

## **FUSTIPEN Topical Meeting**

« Fission-fragments in low-energy fission: a gauge for macroscopic  
and microscopic influences »

# From fission yields to scission properties

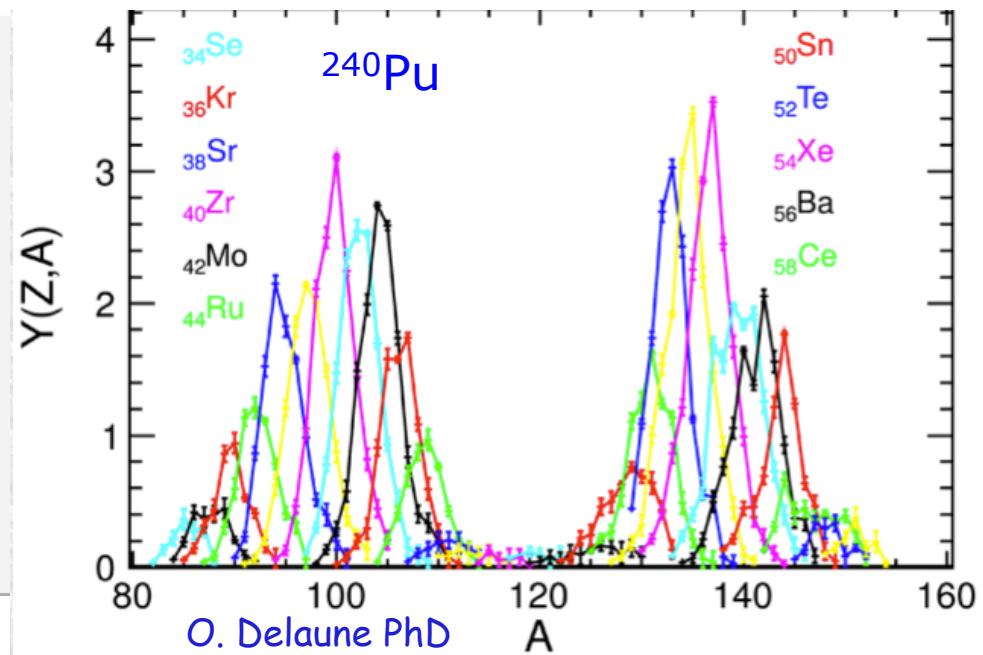
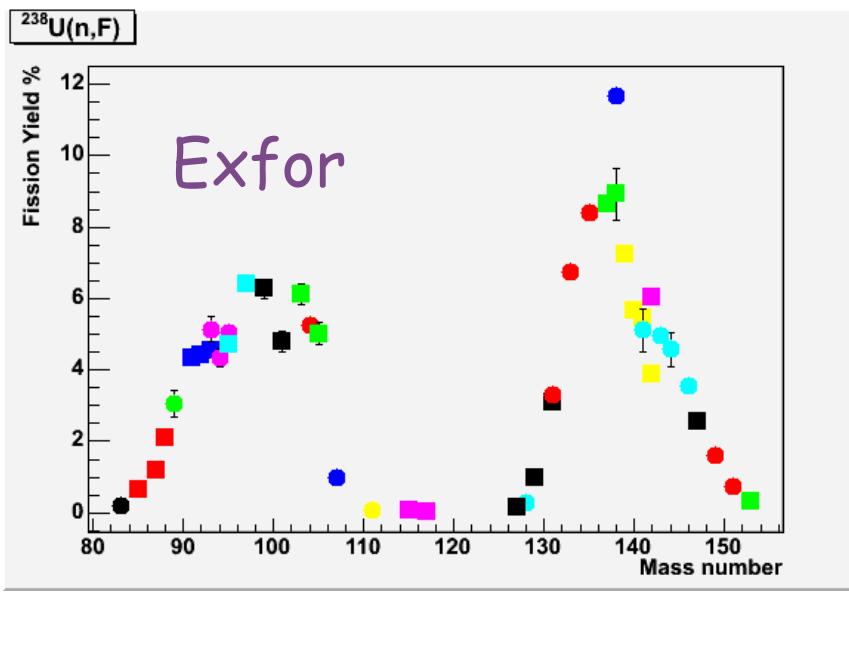
Fanny Farget,

GANIL

October 22<sup>nd</sup>, 2015

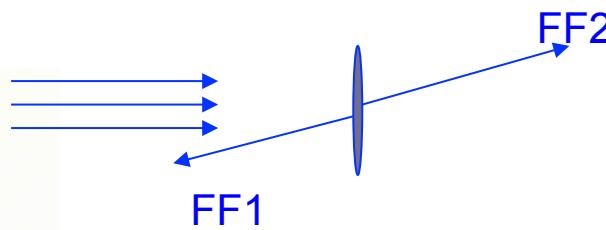
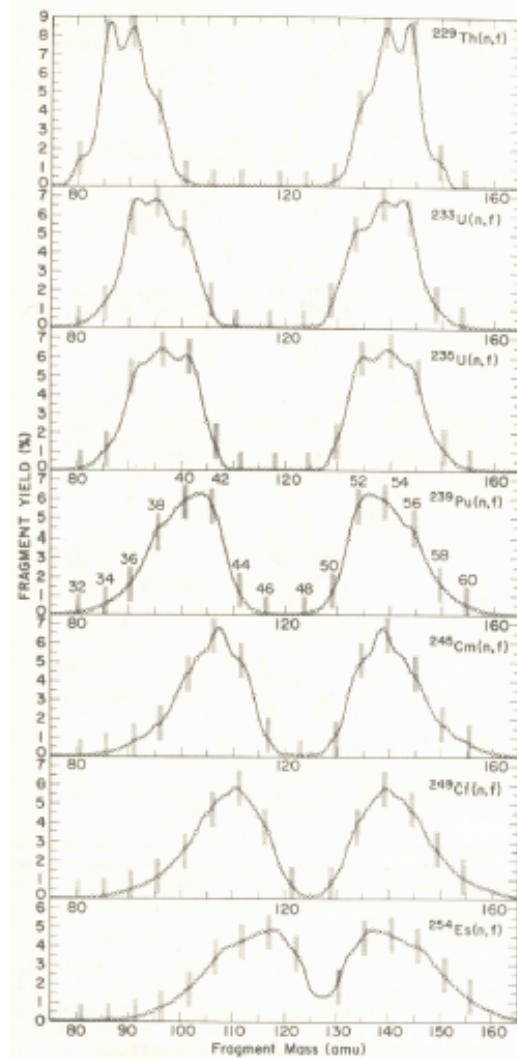
# The strength of inverse kinematics for fission

# Isotopic fission yields



# Fission yields in direct kinematics

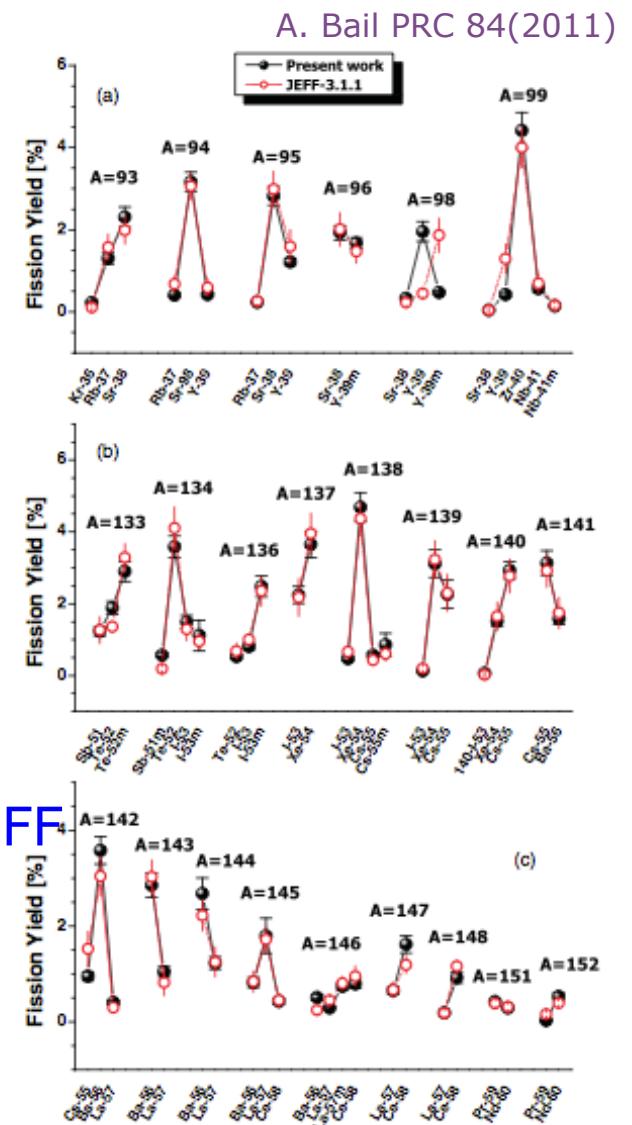
J.P. Unik, IAEA -SM-174/209



Mass distribution: OK

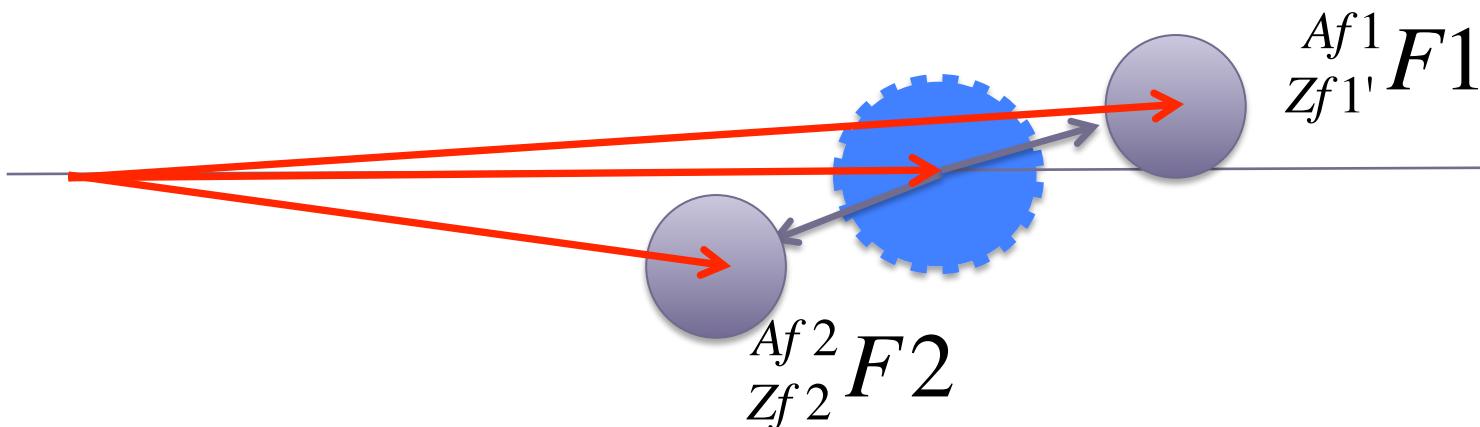
Isotopic distribution:  
prompt or  $\beta$ -delayed  
spectroscopy

Limited by the  
-lifetime of the FF  
-unknown level scheme of FF

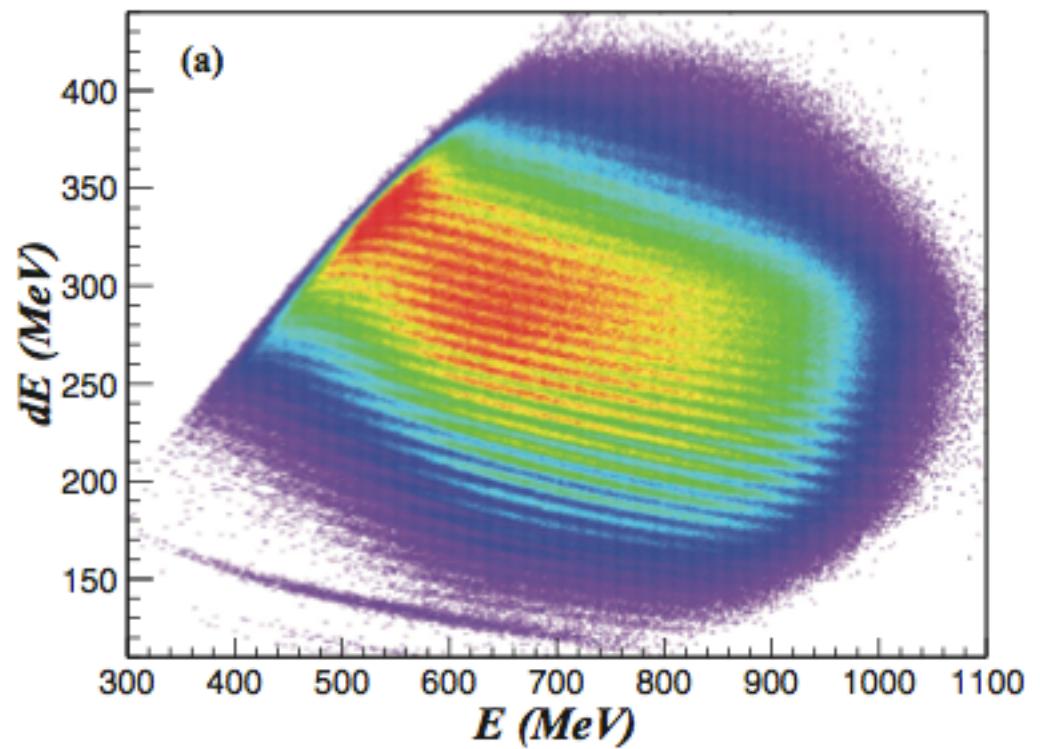


A. Bail PRC 84(2011)

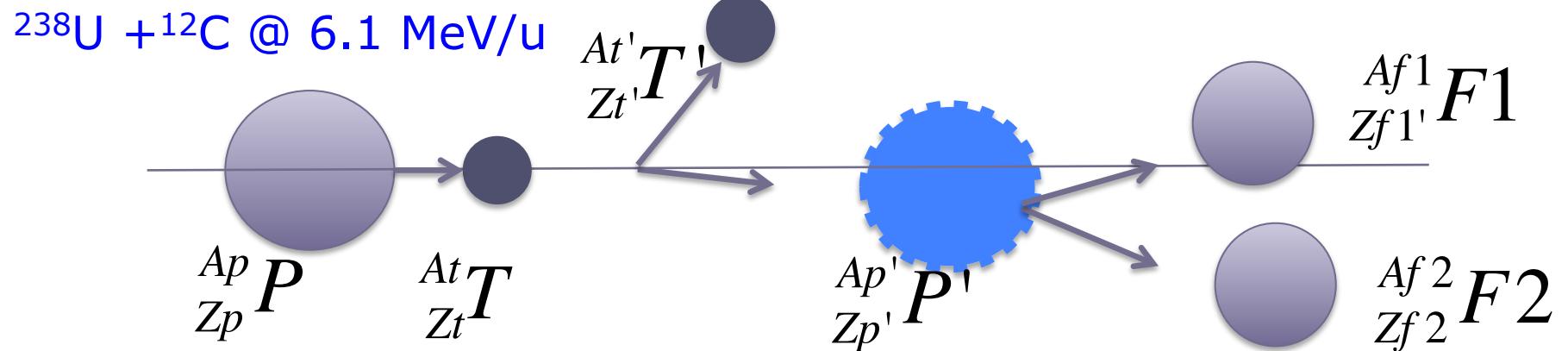
## Fission in inverse kinematics: kinematical boost for a direct identification of the fission fragments



We will see in the following slides  
(and presentations)  
That inverse kinematics brings  
more than isotopic distribution



# Transfer-induced fission in inverse kinematics @ GANIL



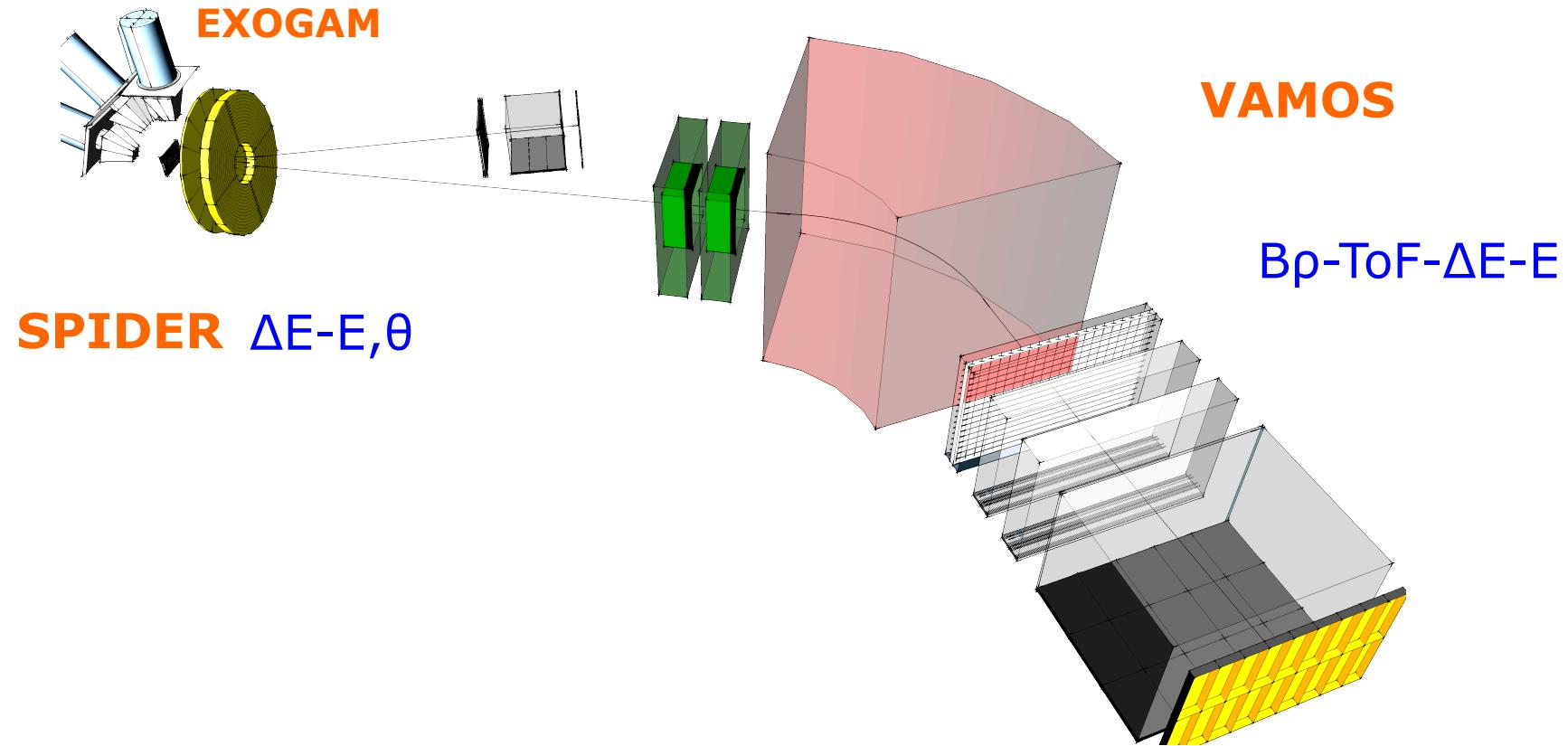
$^{242}\text{Cf}$	$^{243}\text{Cf}$	$^{244}\text{Cf}$	$^{245}\text{Cf}$	$^{246}\text{Cf}$	$^{247}\text{Cf}$	$^{248}\text{Cf}$	$^{249}\text{Cf}$	$^{250}\text{Cf}$	$^{251}\text{Cf}$	$^{252}\text{Cf}$
$^{241}\text{Bk}$	$^{242}\text{Bk}$	$^{243}\text{Bk}$	$^{244}\text{Bk}$	$^{245}\text{Bk}$	$^{246}\text{Bk}$	$^{247}\text{Bk}$	$^{248}\text{Bk}$	$^{249}\text{Bk}$	$^{250}\text{Bk}$	$^{251}\text{Bk}$
$^{240}\text{Cm}$	$^{241}\text{Cm}$	$^{242}\text{Cm}$	$^{243}\text{Cm}$	$^{244}\text{Cm}$	$^{245}\text{Cm}$	$^{246}\text{Cm}$	$^{247}\text{Cm}$	$^{248}\text{Cm}$	$^{249}\text{Cm}$	$^{250}\text{Cm}$
$^{239}\text{Am}$	$^{240}\text{Am}$	$^{241}\text{Am}$	$^{242}\text{Am}$	$^{243}\text{Am}$	$^{244}\text{Am}$	$^{245}\text{Am}$	$^{246}\text{Am}$	$^{247}\text{Am}$	$^{248}\text{Am}$	$^{249}\text{Am}$
$^{238}\text{Pu}$	$^{239}\text{Pu}$	$^{240}\text{Pu}$	$^{241}\text{Pu}$	$^{242}\text{Pu}$	$^{243}\text{Pu}$	$^{244}\text{Pu}$	$^{245}\text{Pu}$	$^{246}\text{Pu}$	$^{247}\text{Pu}$	
$^{237}\text{Np}$	$^{238}\text{Np}$	$^{239}\text{Np}$	$^{240}\text{Np}$	$^{241}\text{Np}$	$^{242}\text{Np}$	$^{243}\text{Np}$	$^{244}\text{Np}$			
$^{236}\text{U}$	$^{237}\text{U}$	$^{238}\text{U}$	$^{239}\text{U}$	$^{240}\text{U}$	$^{241}\text{U}$	$^{242}\text{U}$				

- 10 actinides produced
- $E^*$  distribution
- Full resolution in (Z,A) of fragments
- TKE
- Détermination of scission fragments

Can't choose your actinide  
Can't choose your  $E^*$

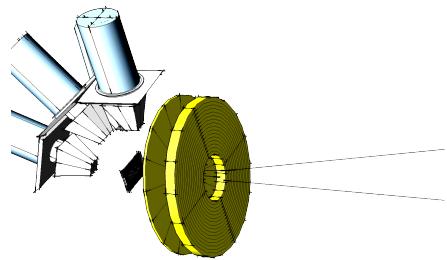
See talk of D. Ramos

# Transfer-induced fission in inverse kinematics



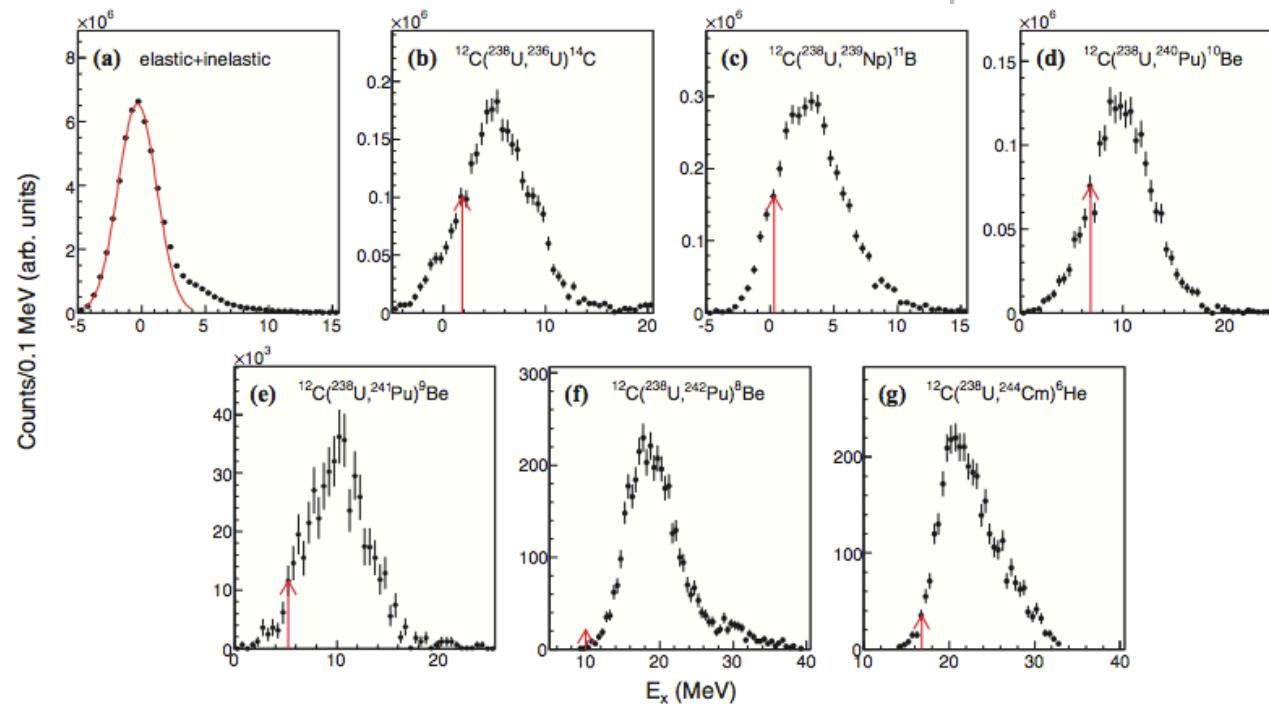
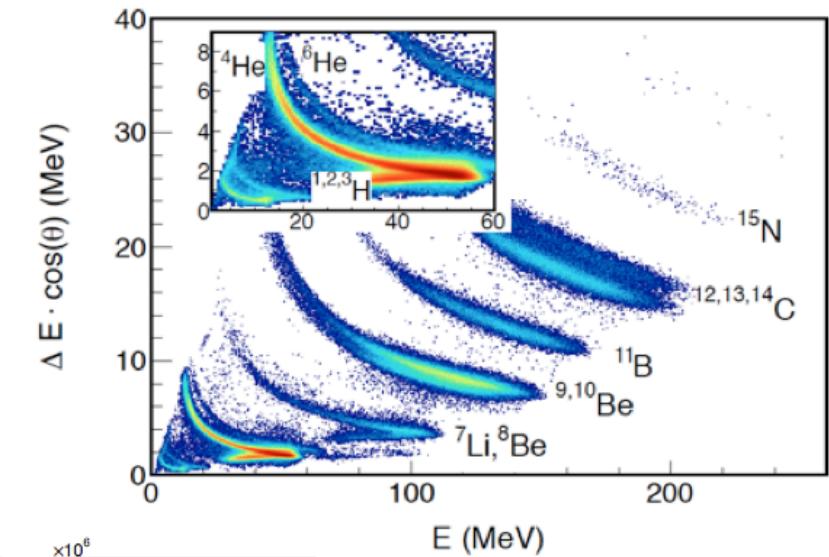
S. Pullanhiotan et al., NIM 593 (2008) 343  
M. Rejmund et al., NIMA 646 (2011) 184

# Transfer-induced fission in inverse kinematics

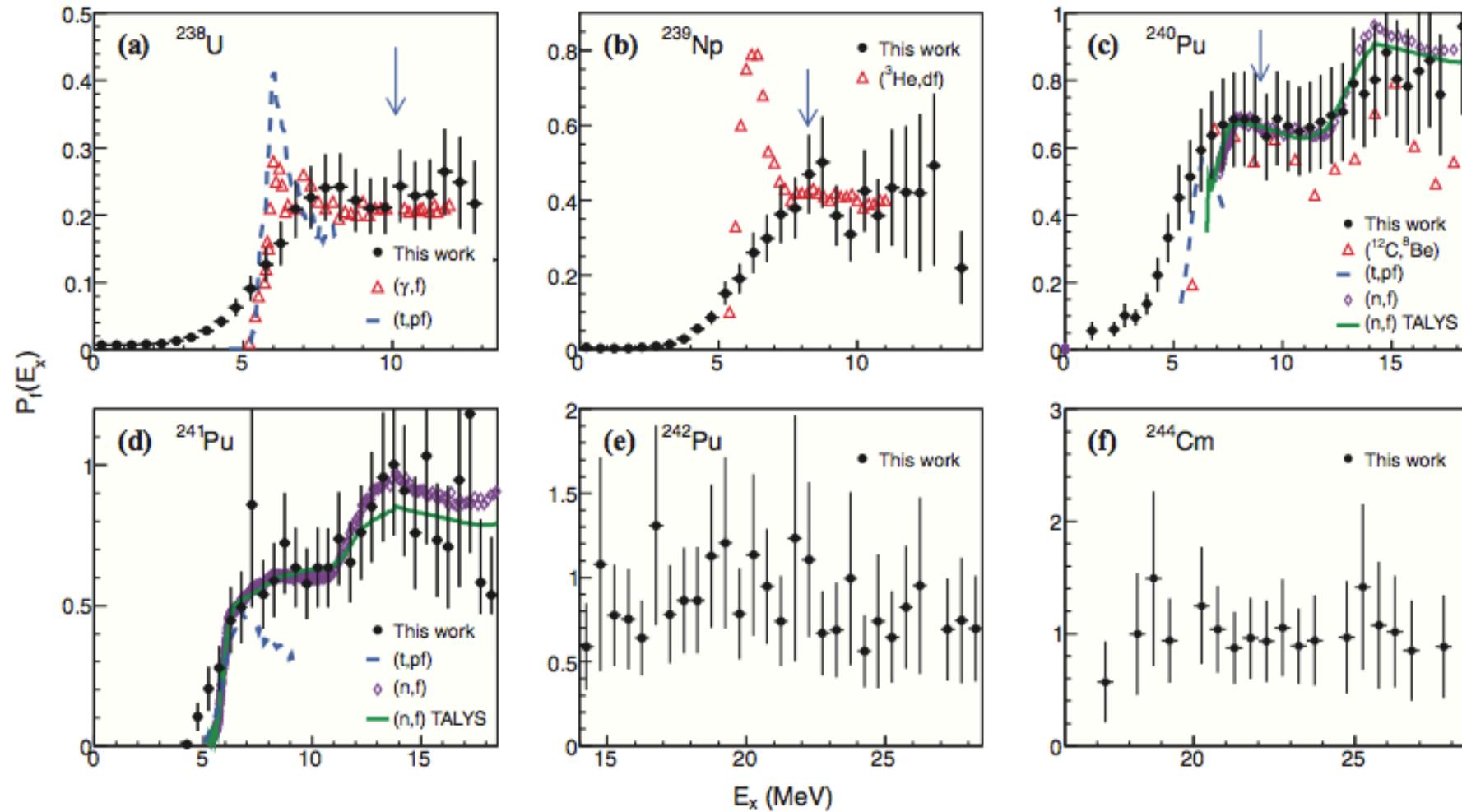


**SPIDER  $\Delta E - E, \theta$**

C. Rodriguez-Tajes et al., PRC (2014) 024614

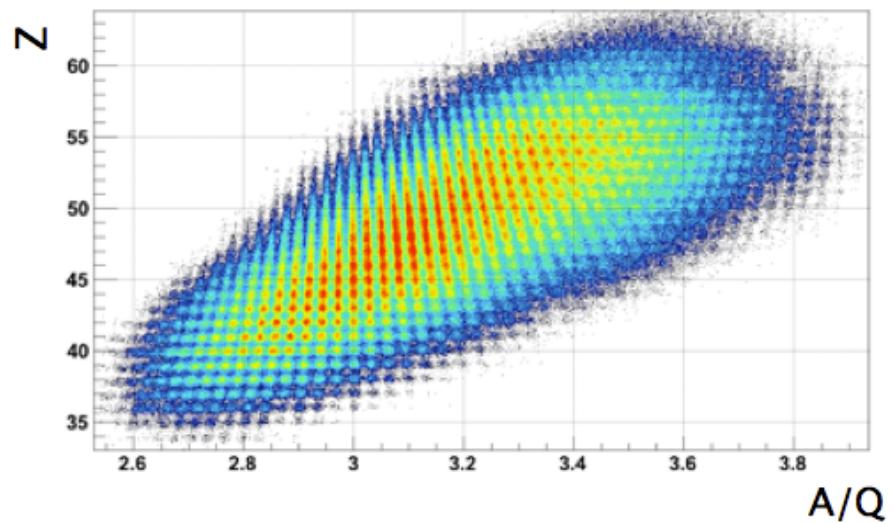
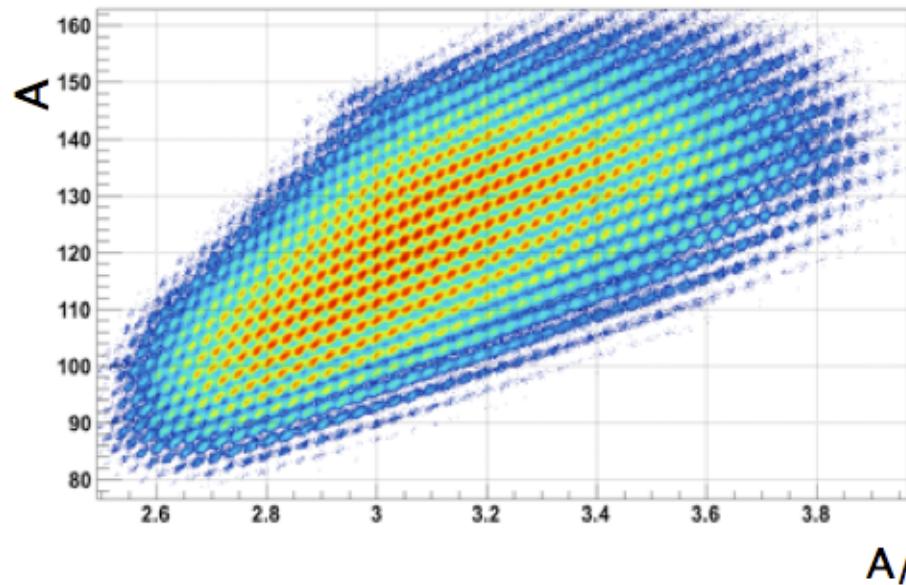
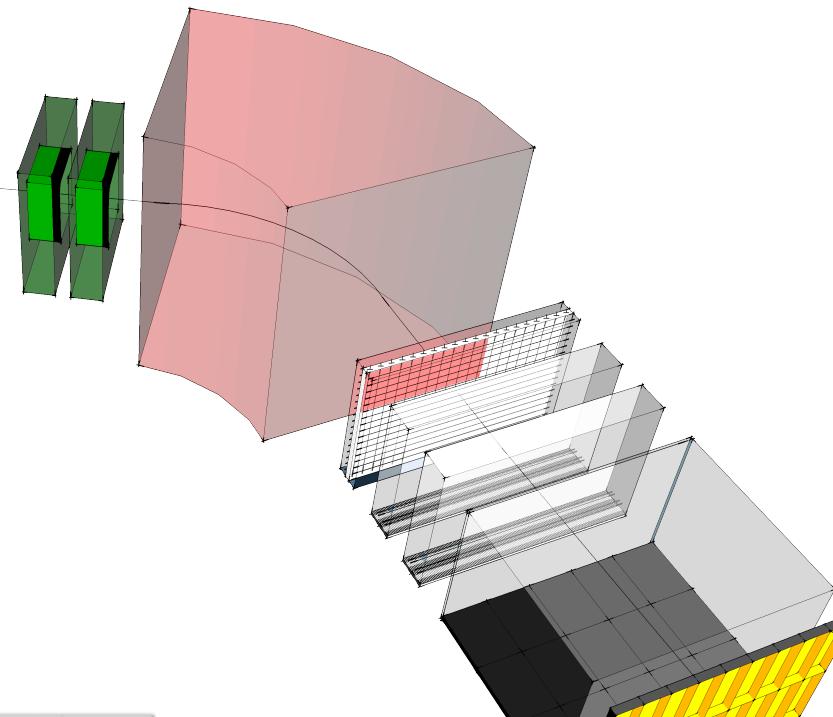
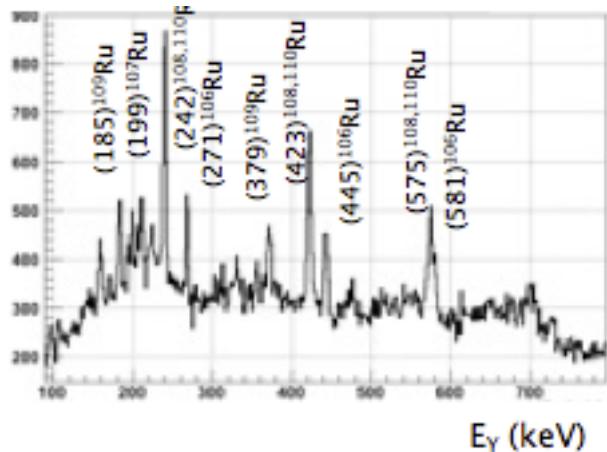
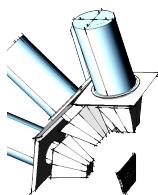


# Fission probabilities



Limitations : excitation energy resolution  
angular momentum effects

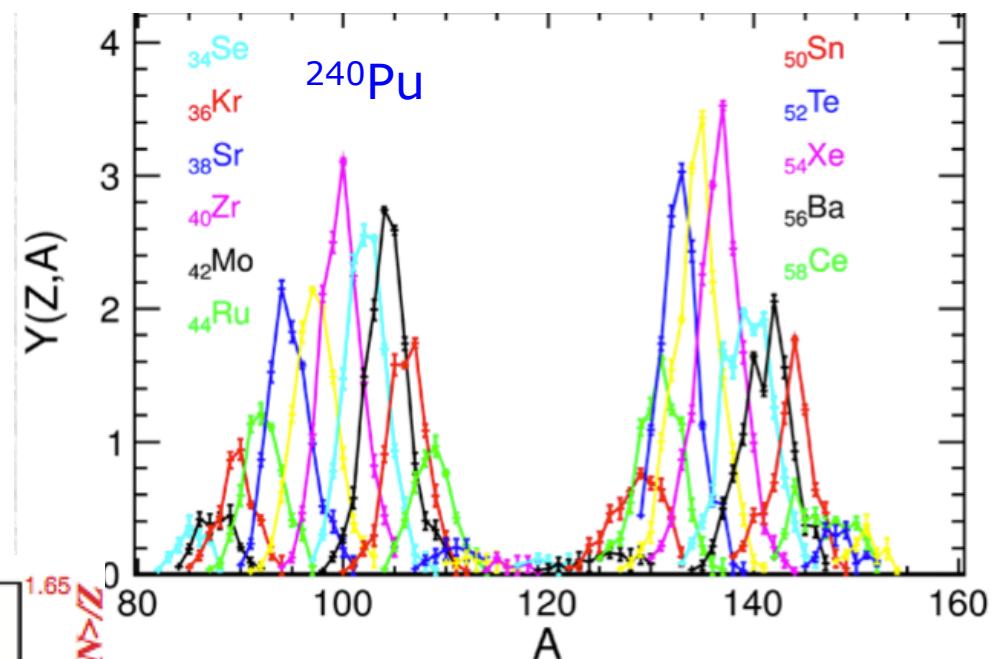
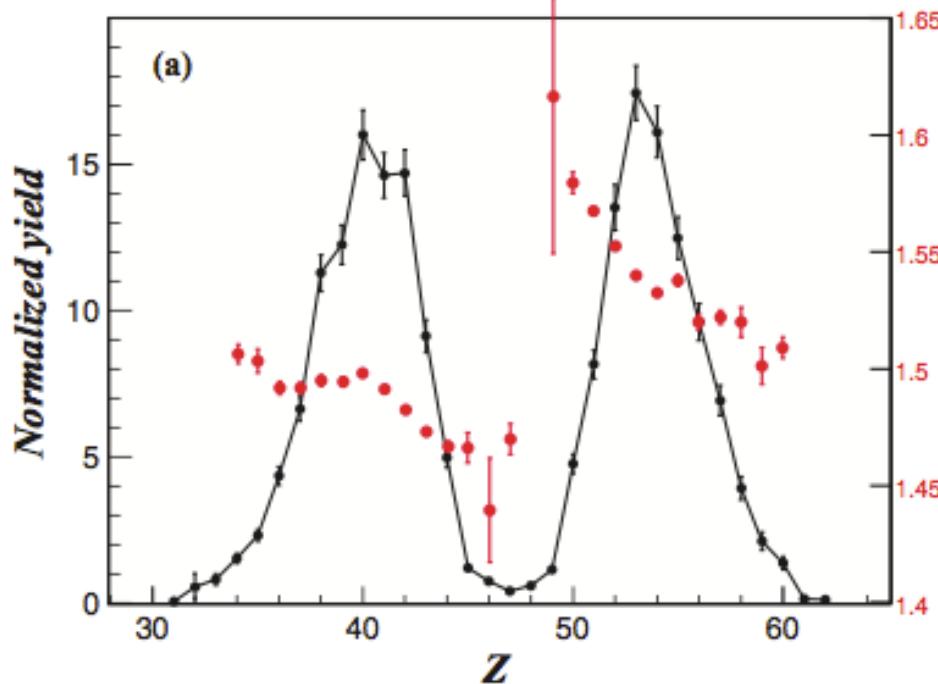
# Fission fragment identification



# Isotopic distributions of fission fragments induced in 2 proton transfer

$^{238}\text{U}(\text{C}^{12}, \text{Be}^{10})^{240}\text{Pu}$   $E^* \sim 9 \text{ MeV}$

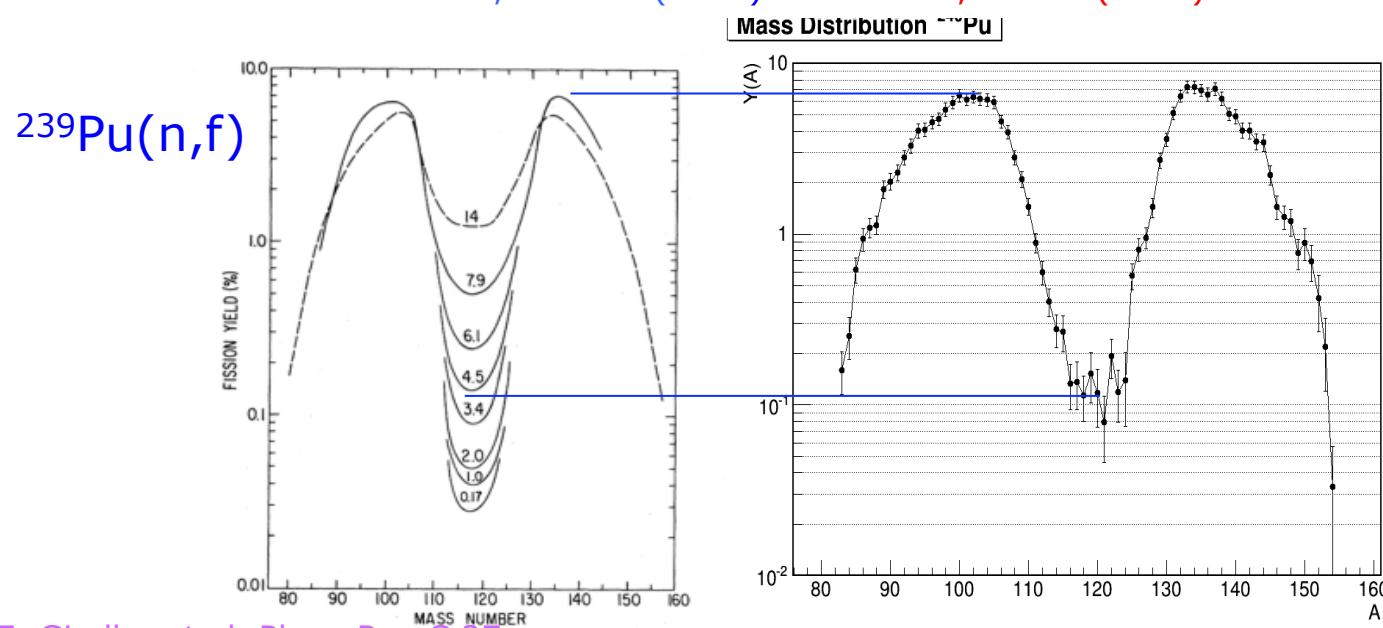
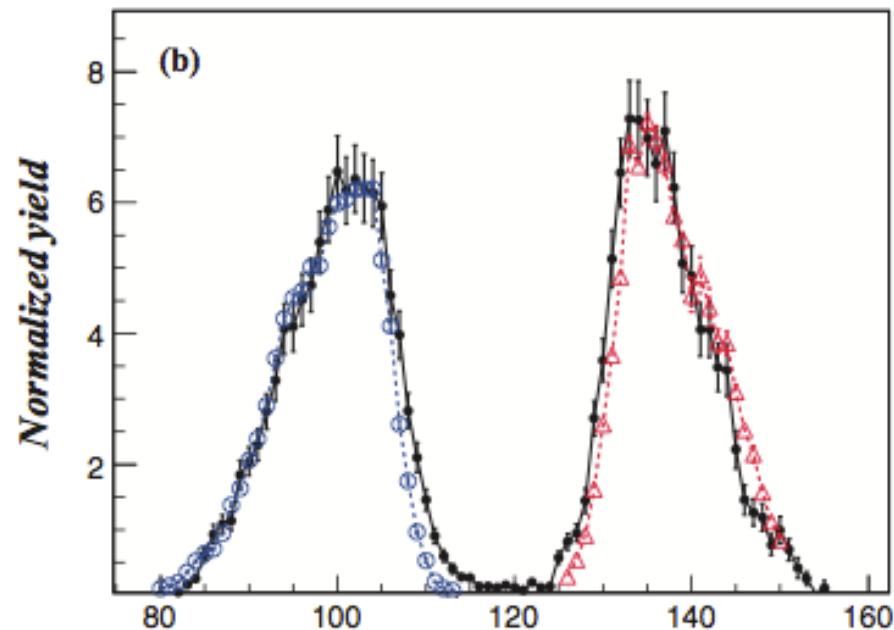
M. Caamaño et al., PRC 88 (2013) 024605



Neutron excess

$$\langle N \rangle(Z) = \frac{\sum_A A Y(Z, A)}{\sum_A Y(Z, A)} - Z.$$

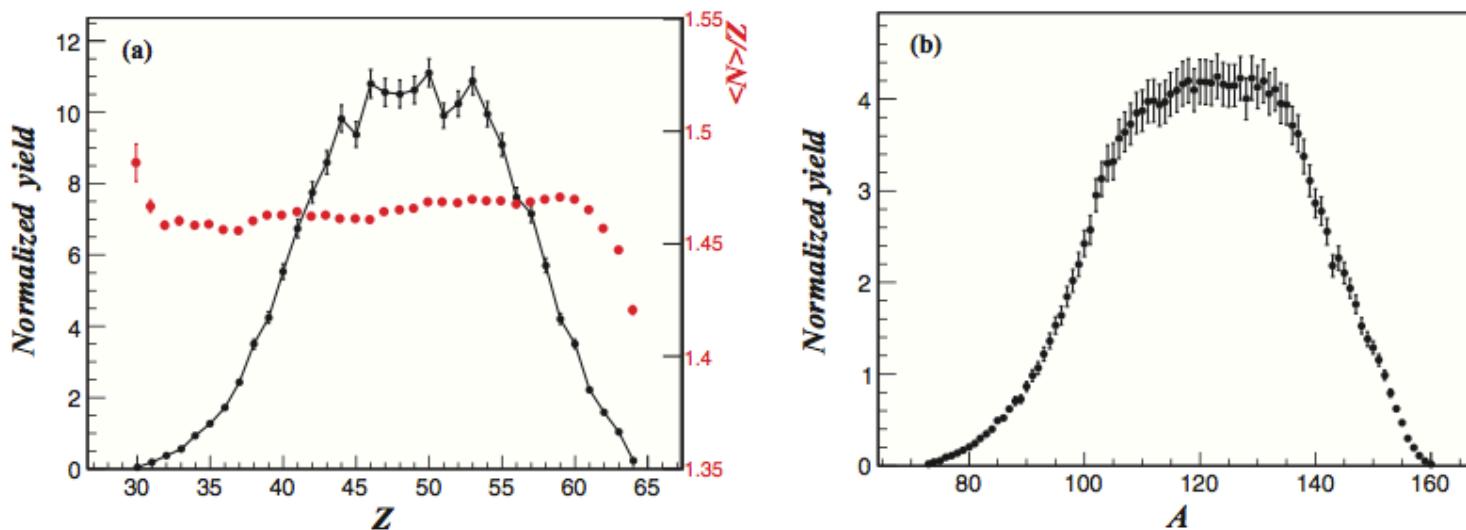
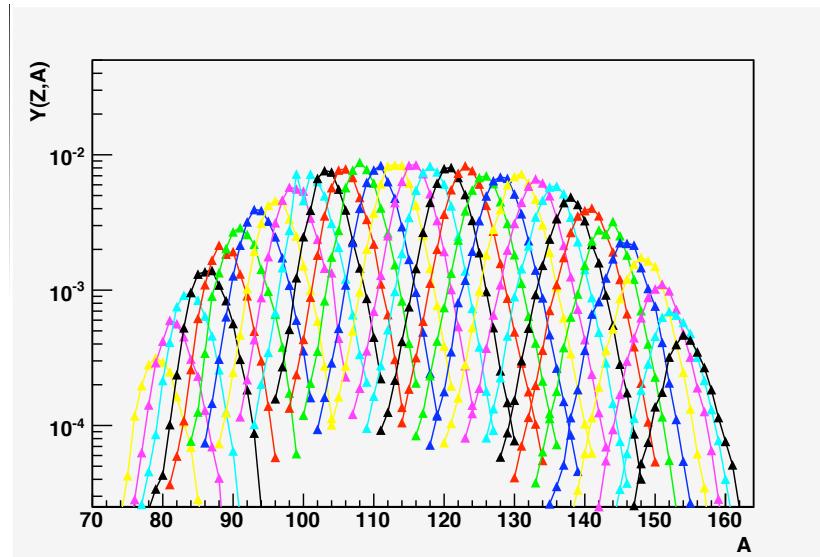
## Mass distribution for $^{240}\text{Pu}$ $E^* \sim 9$ MeV



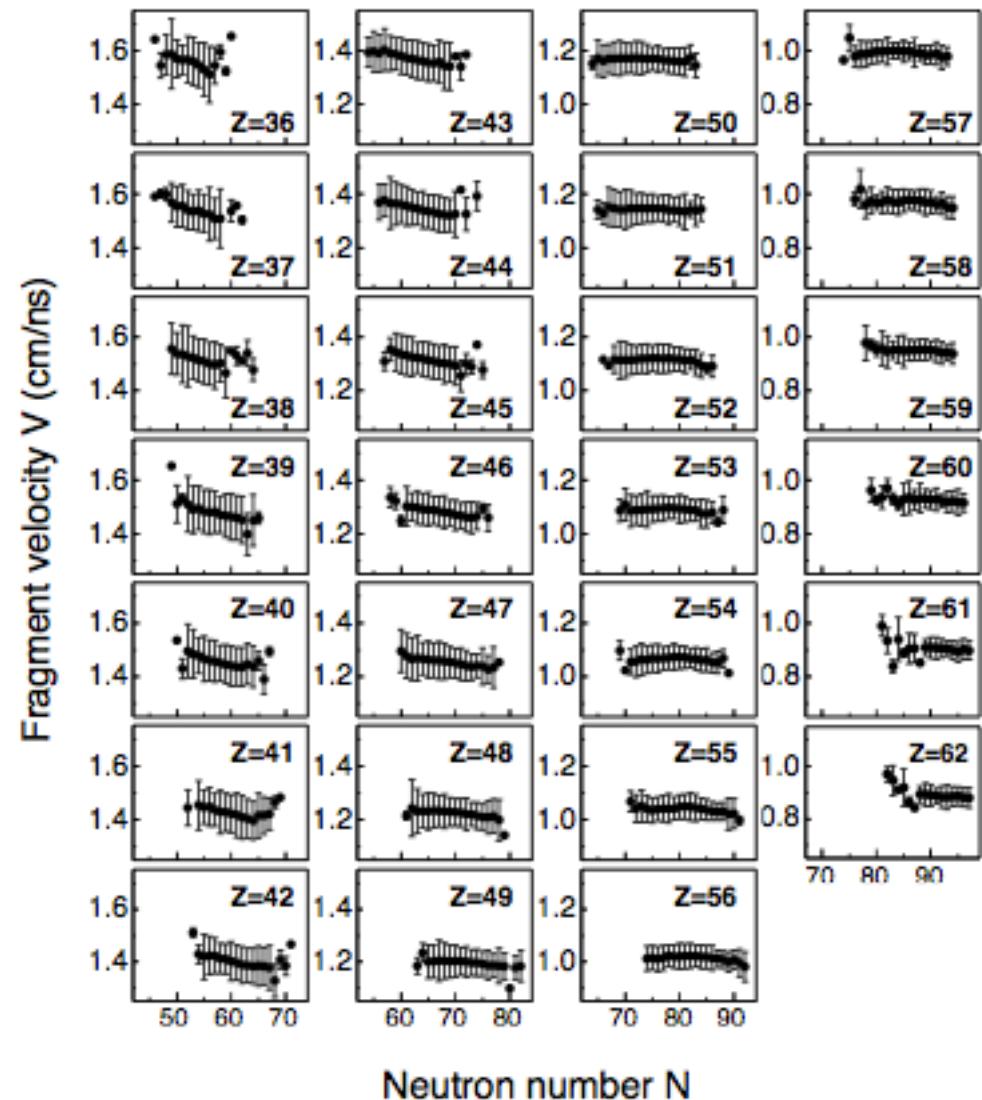
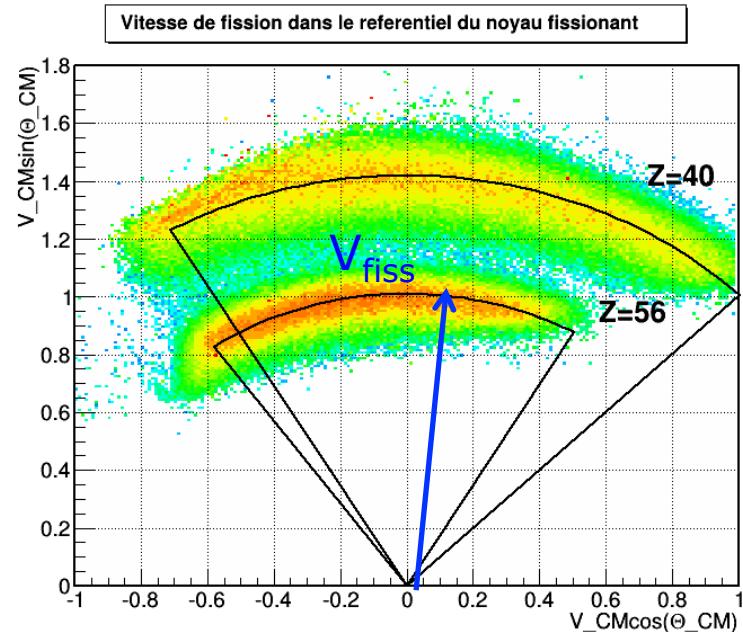
$E^* \sim 9$  MeV  
 $Sn = 5.6$  MeV  
 $En \sim 4$  MeV

# Isotopic distributions of fission fragments induced in fusion

$^{238}\text{U} + ^{12}\text{C} \rightarrow ^{250}\text{Cf}$   $E^* \sim 45$  MeV

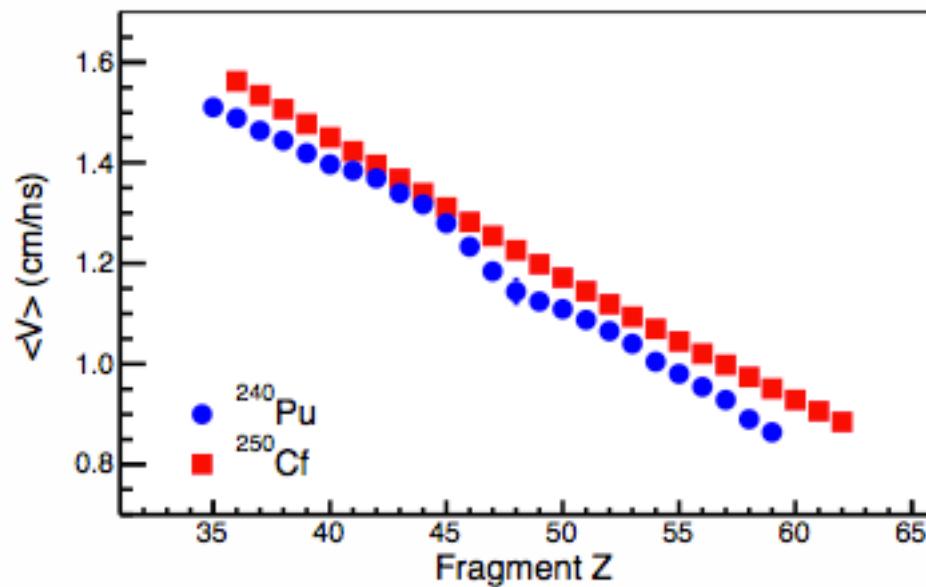


# Assets of the experimental set-up: Reconstruction of kinematical properties

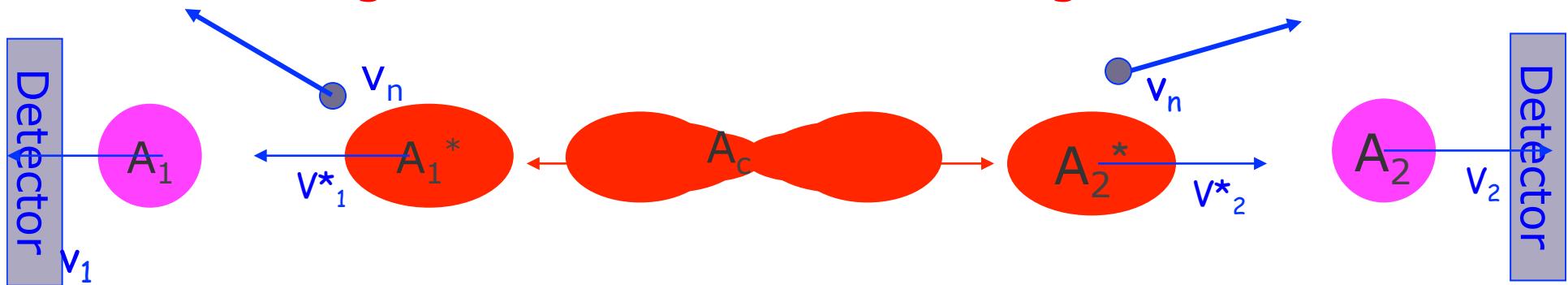


## Average velocity of fission fragments

$$\langle V \rangle (Z) = \frac{\sum_A Y(A,Z)V(Z,A)}{\sum_A Y(A,Z)}$$



## Recovering scission masses from fragment velocities



$$A_1^* v_1^* = A_2^* v_2^*$$

Momentum conservation

$$A_1^* + A_2^* = A_c$$

Mass conservation

$$\langle v_{1,2}^* \rangle = \langle v_{1,2} \rangle$$

Isotropic evaporation

$$\langle v_1 \rangle / \langle v_2 \rangle = A_2^* / A_1^* \quad A_1^* = A_c (v_1 / (v_1 + v_2))$$

$$A_1^* = A_c (v_1 / (v_1 + v_2))$$

$$A_2^* = A_c - A_1^*$$

$$A_1^* = A_c (v_1 / (v_1 + v_2))$$

# Reconstruction of the scission fragment masses $A^*$

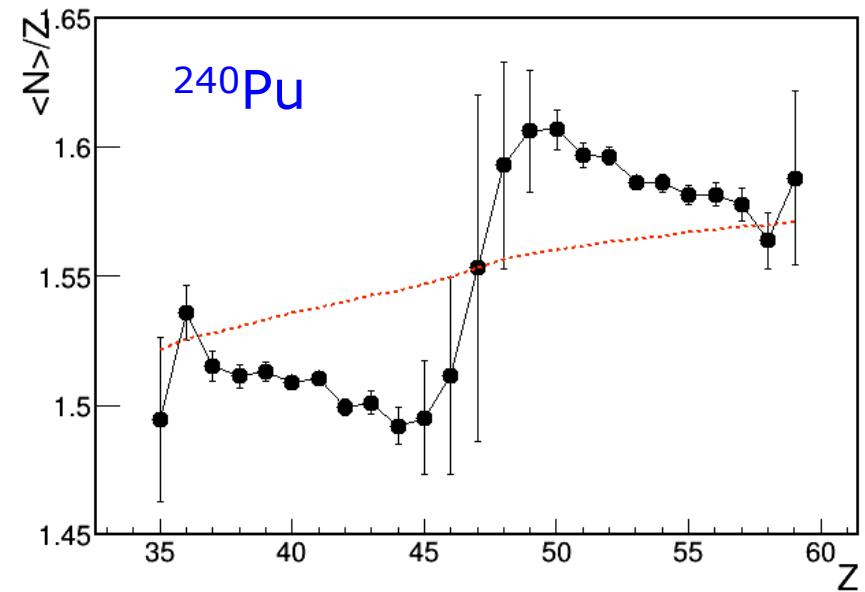
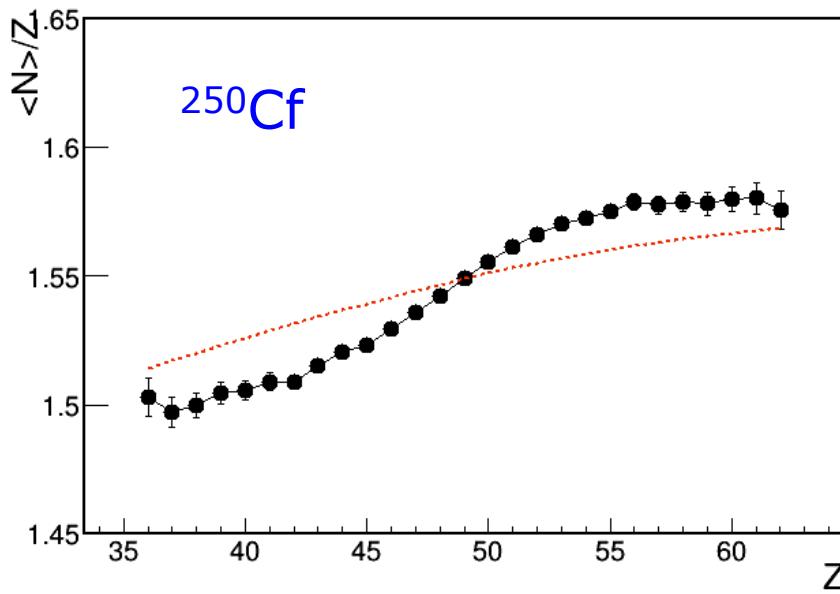
$$\frac{V_1}{V_2} = \frac{A_2^*}{A_1^*}$$

Momentum conservation

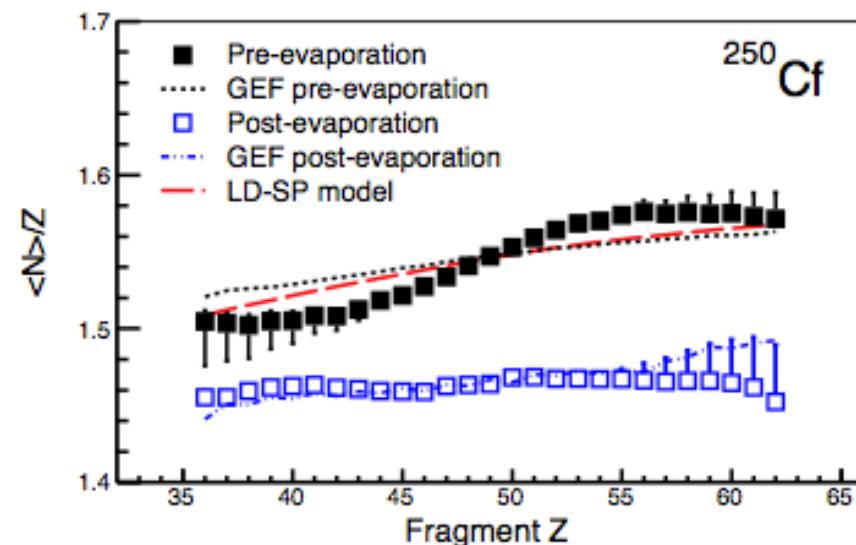
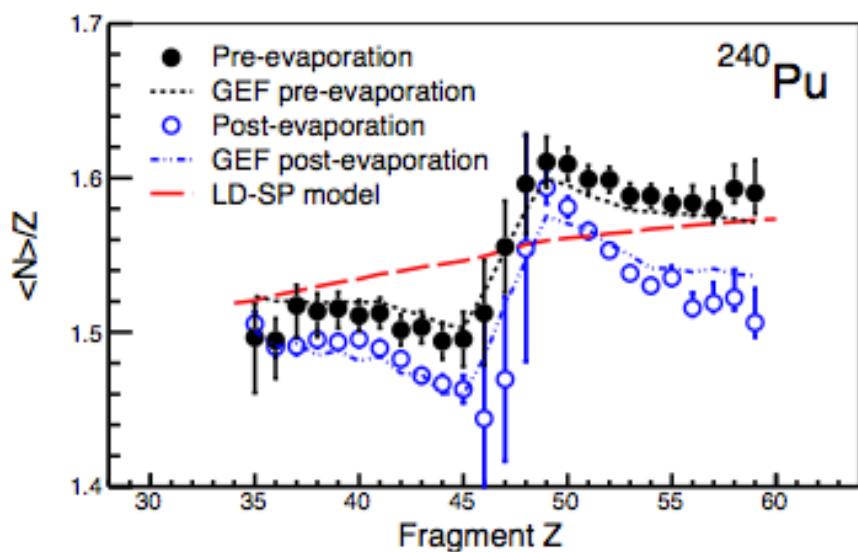
$$\begin{aligned} < A^* >_1 &= A_{FS} \frac{< V_2 >}{< V_1 >} \\ < A^* >_2 &= A_{FS} - < A^* >_1 \end{aligned}$$

$$Z_2 = Z_{FS} - Z_1 \quad \text{Charge conservation}$$

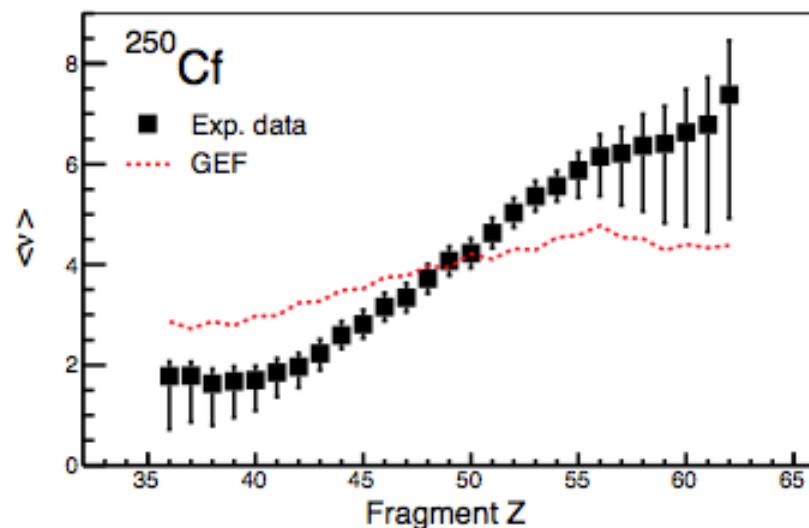
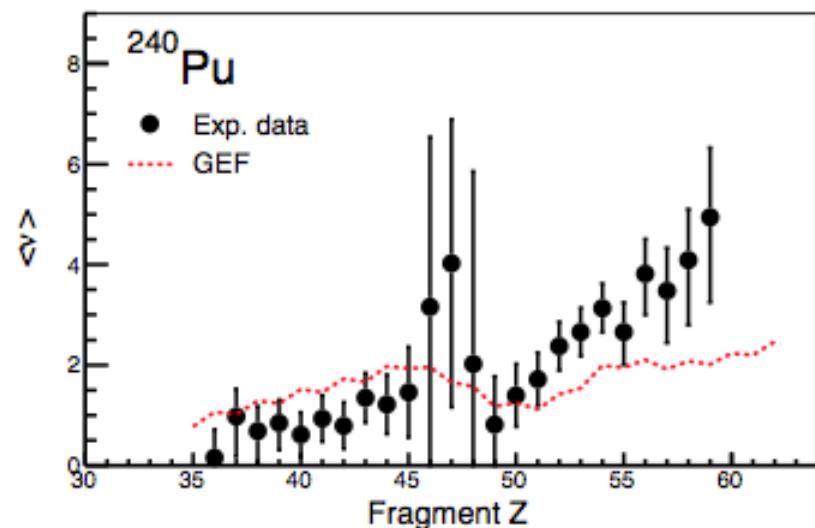
Neutron excess of the fragments at scission



## Average neutron excess @ scission



## Average neutron multiplicities @ scission



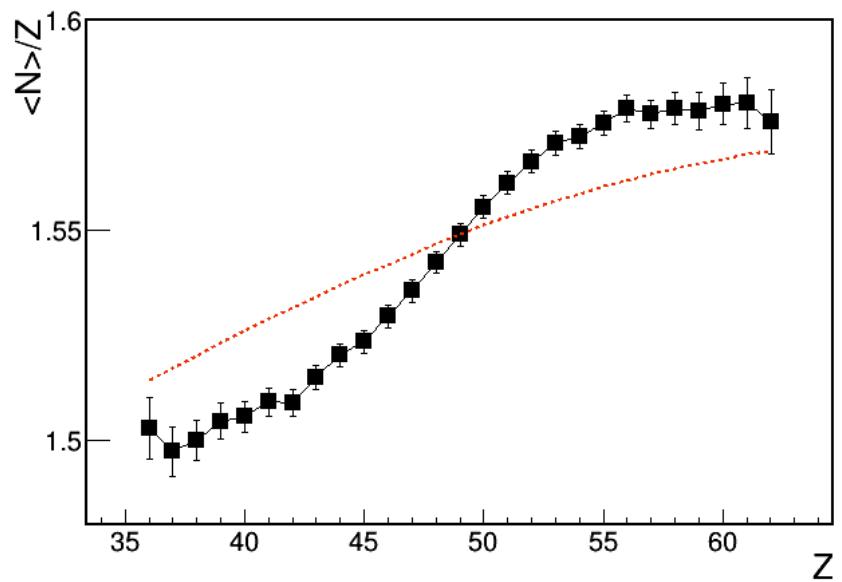
## Scission point model: minimization of the total potential energy

$$\begin{aligned}
 V(N_1, Z_1, \beta_1, N_2, Z_2, \beta_2, \tau, d) = & V_{LD_1}(N_1, Z_1, \beta_1) + V_{LD_2}(N_2, Z_2, \beta_2) \\
 & + S_1(N_1, \beta_1, \tau) + S_1(Z_1, \beta_1, \tau) + S_2(N_2, \beta_2, \tau) + S_2(Z_2, \beta_2, \tau) \\
 & + P_1(N_1, \beta_1, \tau) + P_1(Z_1, \beta_1, \tau) + P_2(N_2, \beta_2, \tau) + P_2(Z_2, \beta_2, \tau) \\
 & + V_C(N_1, Z_1, \beta_1, N_2, Z_2, \beta_2, d) + V_n(N_1, Z_1, \beta_1, N_2, Z_2, \beta_2, d),
 \end{aligned}$$

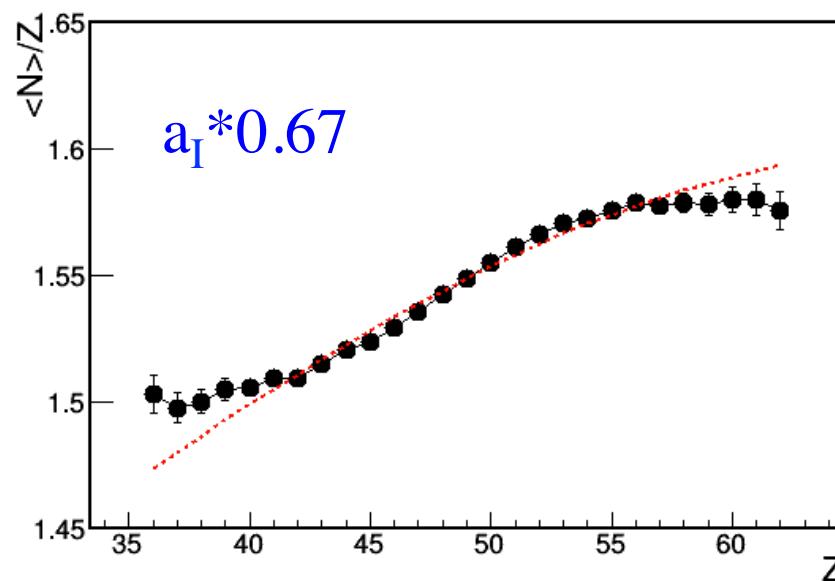
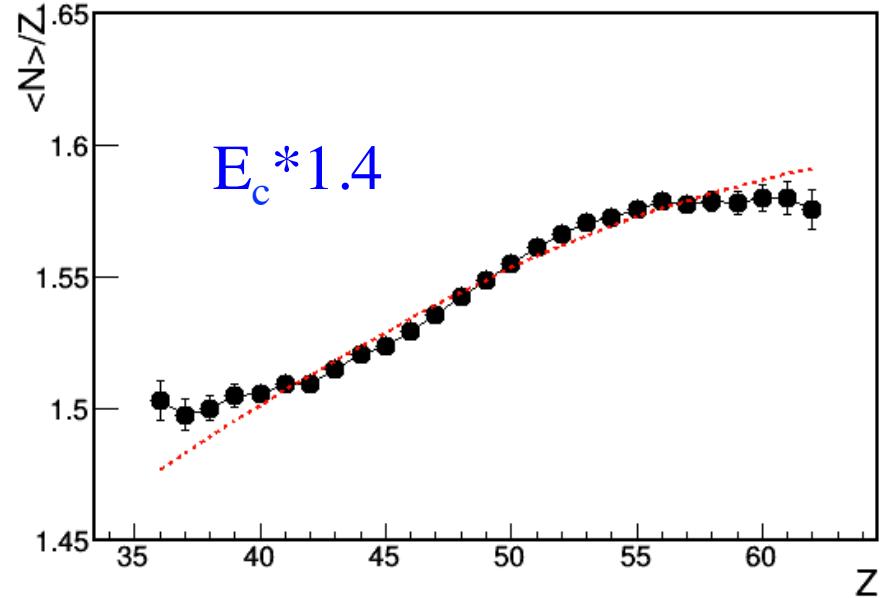
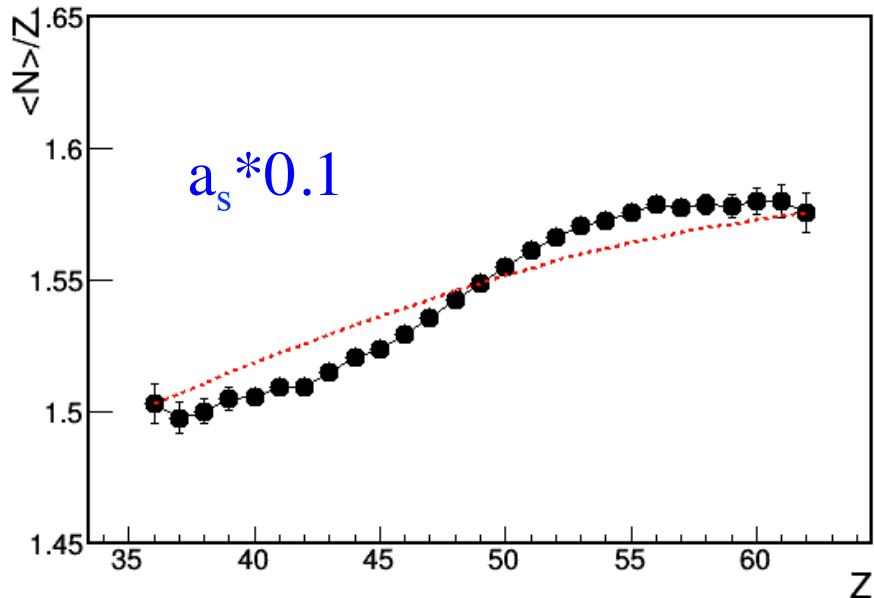
$^{250}\text{Cf}$   $E^*=45$  MeV :  
only liquid-drop terms play a role (shell effects disappeared)

$$\begin{aligned}
 V_{LD}(Z, N, \beta) = & a_a A - a_s A^{2/3} (1 + 0.4 \alpha^2) \\
 & - 1.78 I^2 (a_a A - a_s A^{2/3} (1 + 0.4 \alpha^2)) \\
 & + Z^2 ((0.705/A^{1/3}) (1 - 0.2 \alpha^2) - 1.15/A)
 \end{aligned}$$

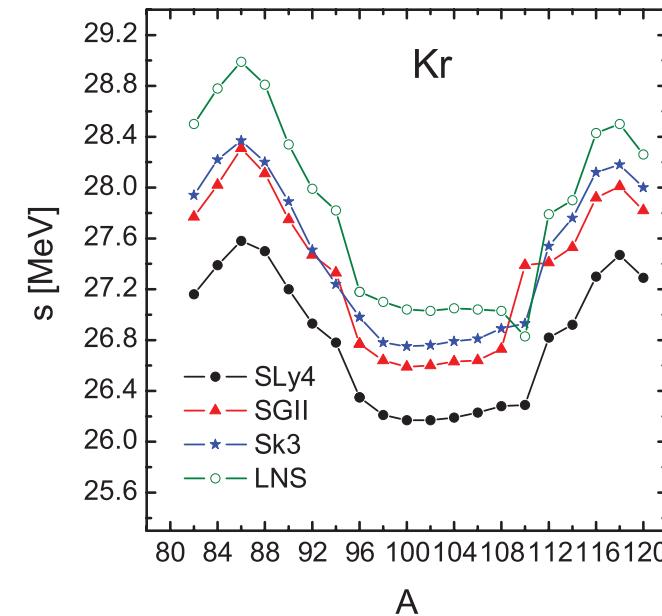
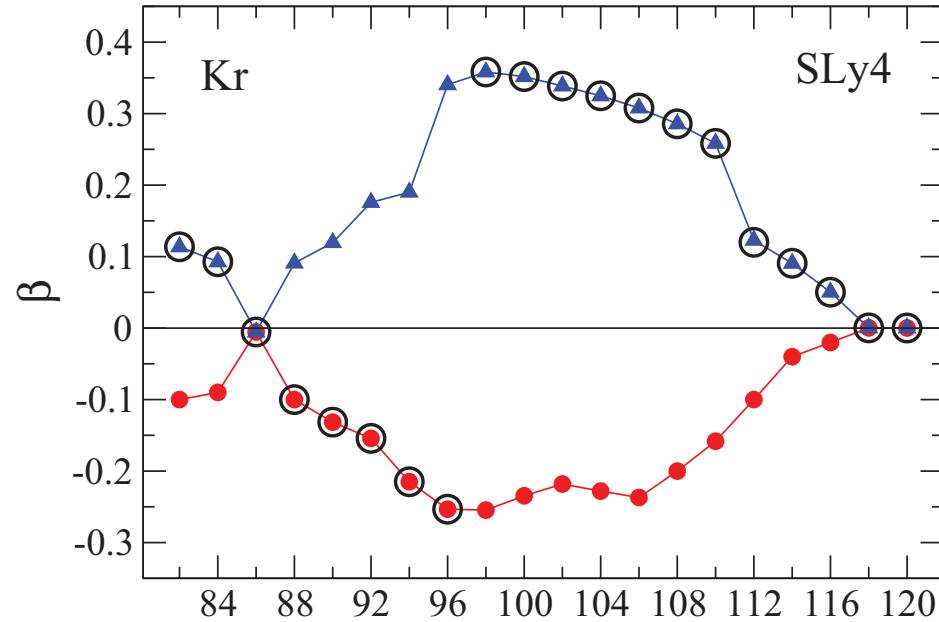
W.D. Myers, and W.J. Swiatecki, Ark. Fys., 36, 343, (1967)



## Scission point model: influence of different mass terms



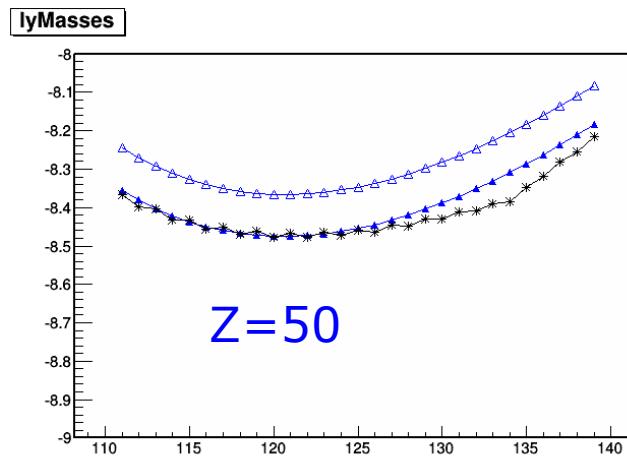
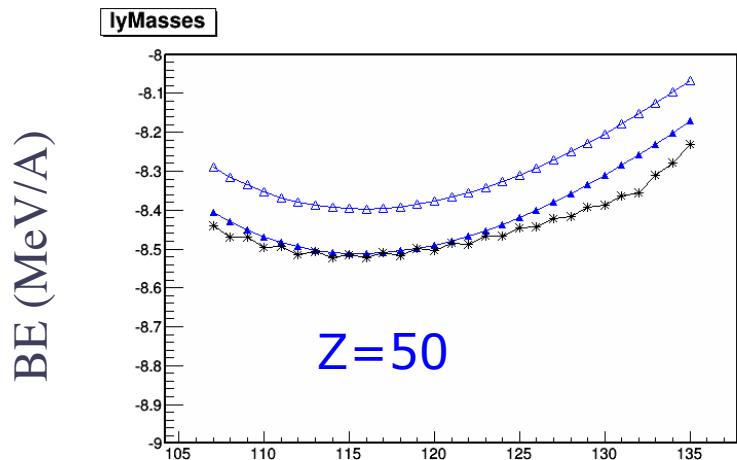
## Diminution of symmetry energy with deformation ?



Gaidarov et al., PRC 85 (2012) 064319

A diminution of 10% is predicted when deformation increases  
From 0 to 0.4  
⇒ What happens at scission deformation ??  
⇒ Effect of density ??

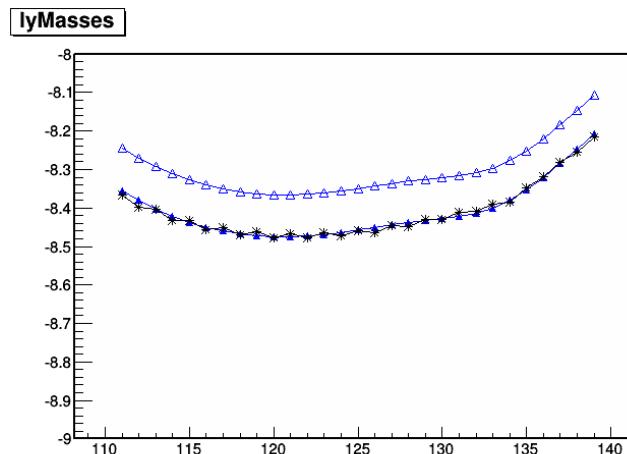
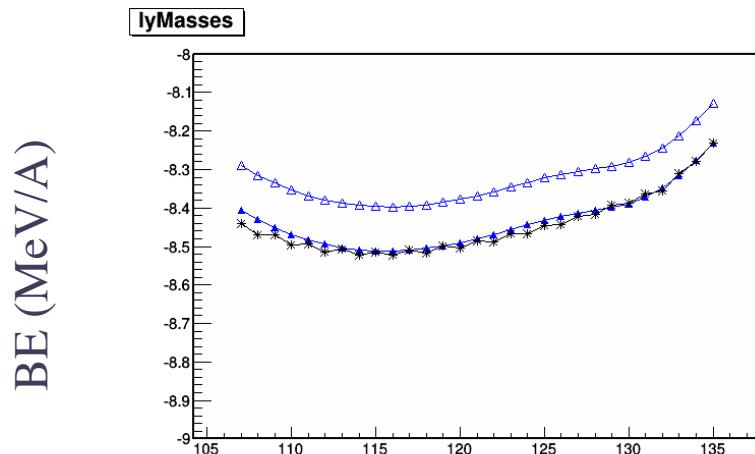
## Other explanation: Remaining of shell effects in BE



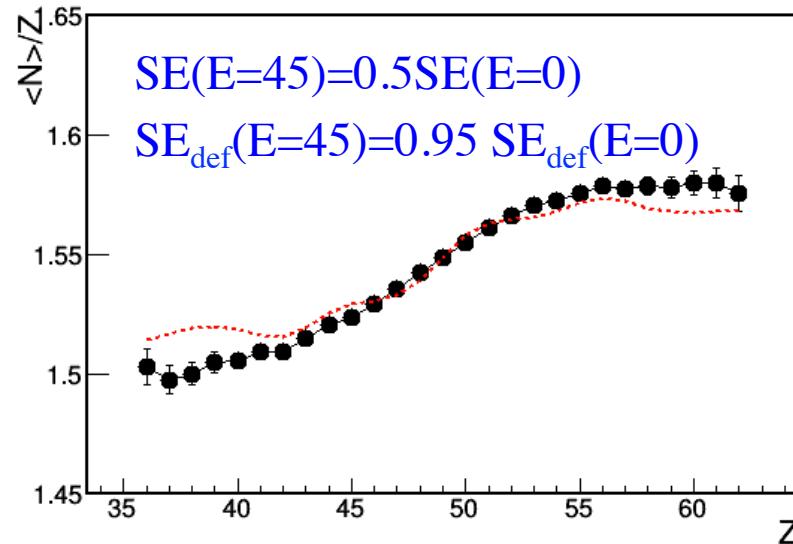
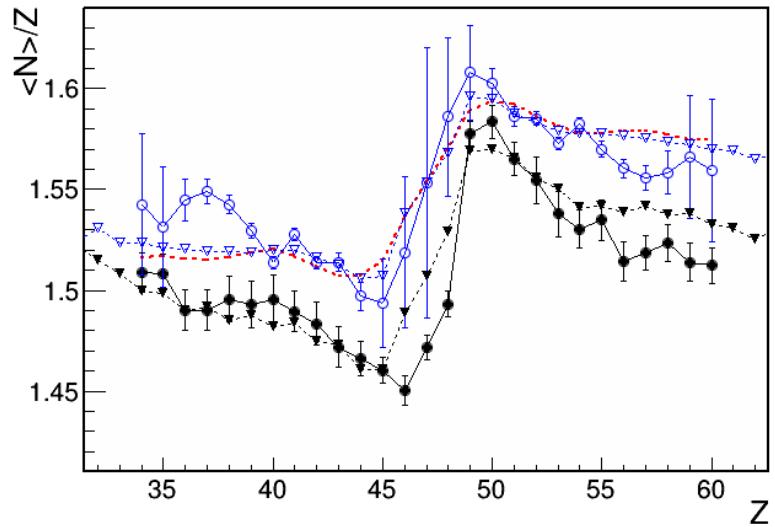
$$BE = BE + SE + SE_{def}$$

$$SE = 3 * \exp(-(Z-50)^2/(2*3^2)) * 4 * \exp(-(N-82)^2/(2*3.5^2))$$

$$SE_{def} = 3 * \exp(-(Z-54)^2/(2*3^2)) * 4 * \exp(-(N-90)^2/(2*3.5^2))$$

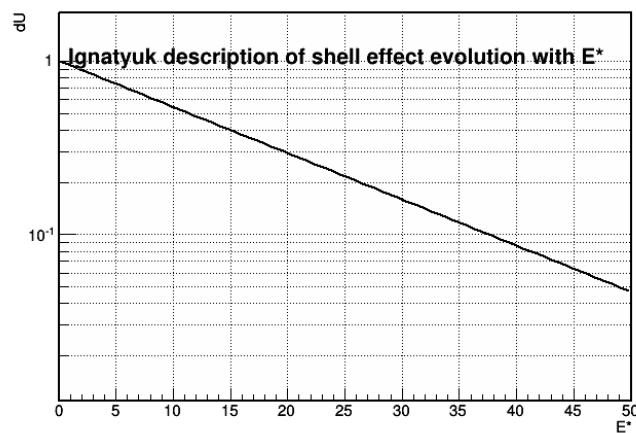


## Scission-point model with shell effects



Shell effects remain quite strong, even at  $E^*=45$  MeV ??

$$\exp(-1*(132/(3.2*\text{pow}(132,4./3)))*x[0])$$



# CONCLUSIONS

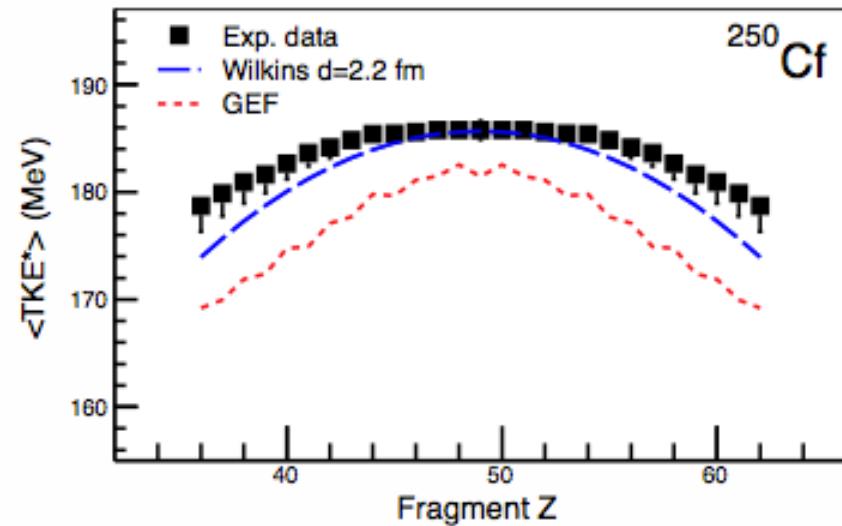
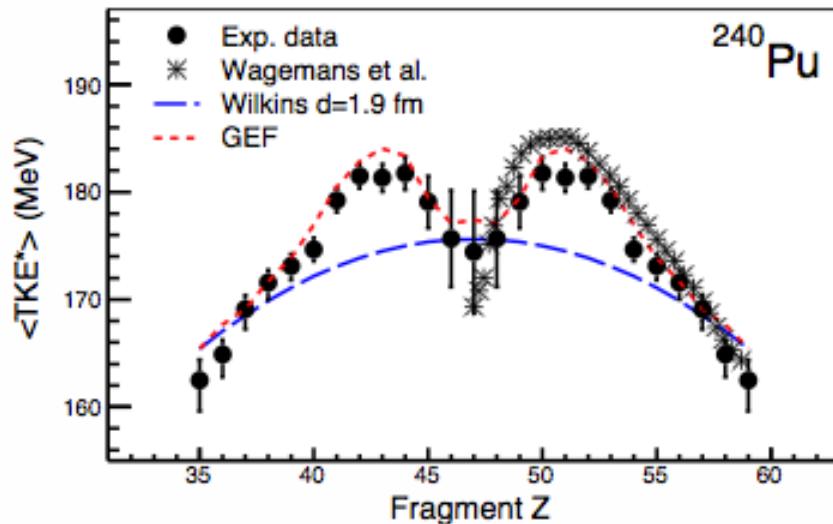
- Inverse kinematics is a powerful method
  - Broad range of actinides produced
  - Isotopic distribution
  - Kinematical properties
  - Access to the scission point !!
    - Neutron evaporation multiplicity
    - Neutron and proton sharing
  - Evidence for charge polarisation, even at moderate (high) excitation energy
  - Possibility of investigate the influence of excitation energy

Talks of  
L. Audouin,  
D. Ramos

Talk of  
D. Ramos

## Total kinetic energies

$$TKE(Z) = \frac{1}{2} < A_1^* > < V_1 >^2 + \frac{1}{2} < A_2^* > < V_2 >^2$$



$$TKE = 1.44 \frac{Z_h Z_l}{r_0(A_h^{1/3}(1 + \frac{2}{3}\beta_h) + A_l^{1/3}(1 + \frac{2}{3}\beta_l)) + d}.$$

Deformation at scission !!

Talk of  
M. Caamaño