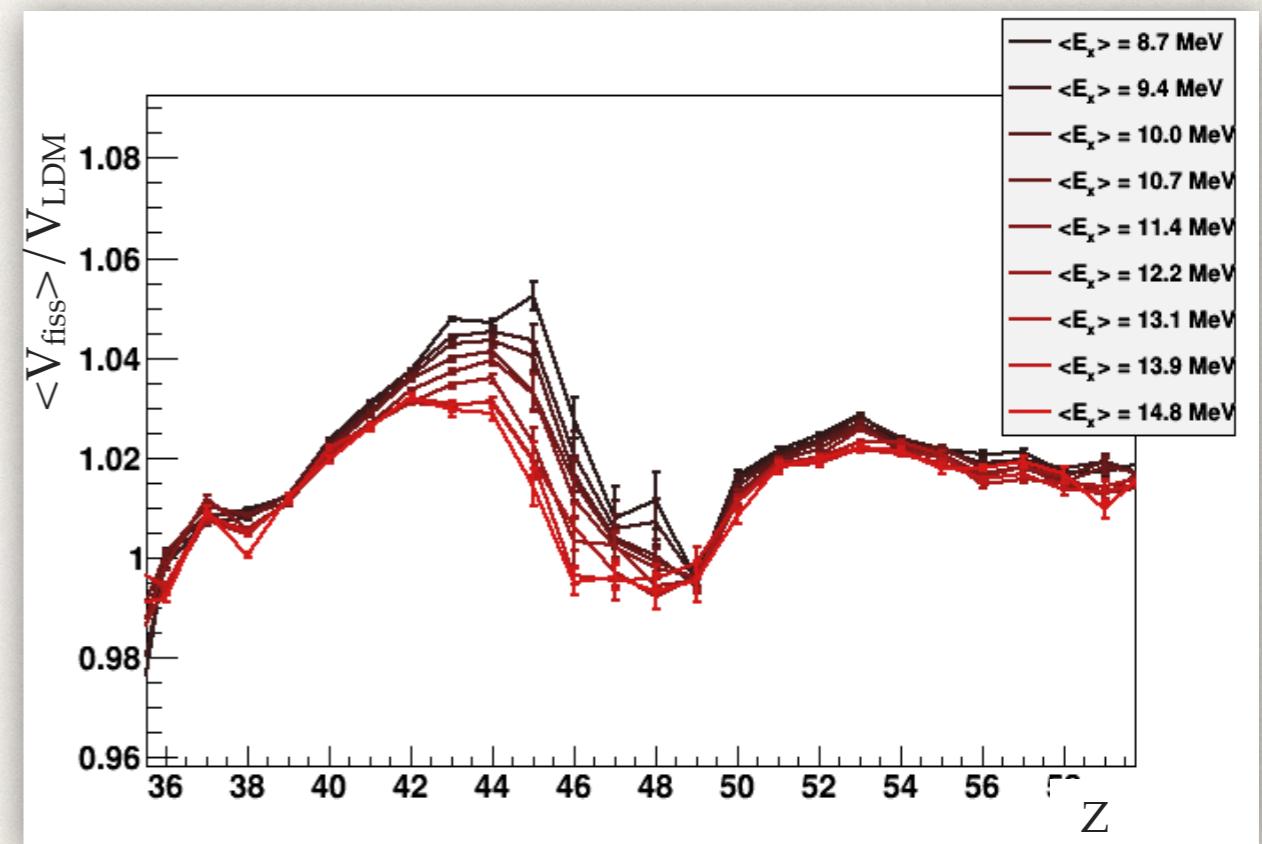
Average Velocity

PRELIMINARY



In the asymmetric region, the light fragment is emitted with a higher velocity compared with the LDM

$$\langle V_{fiss} \rangle = \frac{\sum_A V_{fiss}(Z, A) \cdot Y(Z, A)}{\sum_A Y(Z, A)}$$

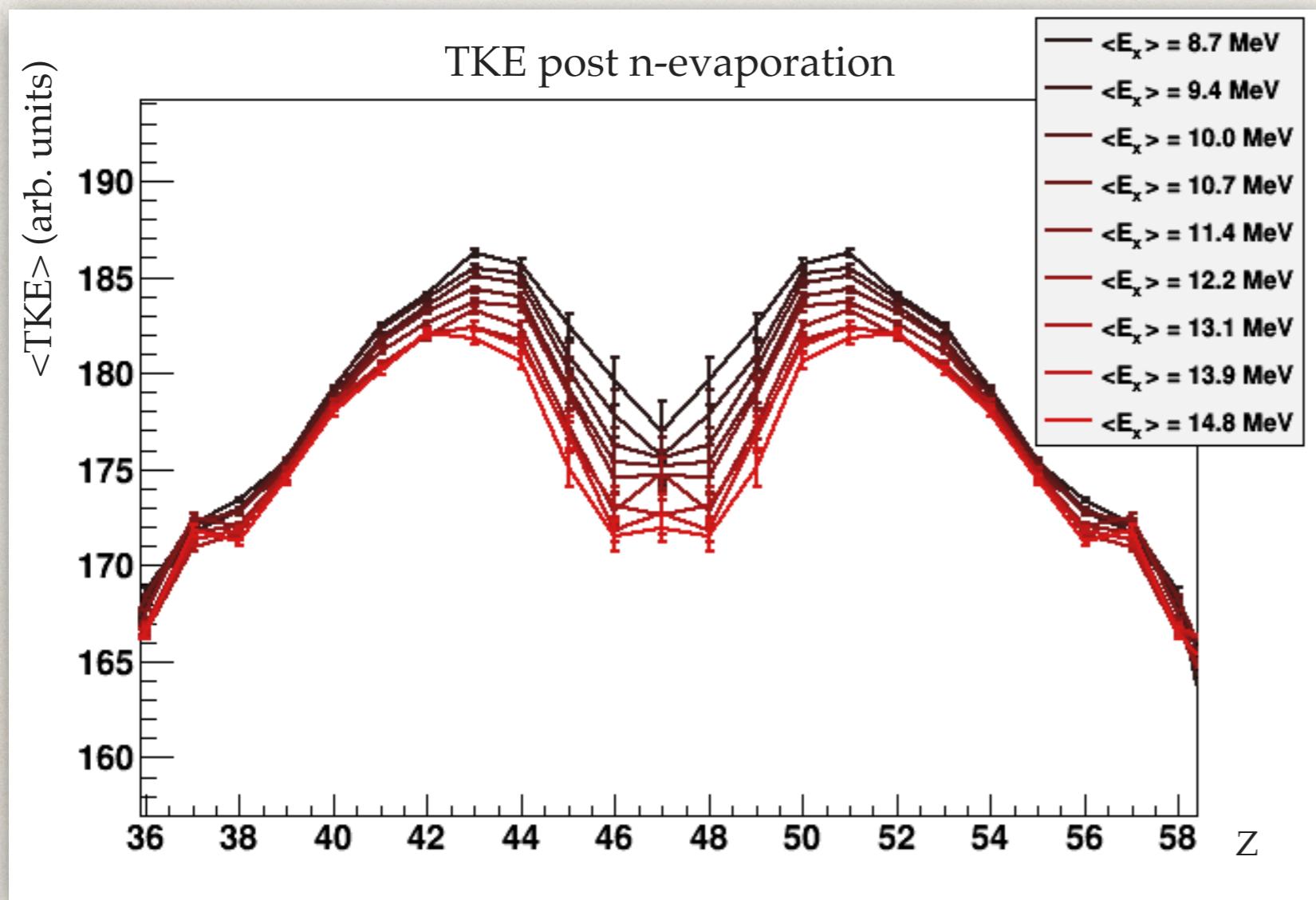


Velocity decrease with higher E_x
The distance between both fragments at the scission point is larger with higher E_x

Average Total Kinetic Energy

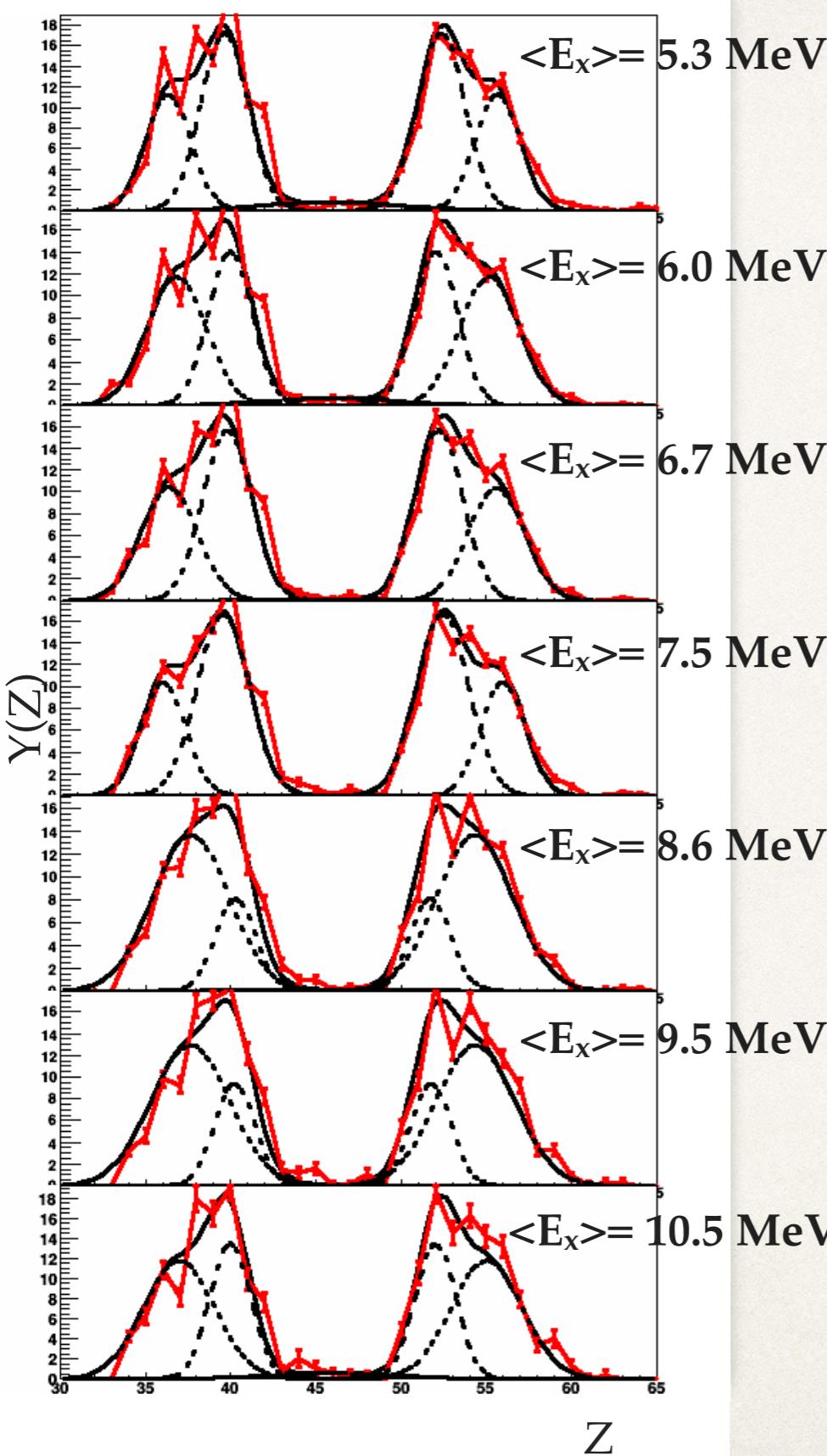
PRELIMINARY

$$TKE = u \cdot \langle A \rangle_Z \cdot (\langle \gamma \rangle_Z - 1) + \\ u \cdot \langle A \rangle_{Z_{Act} - Z} \cdot (\langle \gamma \rangle_{Z_{Act} - Z} - 1)$$

TKE values decrease with higher E_x

- Larger distance at scission point
- Larger neutron evaporation

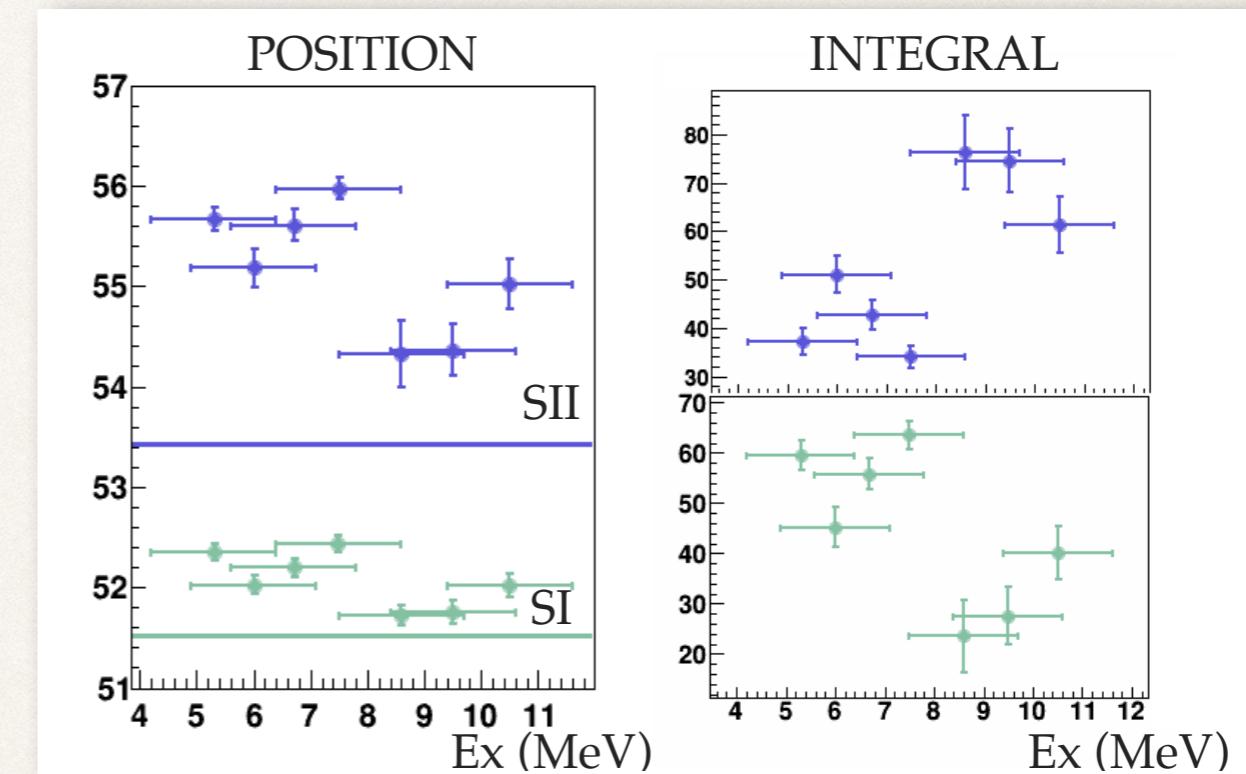
Excitation Energy Dependence



- Super Long (A, μ and σ fixed)
- Standard I
- Standard II

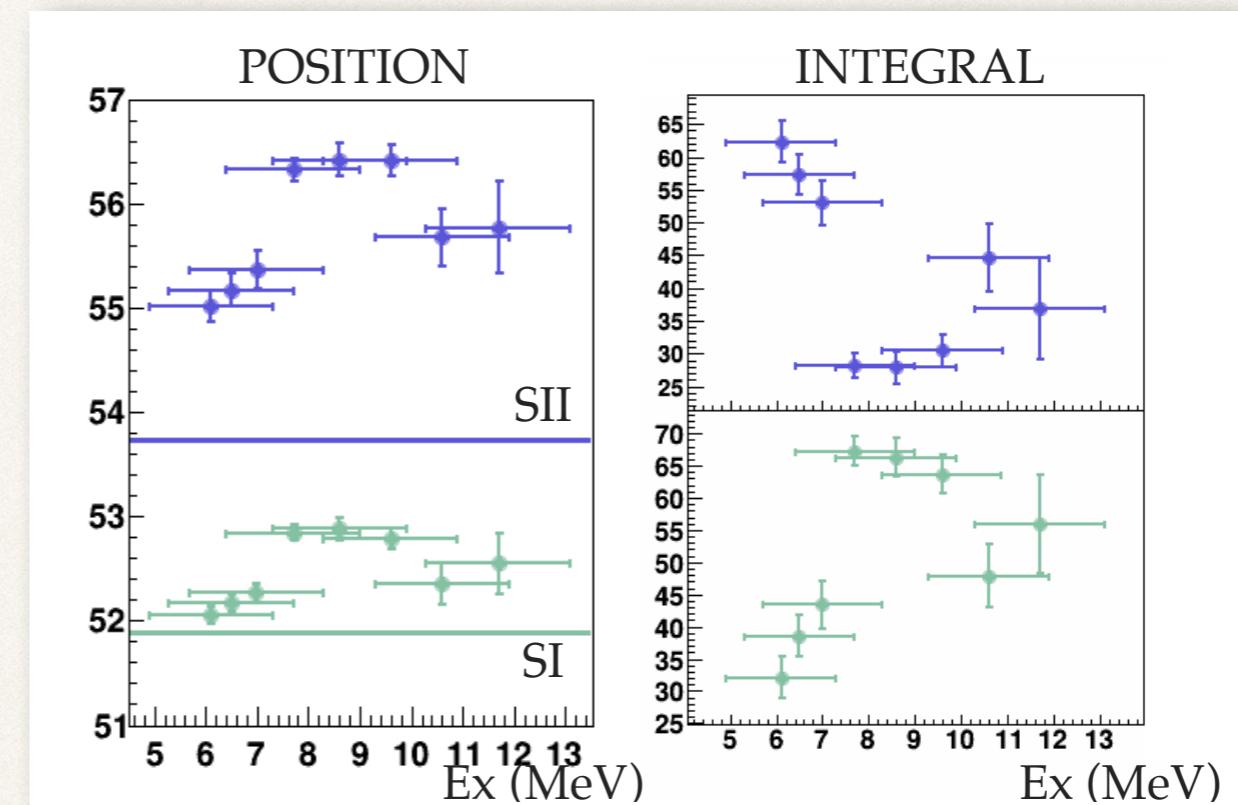
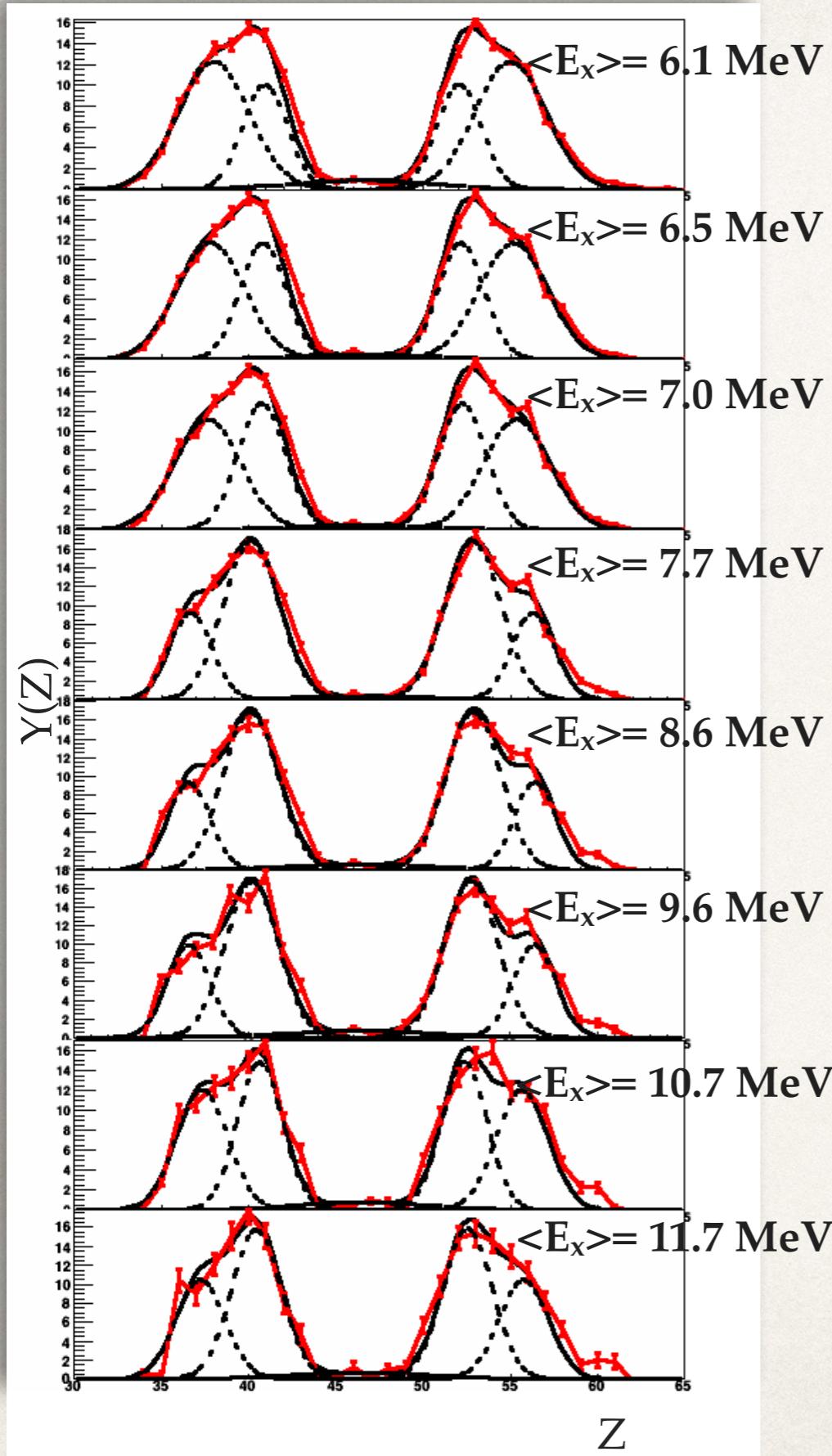
Both Standard positions appear above the K.-H. Schmidt prescription

K.-H. Schmidt et al., Nucl. Phys. A 665, 221 (2000)



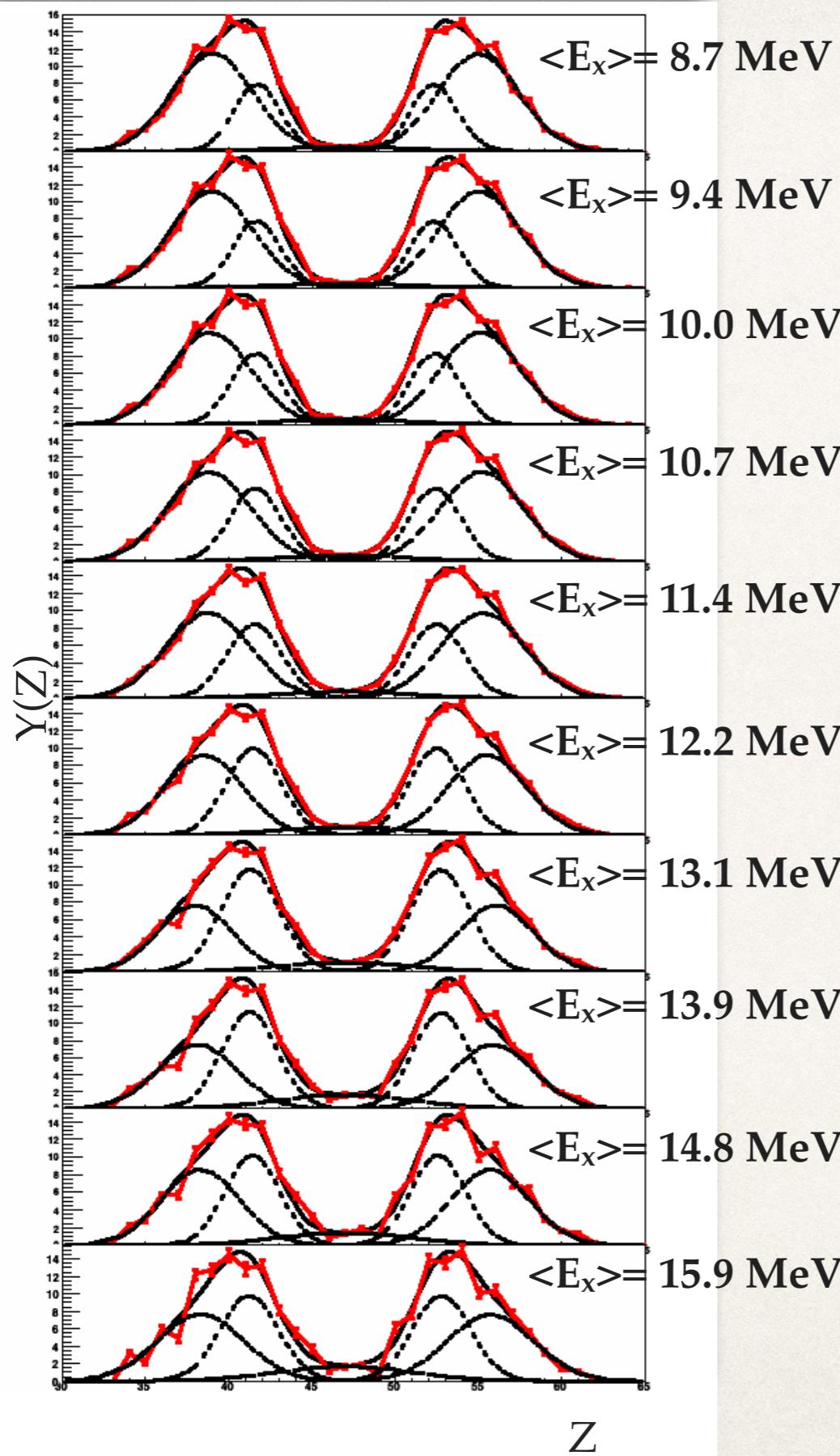
We observe an evolution of the position and the integrals of the asymmetric modes Standard I and Standard II

Excitation Energy Dependence

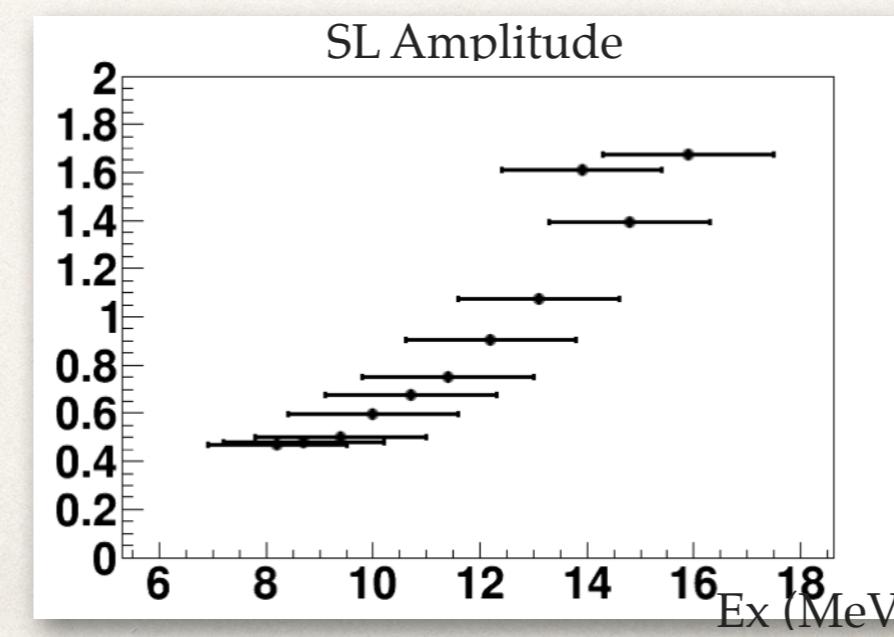
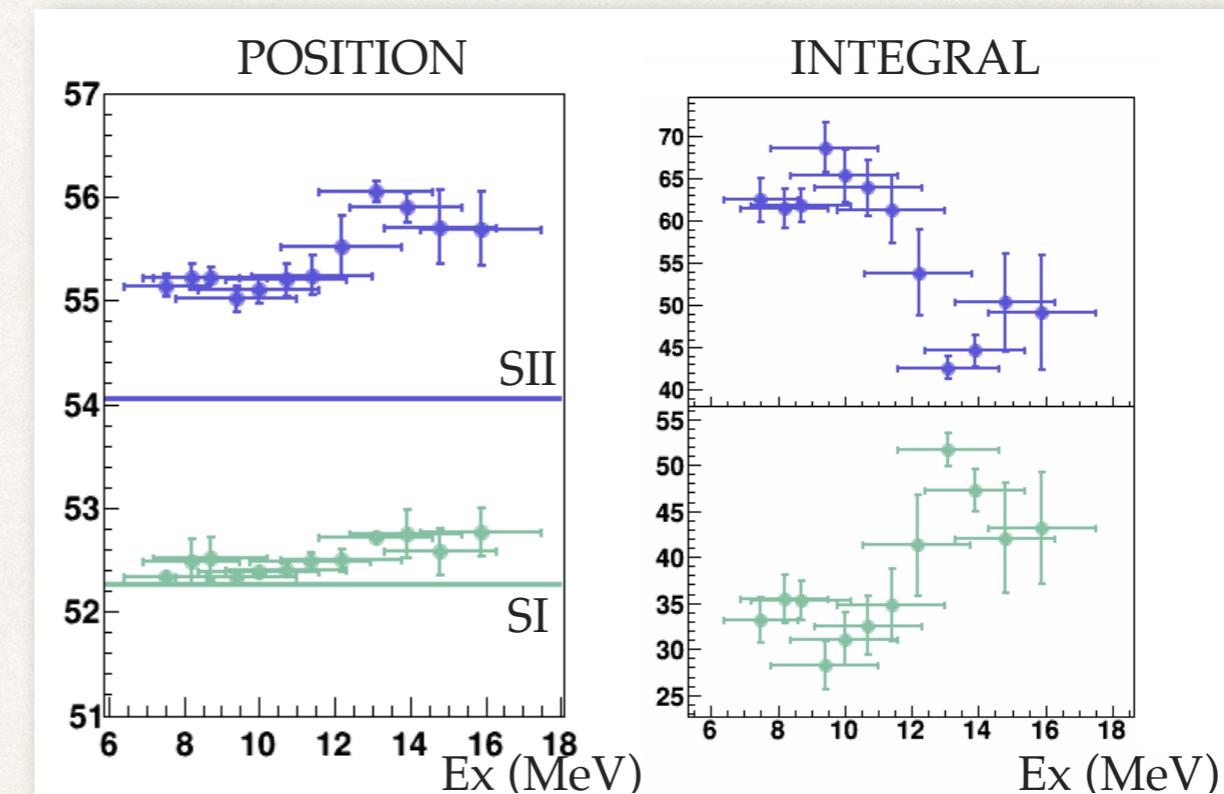


The evolution of the fission channels becomes more clear in this case

Excitation Energy Dependence



Fissioning system with higher statistics
Both positions increase with the E_x for SI and SII. The evolution of the integrals is also clear

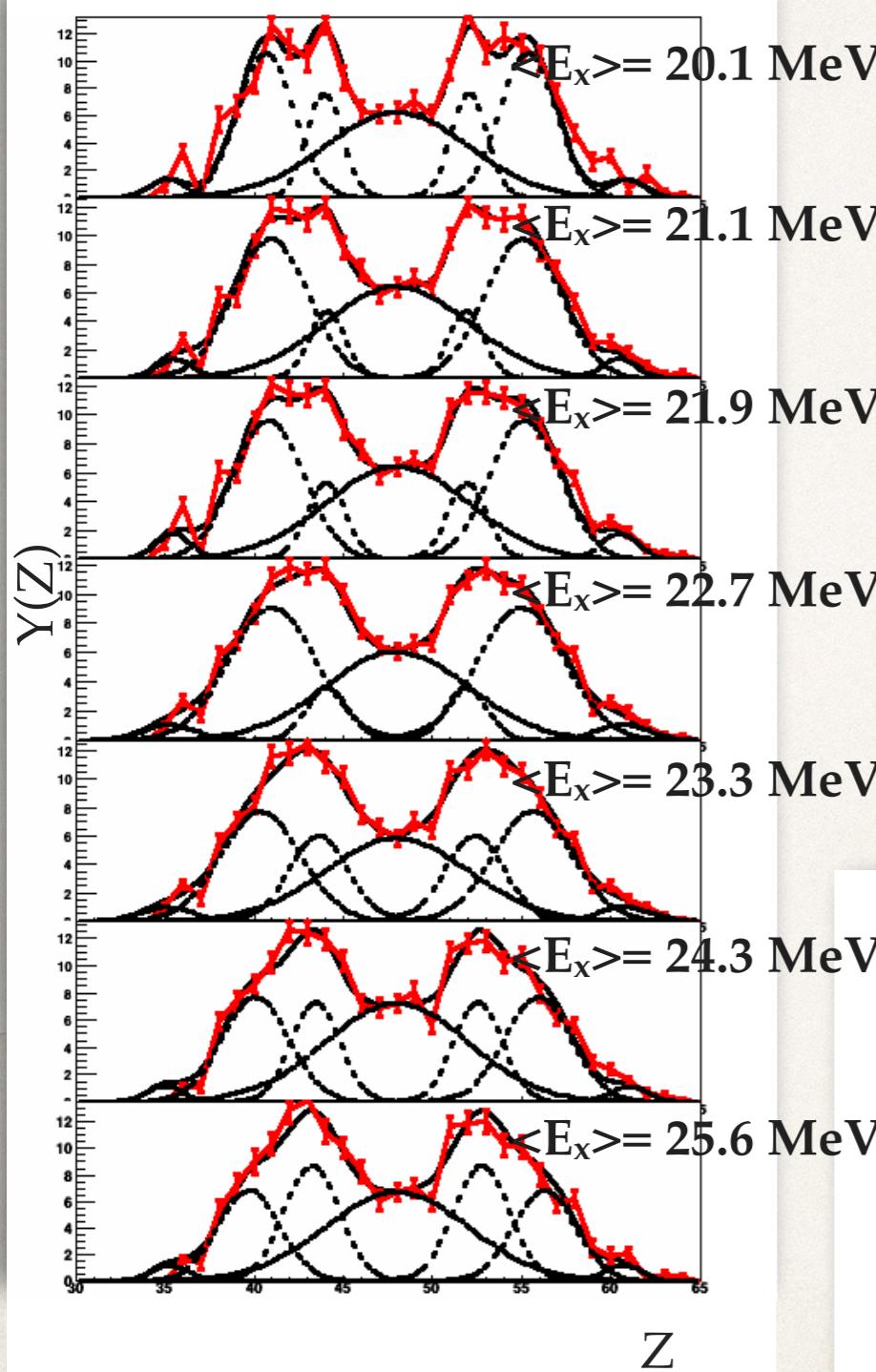


The SuperLong channel becomes more important for higher excitation energy

^{244}Cm

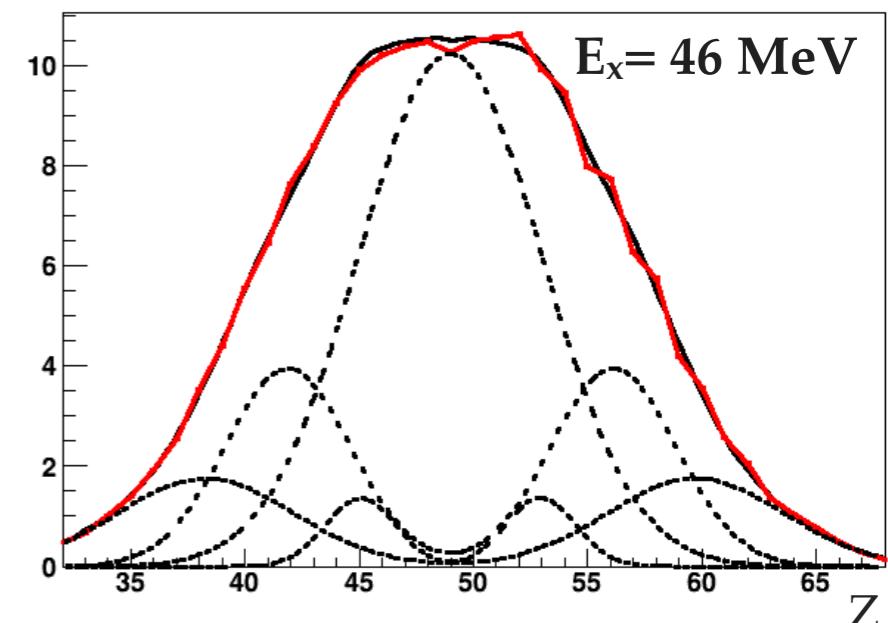
Excitation Energy Dependence

^{250}Cf

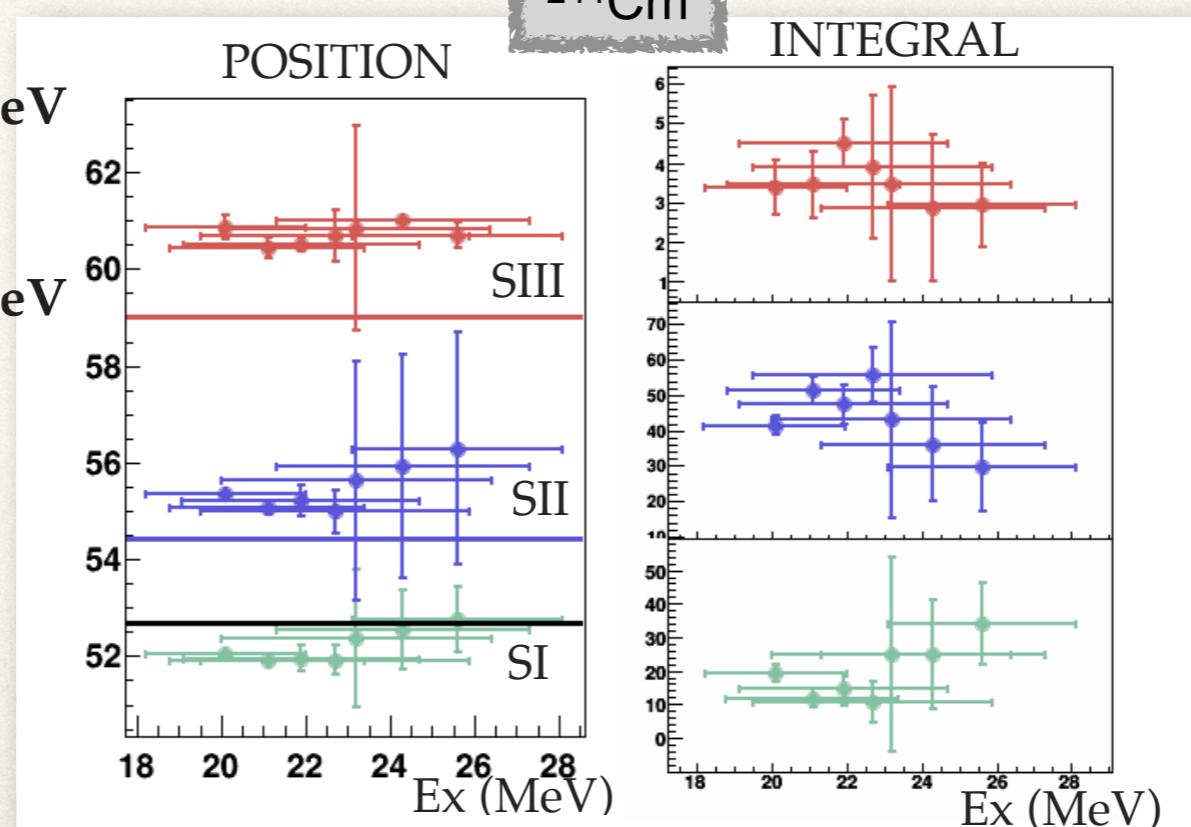


The Standard III is needed to explain the shape of the distribution of ^{244}Cm and ^{250}Cf

SuperLong Channel becomes relevant for these ranges of E_x



SI : $\bar{Z} = 52.9$ I=5.22
SII : $\bar{Z} = 56.1$ I=27.13
SIII : $\bar{Z} = 59.7$ I=16.80



Conclusions

Transfer-induced fission in inverse kinematics coupled to the VAMOS spectrometer allowed:

- The study the fission of different fissioning systems (^{238}U , ^{239}Np , ^{240}Pu , ^{244}Cm and ^{250}Cf).
- The excitation energy measurement distribution through the transfer reaction kinematics.
- The full isotopic identification of fission fragments using the VAMOS spectrometer.

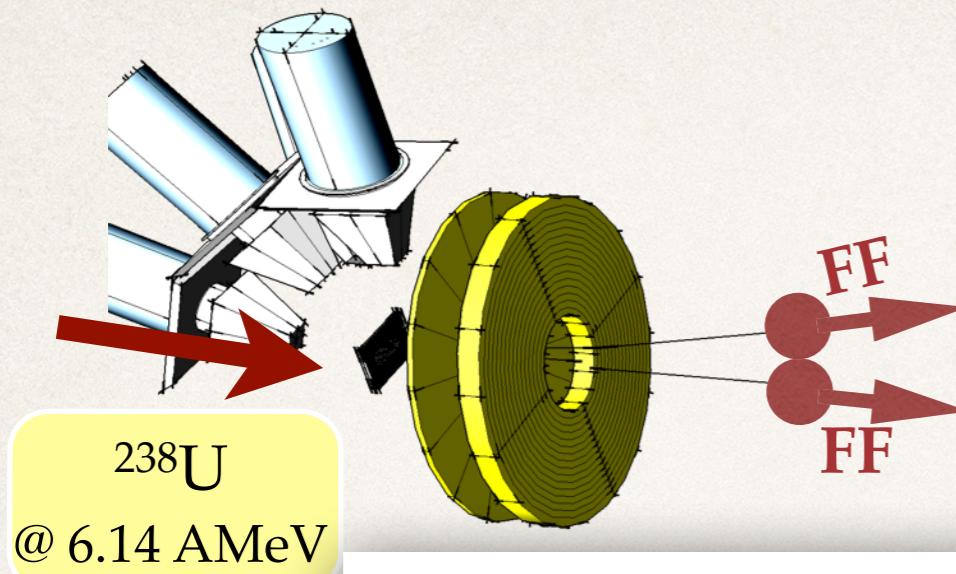
The effect of shell structure was observed as a function of the excitation energy.

- The evolution of the fission fragments from asymmetric to symmetric distributions.
- The $\langle N \rangle/Z$ ratio at $Z \approx 50$ decreases by increasing the excitation energy.
- The TKE decreases by increasing the excitation energy.

Fission Fragment Distribution were investigated in terms of Fission Channels Interpretation

- The SI, SII and SIII are needed to reproduce the almost gaussian shape of ^{250}Cf
- Observed an evolution of the position and the integral of SI and SII with the excitation energy depending on the fissioning system

Fission Barriers



We observe a general agreement with previous data
with small discrepancies

First and second chance are measured
in ^{240}Pu and ^{241}Pu

238U

@ 6.14 AMeV

